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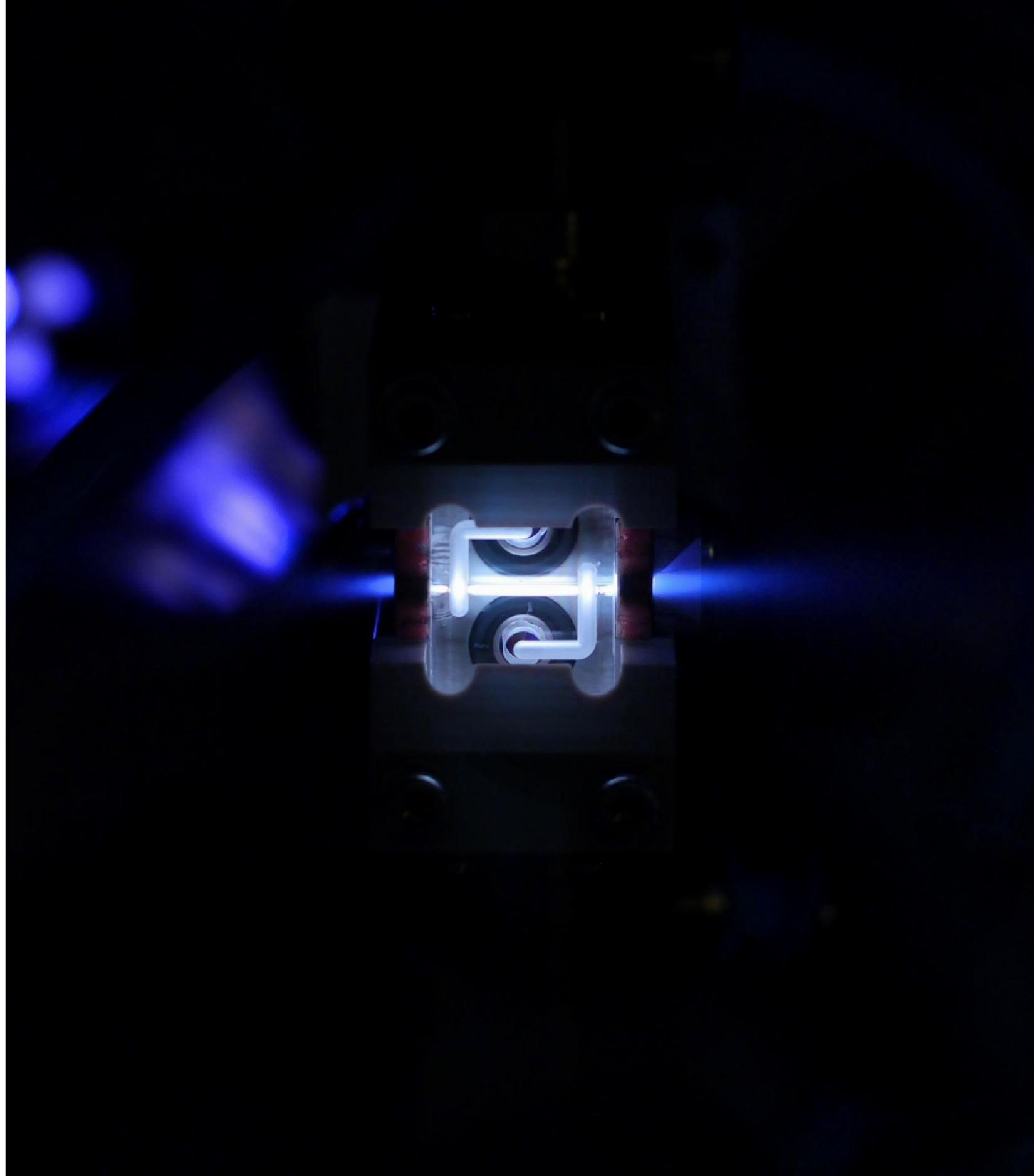
Betatron Radiation and Self-Corrected Energy Spread in Plasma-Wakefield Accelerators

ALEGRO workshop 2025

Daniel Kalvik

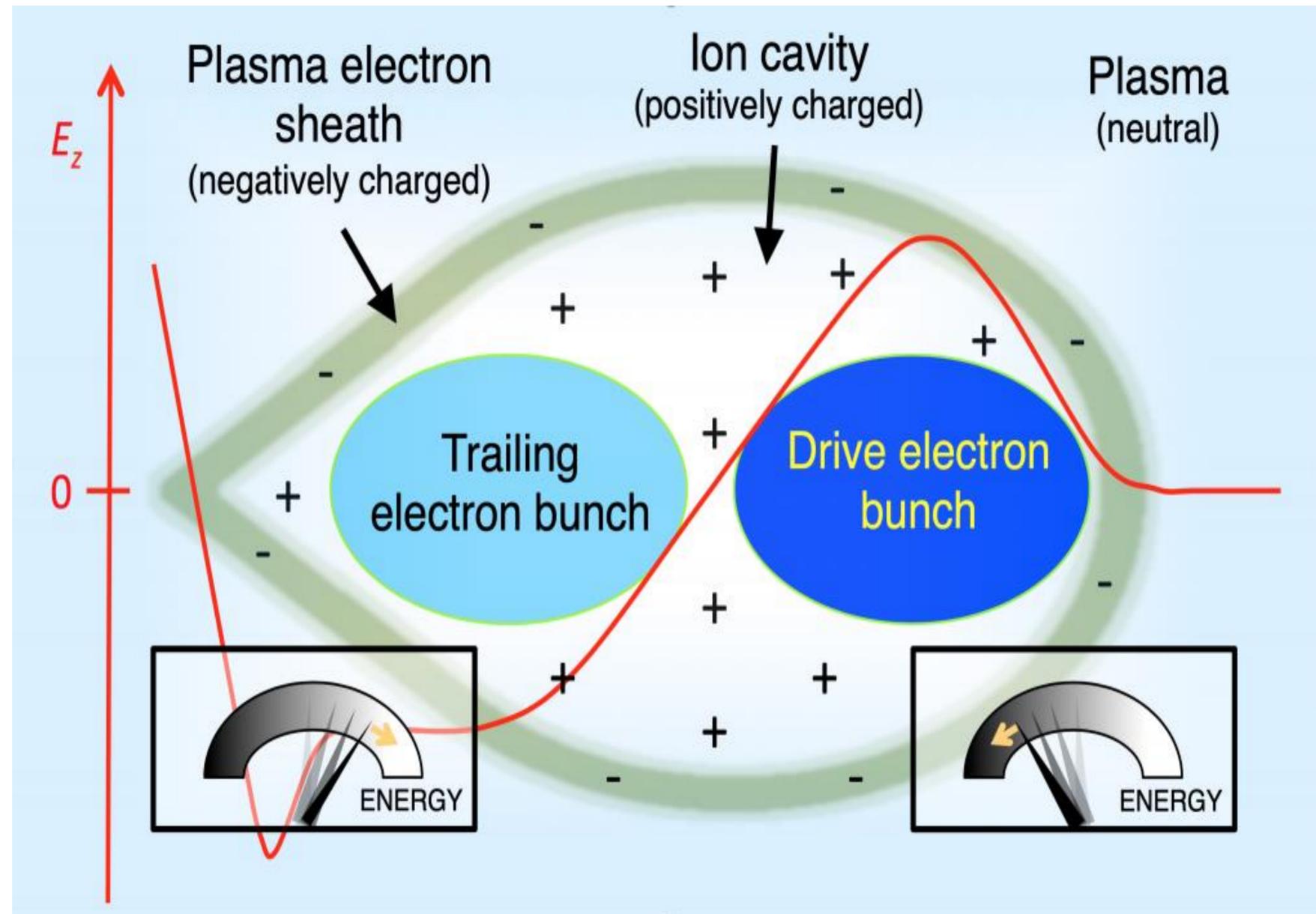
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12.06.24



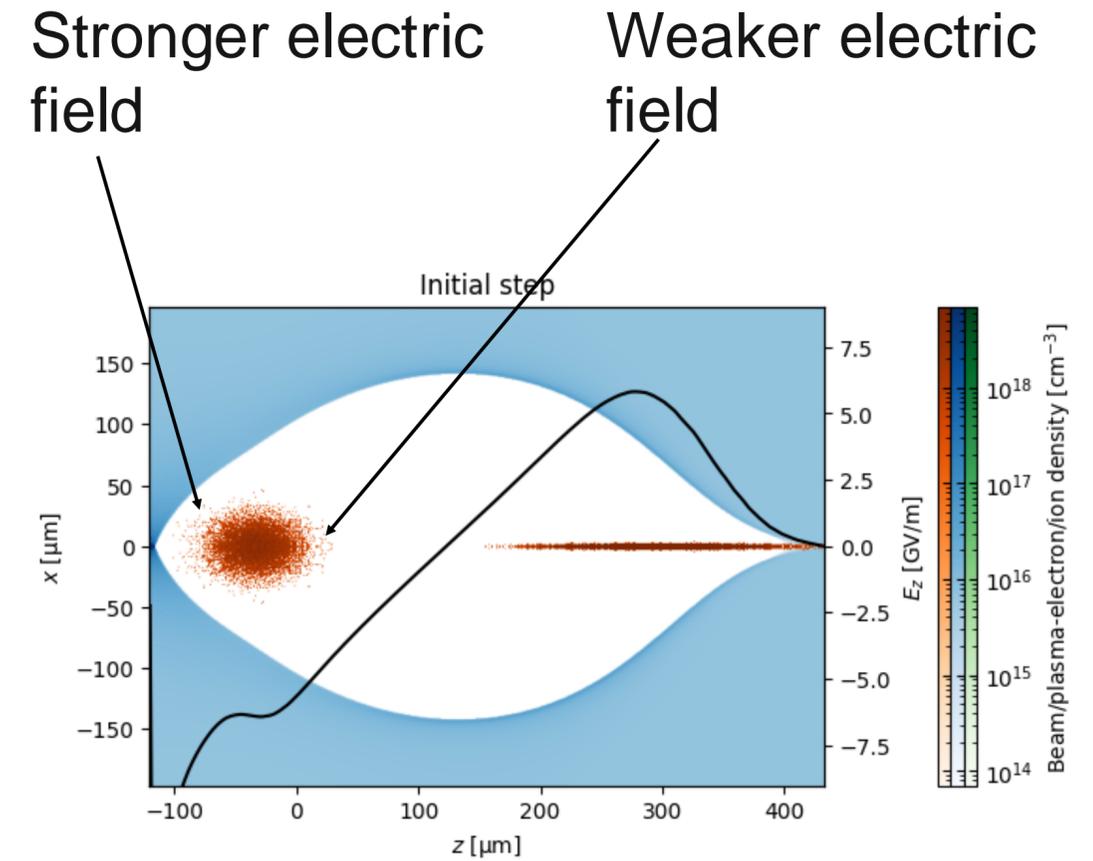
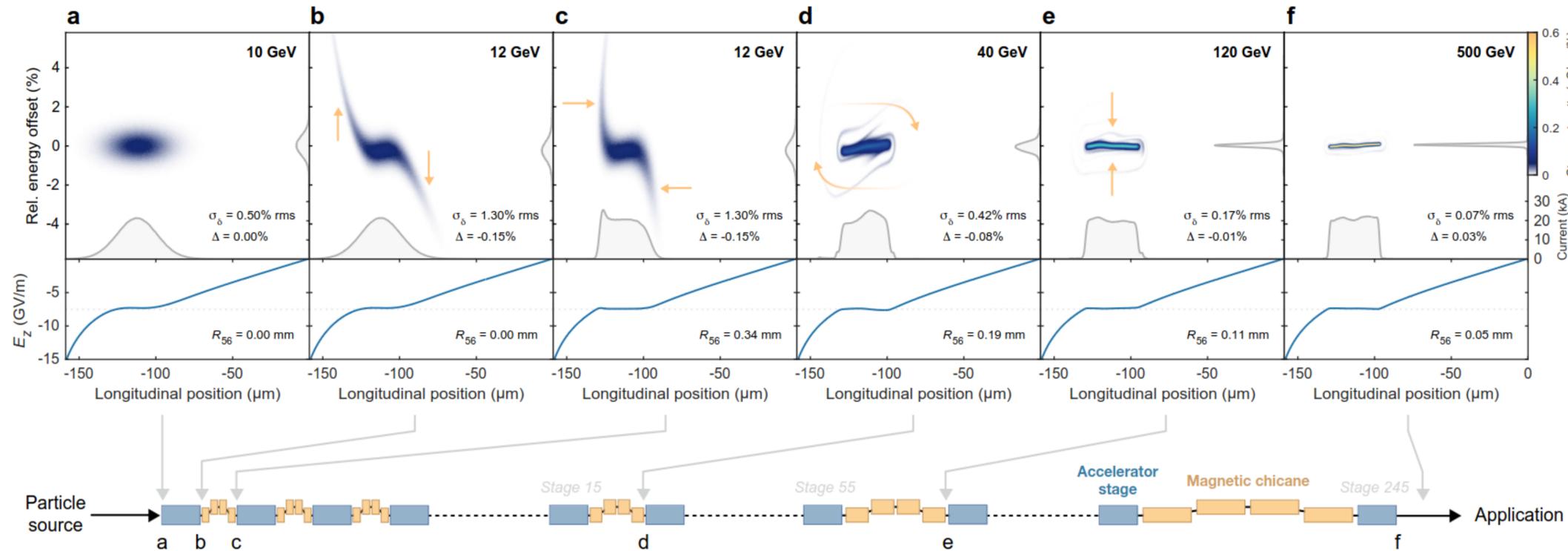
Induced energy spread

- The longitudinal electric field generates a large energy spread.



Self-correction mechanism

How to reduce the energy spread



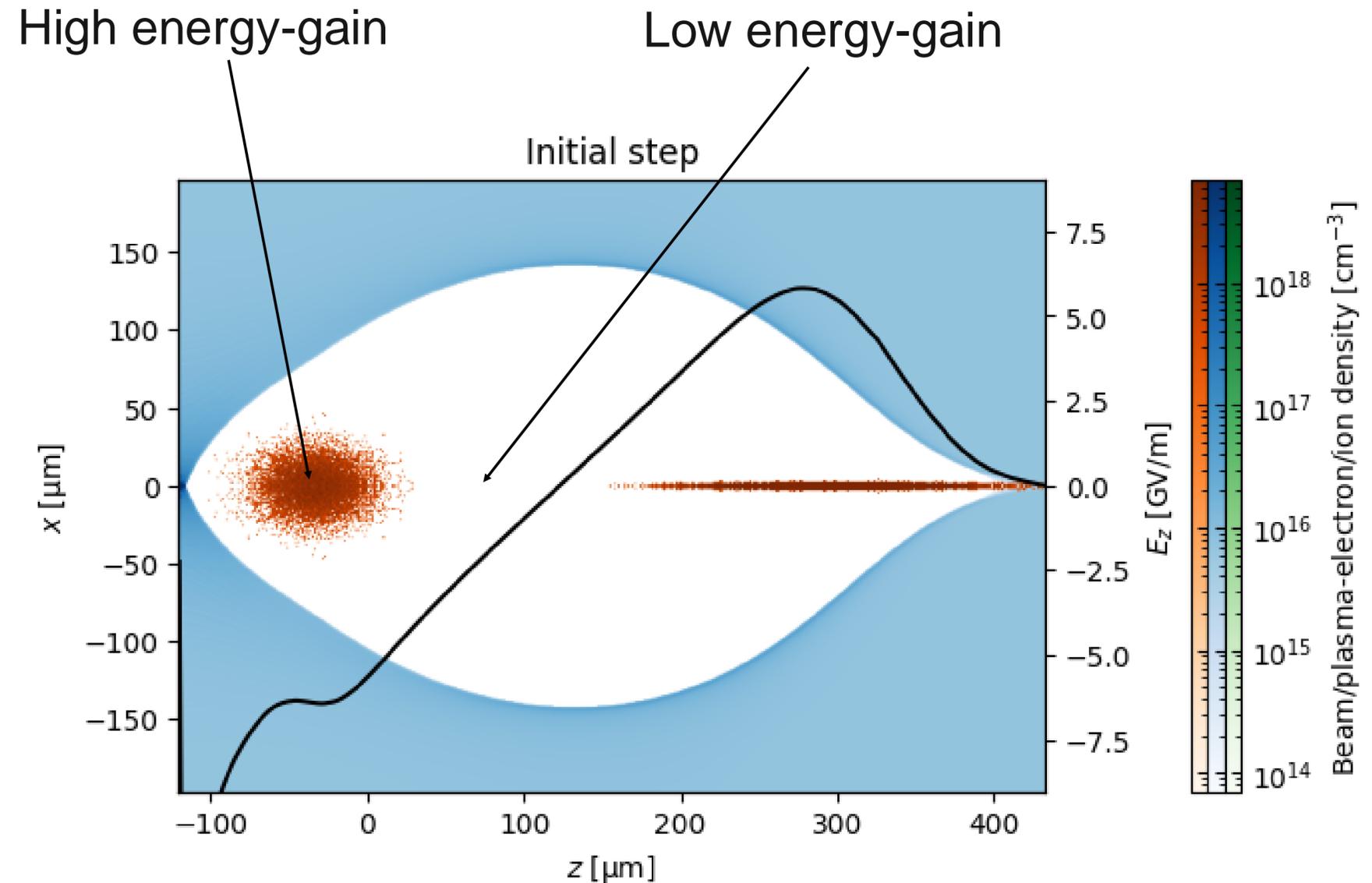
Carl A Lindstrøm. 'Self-correcting longitudinal phase space in a multistage plasma accelerator'. eng. In: *arXiv (Cornell University)* (2021). ISSN: 2331-8422.

- > The longitudinal dynamics in bending magnets can reduce the energy spread.
- > High energy particles are bent less (they take a longer path), low energy particles are bent more, they take a longer path.

Self-correction mechanism

How to reduce the energy spread

- > The strength of the correction is tuned to a specific location in the wake.
- > I.e. the energy gain (decided by location of the beam) is predetermined.



Particle dynamics

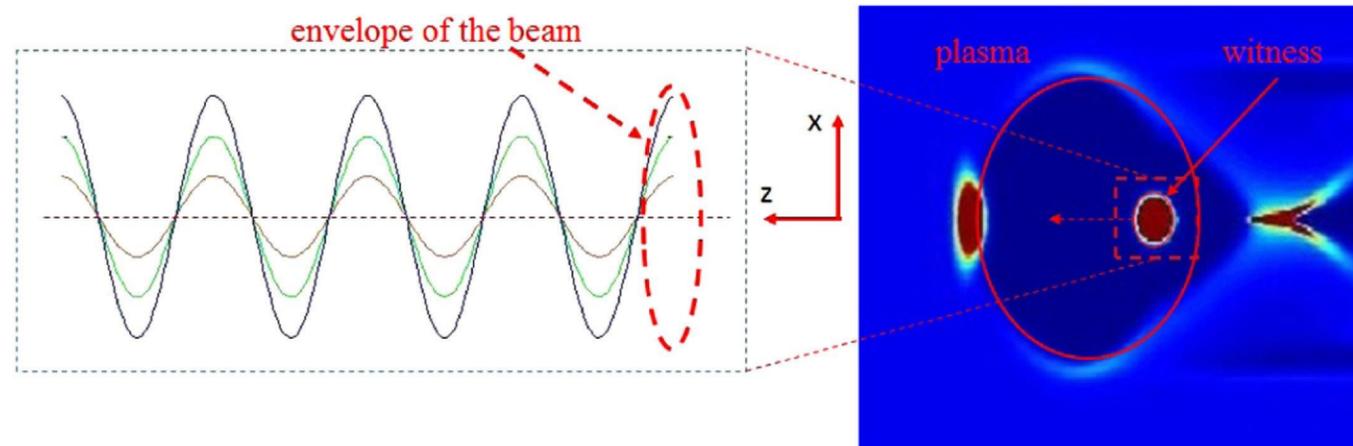
Motion in an ion cavity

> Newton's second law gives the equations of motion

$$\mathbf{a} + \frac{1}{\gamma} \frac{d\gamma}{dt} \mathbf{v} - \frac{\mathbf{F}}{m\gamma} = 0$$

> If the Lorentz factor is increasing, the oscillations are damped. This is known as adiabatic damping.

> In an ion cavity, the equation of motion becomes a harmonic oscillator

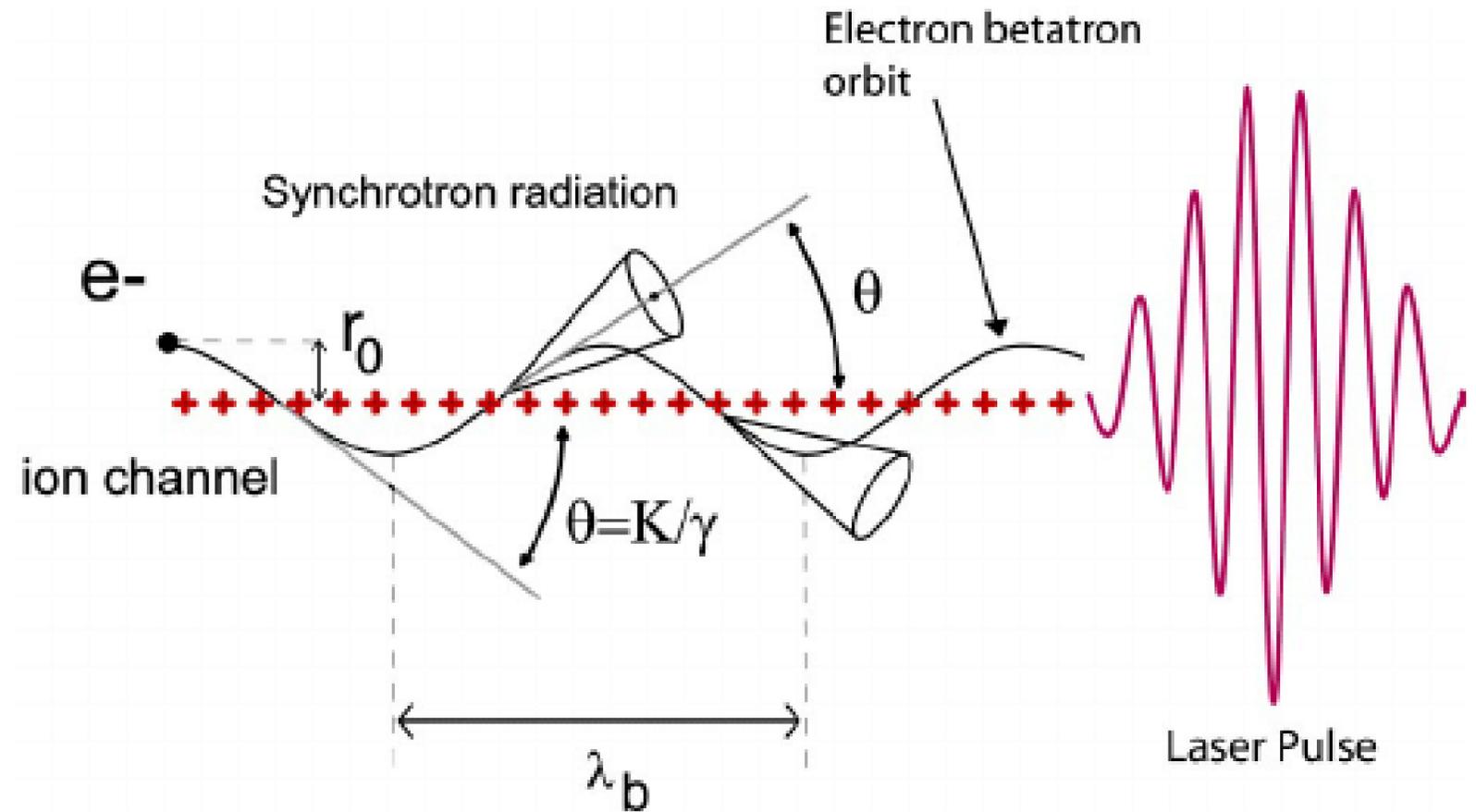


V. Shpakov et al. 'Plasma acceleration limitations due to betatron radiation'. In: *Nuclear Instruments and Methods in Physics Research Section A: Accelerators, Spectrometers, Detectors and Associated Equipment* 909 (2018). 3rd European Advanced Accelerator Concepts workshop (EAAC2017), pp. 463–466. ISSN: 0168-9002. DOI: <https://doi.org/10.1016/j.nima.2018.02.058>. URL: <https://www.sciencedirect.com/science/article/pii/S0168900218302109>.

The radiation reaction

Interaction between a charged particle and its own electromagnetic field

- > As charged particles oscillate, a force is exerted on the particle by its own electromagnetic field; hence it is a self-force.
- > This self-force is called the radiation reaction force, and its effect, the radiation reaction.
- > Particles of different amplitudes radiate differently and contributes therefore to energy spread.



Ta, Phuoc & Burgy, Frédéric & Rousseau, Jean-Philippe & Malka, Victor & Rousse, Antoine & Shah, Rahul & Umstadter, Donald & Pukhov, A. & Kiselev, Sergei. (2005). Laser based synchrotron radiation. *Physics of Plasmas*. 12. 10.1063/1.1842755.

Particle dynamics and the Larmor formula

Radiation for an oscillating charged particle

> The Larmor formula, $P = -\frac{e^2}{6\pi\epsilon_0} \frac{c}{(m_e c^2)^2} \left(\frac{dp^\mu}{d\tau}\right)^2$ *, tells us how much energy is radiated by accelerating charges.

> Simplifies to $P = \frac{e^2}{6\pi\epsilon_0 c} \gamma^4 \left(\frac{d\beta}{dt}\right)^2$, which gives $\bar{P} = r_e m_e c^3 \gamma^2 k_p^4 r_\beta^2 / 12$ **, which is the radiated power, averaged over one betatron period.

Amplitude of oscillations



> The radiation emitted as the particles oscillate, is betatron radiation.

* 'LXIII. On the theory of the magnetic influence on spectra; and on the radiation from moving ions'. eng. In: *The London, Edinburgh and Dublin philosophical magazine and journal of science* (). ISSN: 1941-5982.

* E. Esarey et al. 'Synchrotron radiation from electron beams in plasma-focusing channels'. eng. In: *Physical review. E, Statistical, nonlinear, and soft matter physics* 65.5 (2002), pp. 056505/15–056505. ISSN: 1539-3755.

Radiation in both transverse dimensions

Radiation for oscillating charged particles

- > Expanding to 2D oscillations the new mean radiated power is

$$\bar{P} = \frac{e^2 c}{48\pi\epsilon_0} \gamma^2 k_p^4 (r_x^2 + r_y^2). \text{ The equation is no longer dependent on the betatron}$$

amplitude. The amplitude in x/y give independent contributions.

- > By averaging over all the betatron amplitudes for a bunch of particles, the total radiated power is

$$\langle \bar{P} \rangle = \frac{e^2 c}{6\pi\epsilon_0} k_p^3 \left(\frac{\gamma}{2} \right)^{3/2} (\epsilon_{nx} + \epsilon_{ny}) \longleftarrow \frac{1}{\beta_m} = \frac{k_p}{\sqrt{2}\gamma}, \text{ and } \frac{\sigma_x^2}{\beta_m} = \frac{\epsilon_{nx}}{\gamma}$$

- > And the spread in the radiation is $\sigma_{\bar{P}} = \frac{e^2 c}{24\pi\epsilon_0} (\sqrt{\gamma} k_p)^3 \sqrt{\epsilon_{nx}^2 + \epsilon_{ny}^2}$ (for no energy spread).

Self-consistent motion for accurate simulations

Including the radiation reaction and adiabatic damping

- The equations of motion for each individual particle are (Deng et al.)

$$K^2 = \frac{k_p^2}{2}$$

$$\ddot{x} + \left(\frac{\omega_p}{\gamma} \frac{E_z}{E_0} + \tau_R c^2 K^2 \right) \dot{x} + \frac{c^2 K^2}{\gamma} x = 0$$

$$\ddot{y} + \left(\frac{\omega_p}{\gamma} \frac{E_z}{E_0} + \tau_R c^2 K^2 \right) \dot{y} + \frac{c^2 K^2}{\gamma} y = 0$$

$$\dot{\gamma} = \omega_p \frac{E_z}{E_0} - \tau_R c^2 K^4 \gamma^2 (x^2 + y^2)$$

Self-consistent motion for accurate simulations

Including the radiation reaction and adiabatic damping

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Adiabatic damping

$$\ddot{x} + \left(\frac{\omega_p}{\gamma} \frac{E_z}{E_0} + \tau_R c^2 K^2 \right) \dot{x} + \frac{c^2 K^2}{\gamma} x = 0$$

$$\ddot{y} + \left(\frac{\omega_p}{\gamma} \frac{E_z}{E_0} + \tau_R c^2 K^2 \right) \dot{y} + \frac{c^2 K^2}{\gamma} y = 0$$

Damping from radiation reaction force

$$\dot{\gamma} = \omega_p \frac{E_z}{E_0} - \tau_R c^2 K^4 \gamma^2 (x^2 + y^2)$$

> My results are obtained by solving this set of equations, using a multistep method.

Single-particle motion

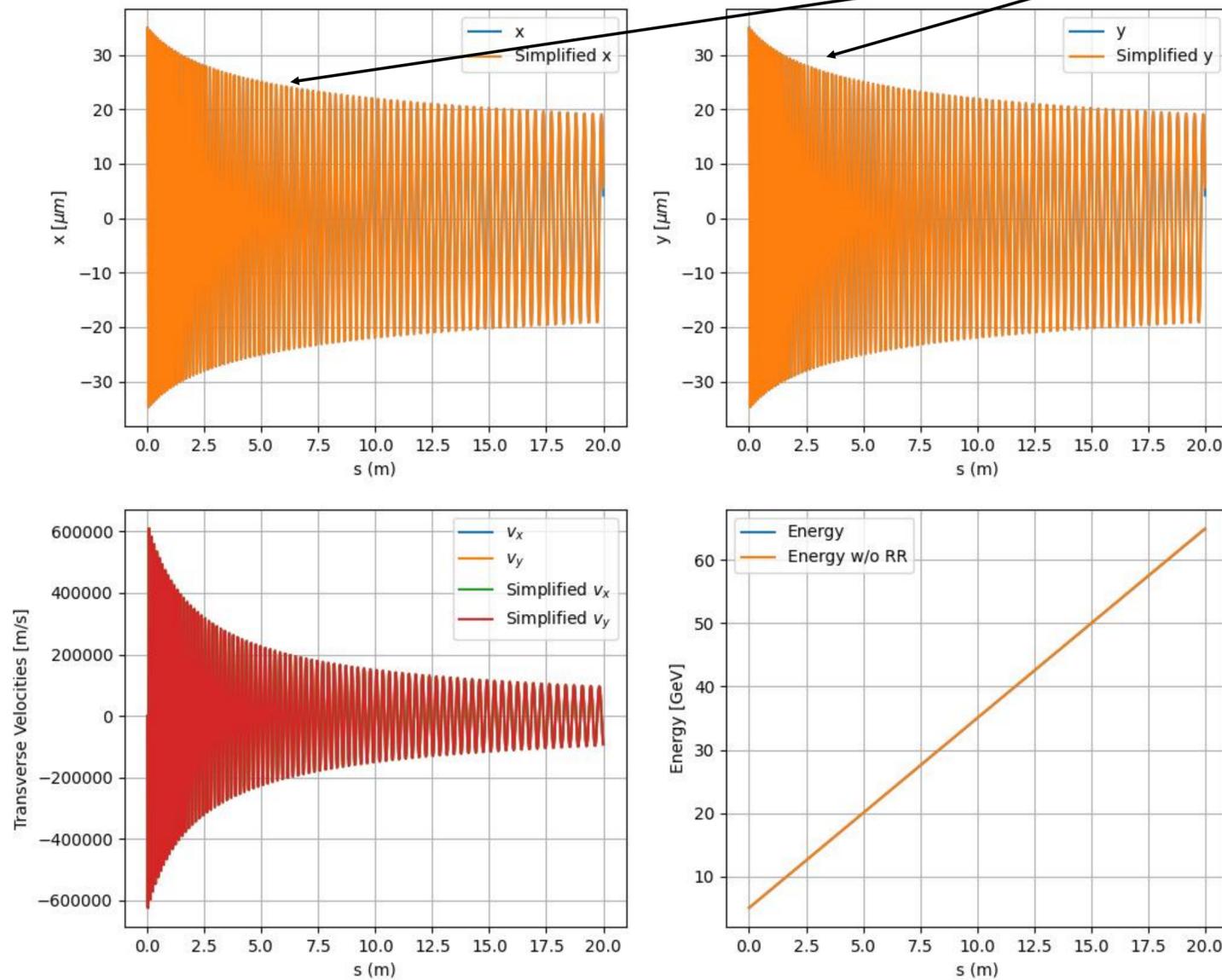
A diagonal oscillation with and without radiation

Evolution of Particle Parameters

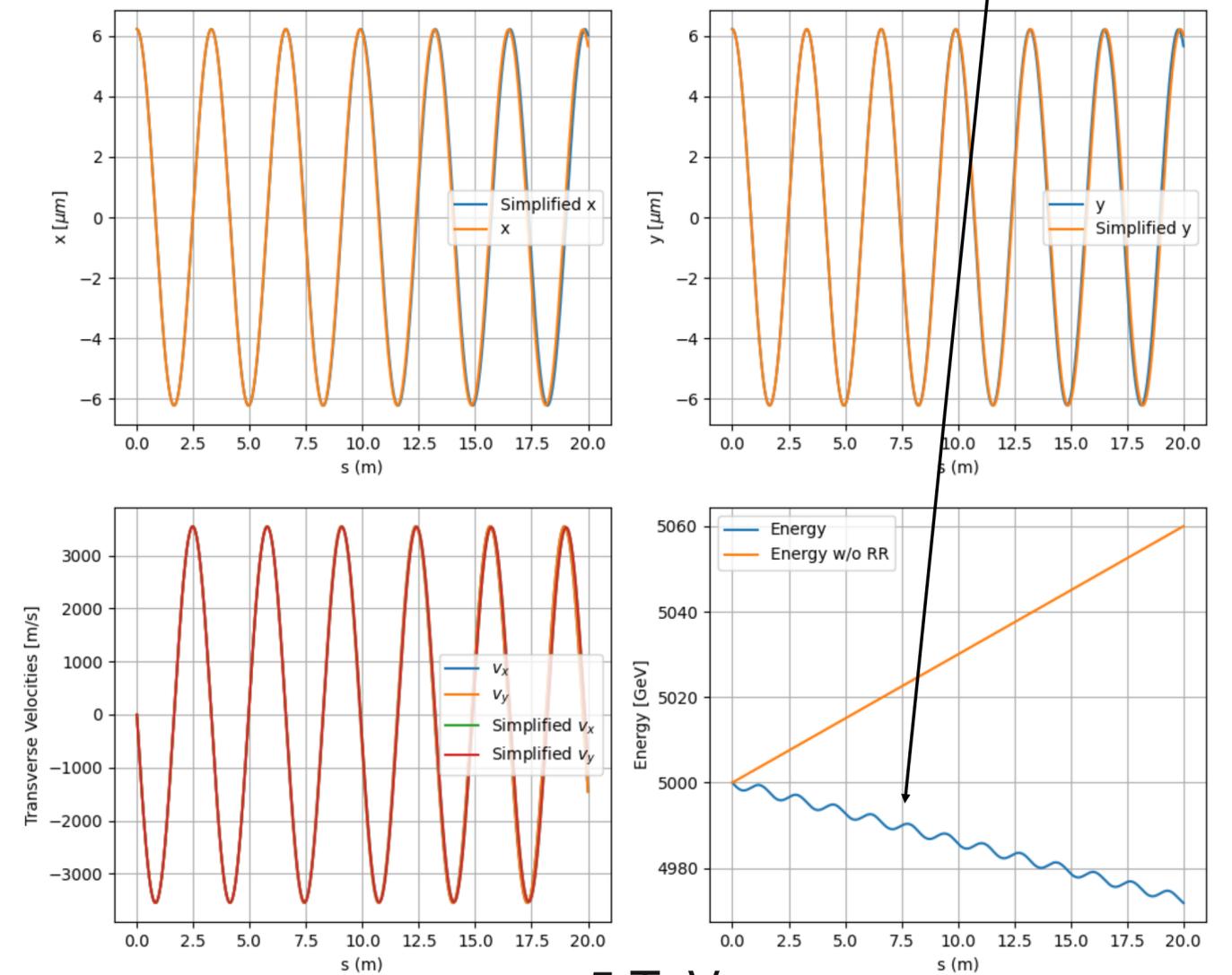
Adiabatic damping due to increase in energy

Evolution of Beam Parameters

Energy loss due to radiation reaction



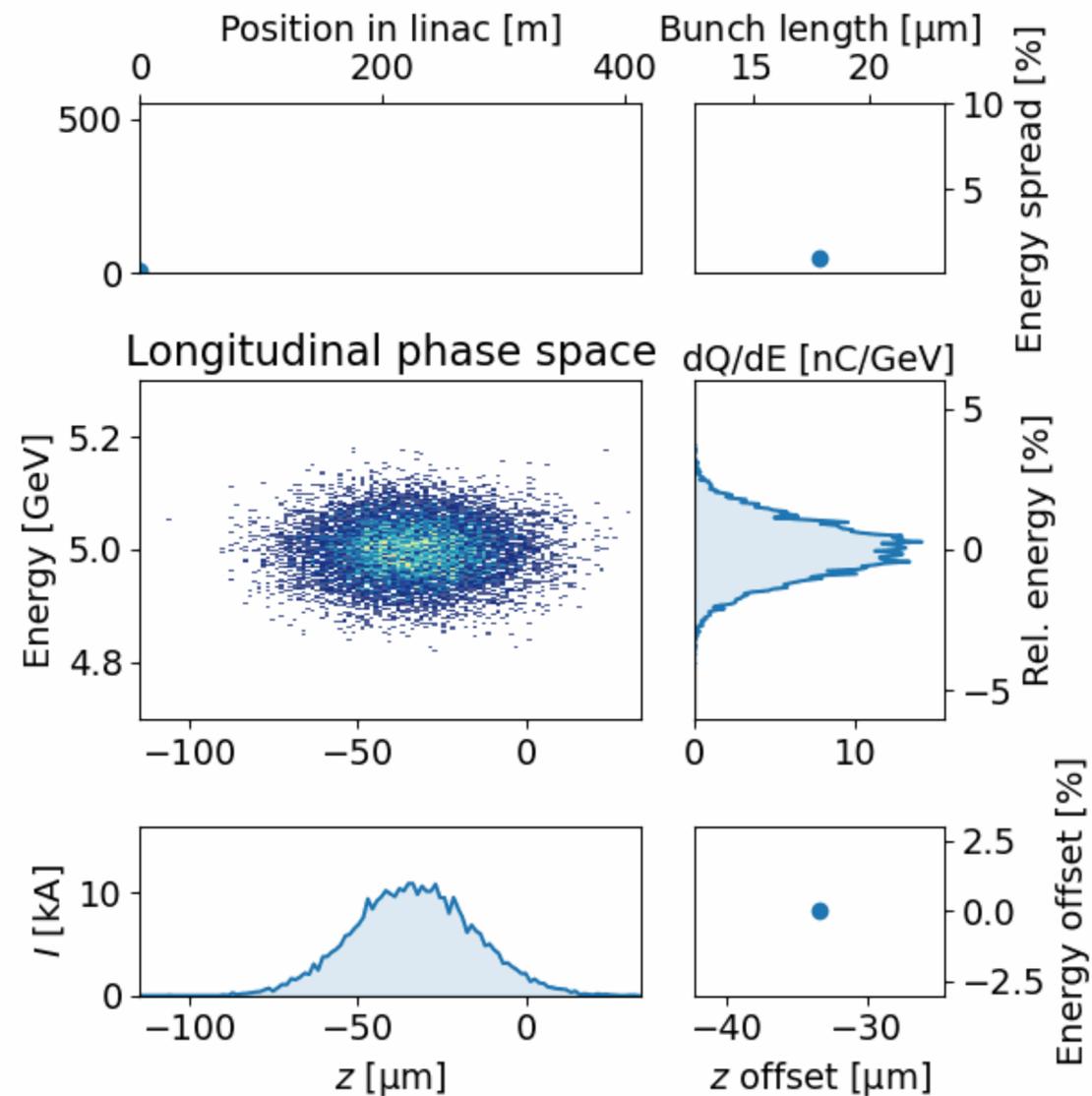
5 GeV



5 TeV

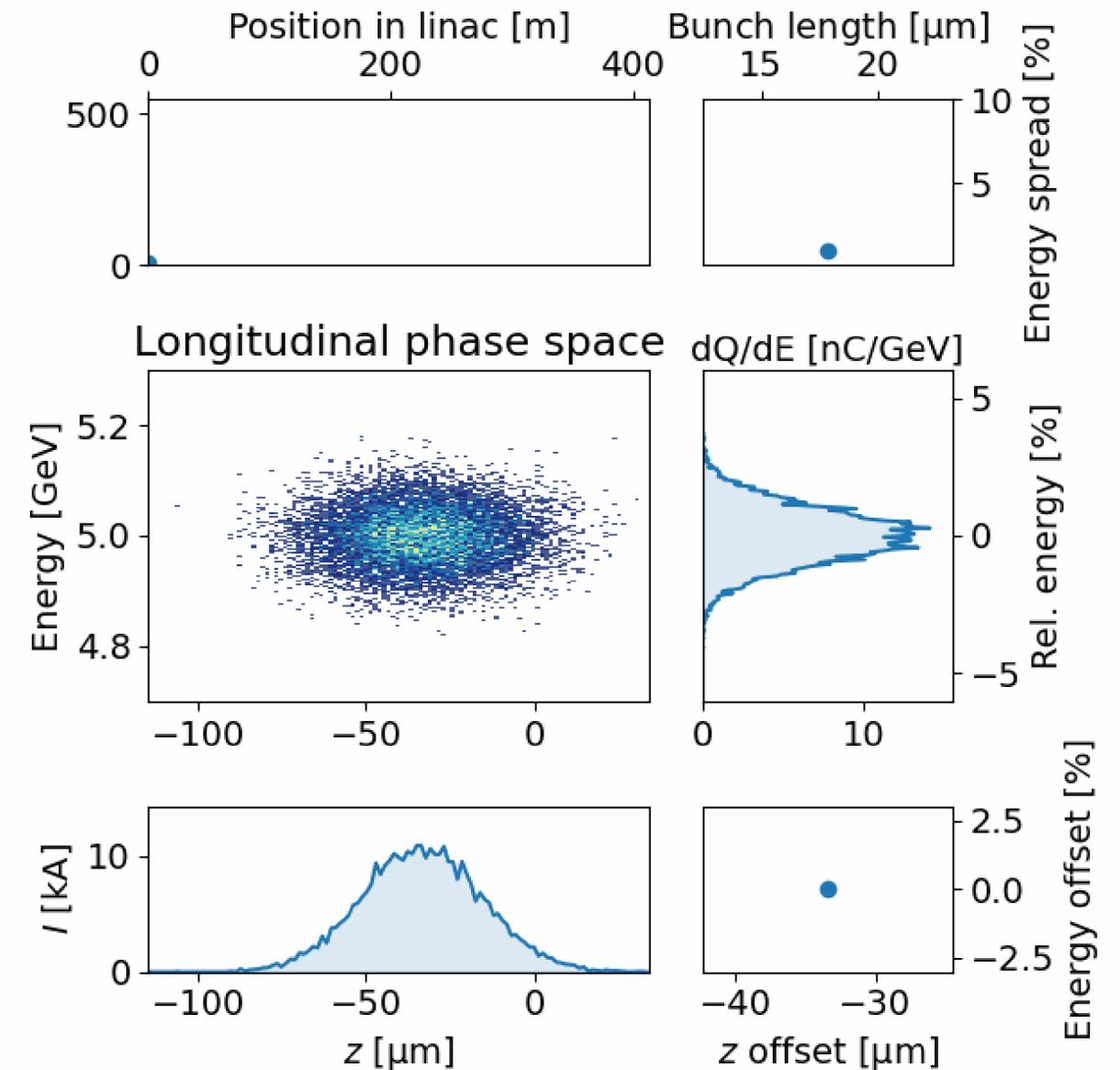
Beam energy loss

What does the effects of radiation reaction look like in the longitudinal phase space?



> The particles with large betatron amplitudes lose more energy.

> This causes a "Bleed" in energy, and an increase in energy spread.

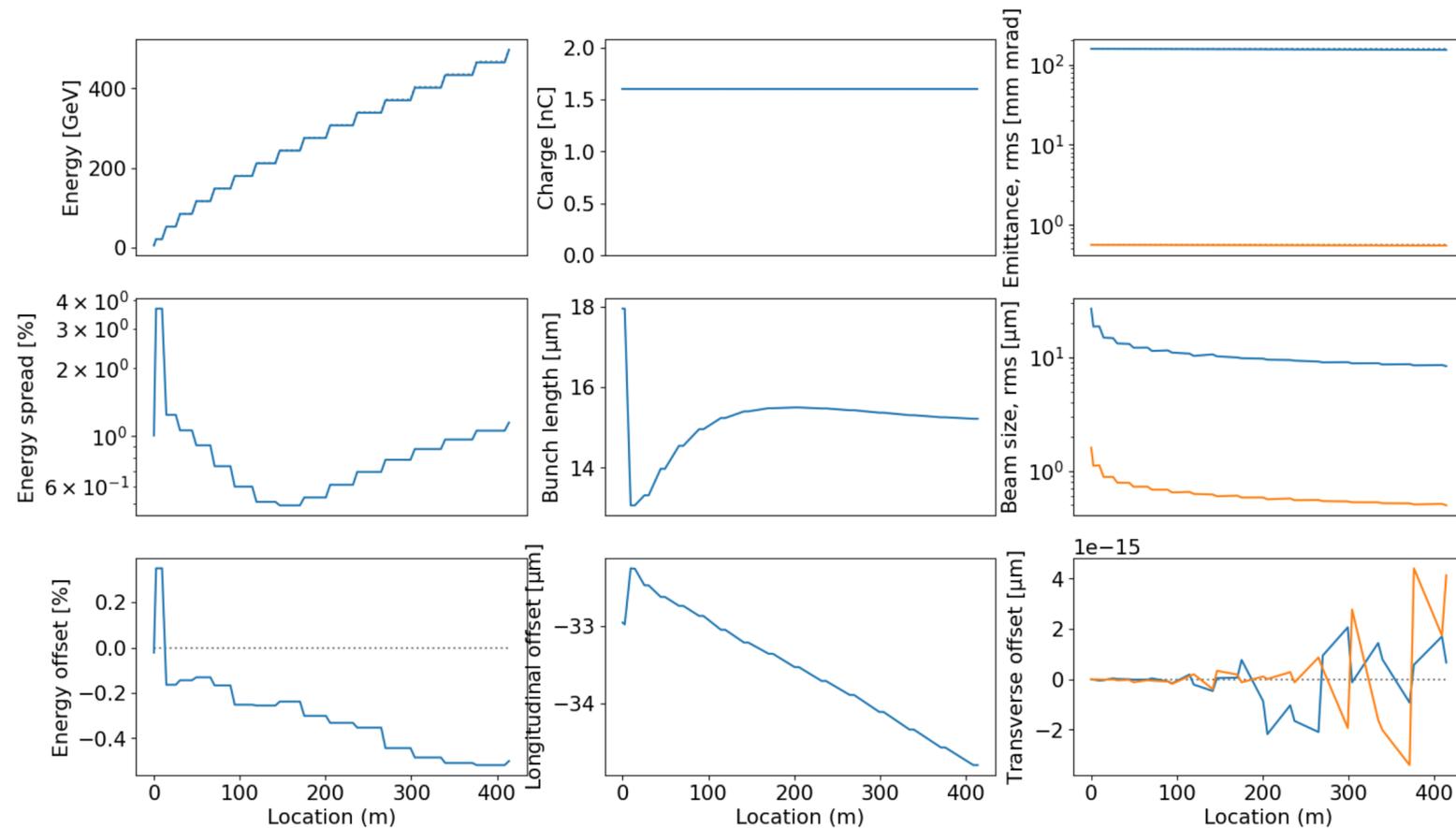


HALHF simulations

Tracking a beam through 16 stages

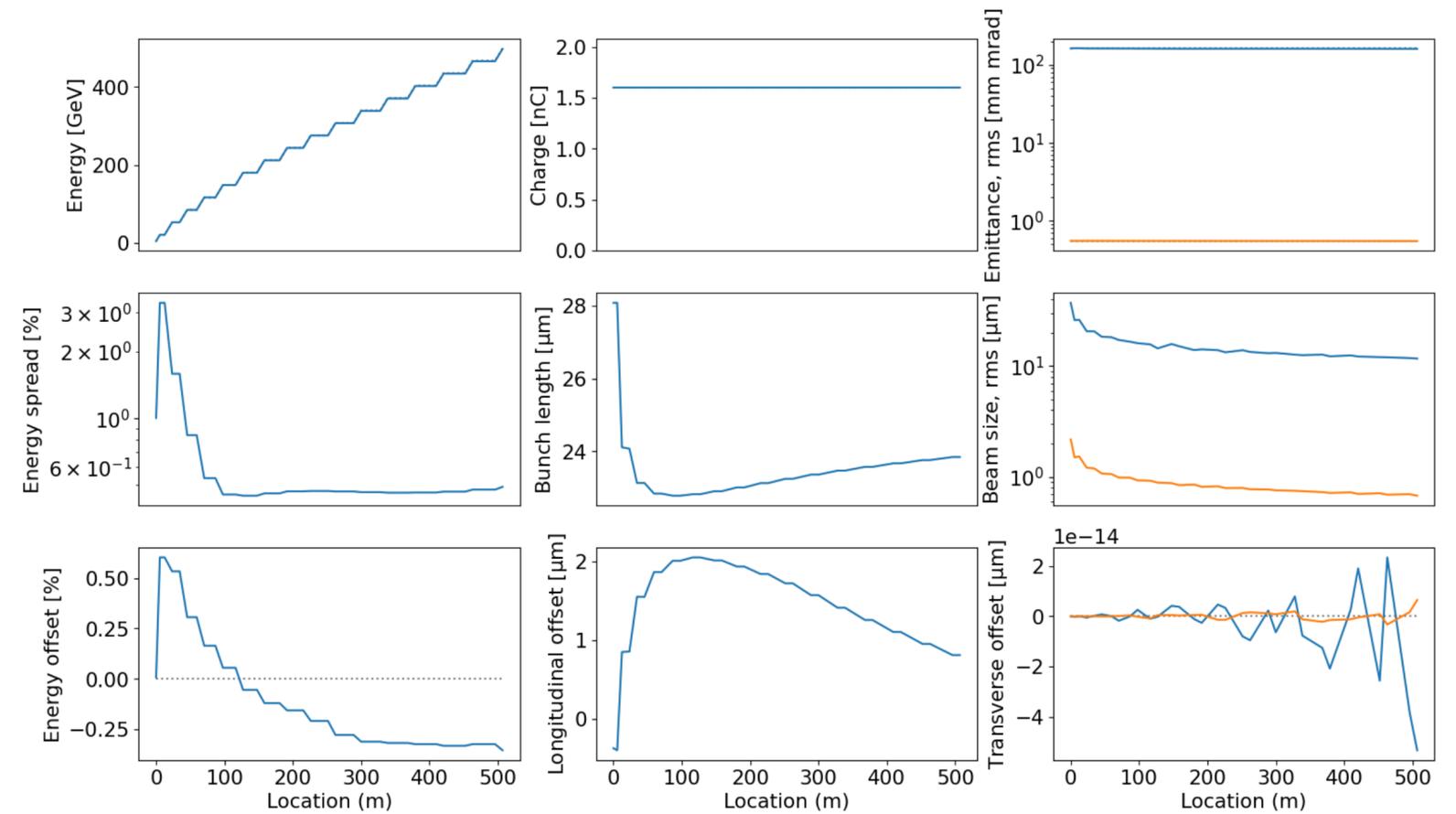
> Already at 500 GeV, we have had to lower the density

Evolution of Beam Parameters



$n_0 = 7e21 \text{ m}^{-3}$

Evolution of Beam Parameters



$n_0 = 2e21 \text{ m}^{-3}$

Goes as $n_0^{3/2}$

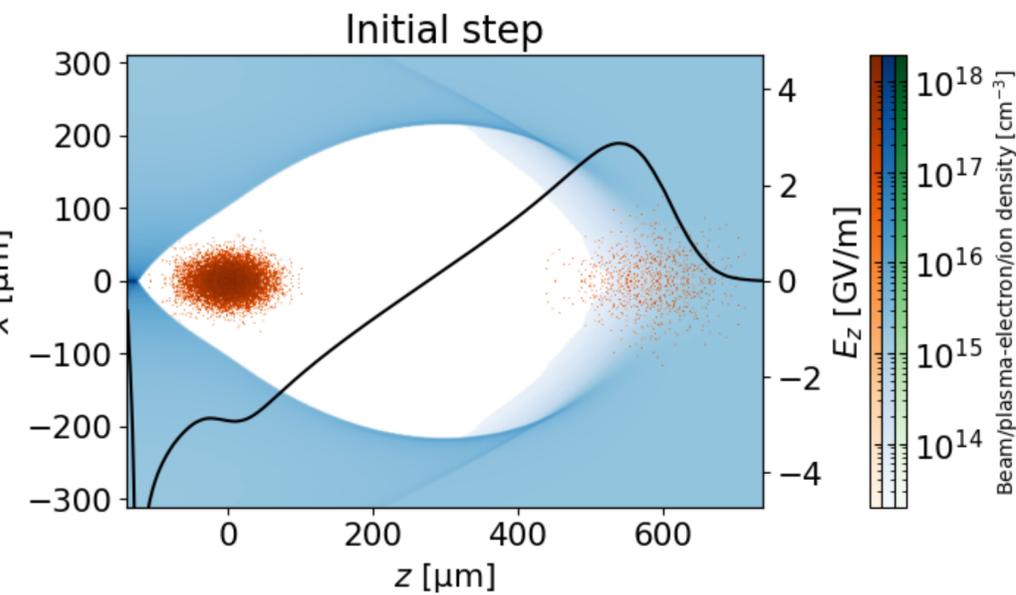
$$\sigma_{\bar{P}} = \frac{e^2 c}{24\pi\epsilon_0} (\sqrt{\gamma} k_p)^3 \sqrt{\epsilon_{nx}^2 + \epsilon_{ny}^2}$$

Plasma wakes

The effects of the radiation reaction + self-correction mechanism

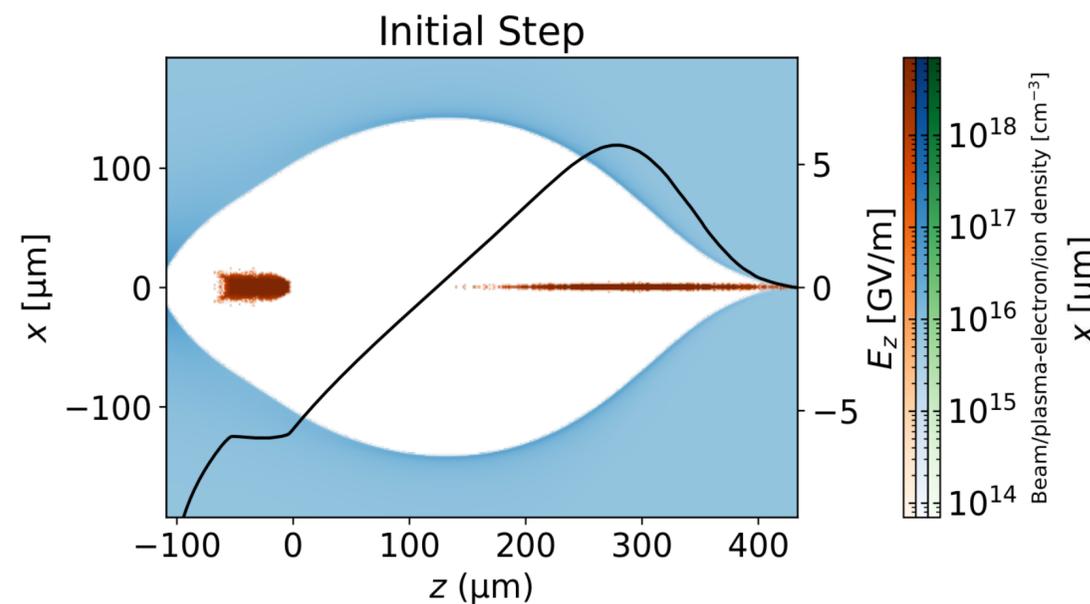
> Along with the self-correction mechanism, the beam shapes into a bullet.

1st stage



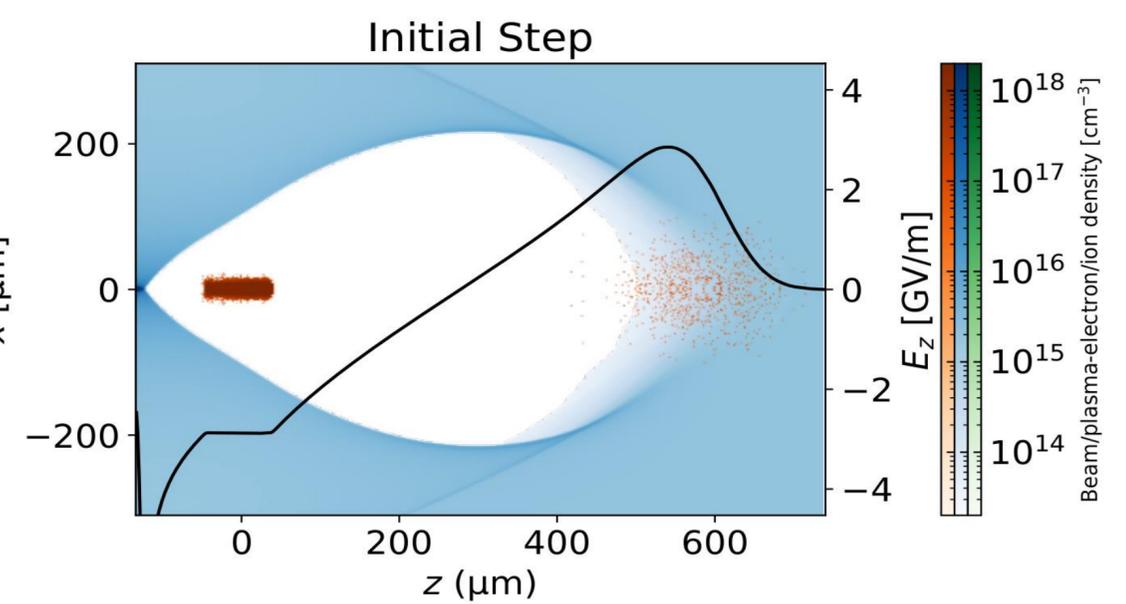
$n_0 = 2e21 \text{ m}^{-3}$

16th stage



$n_0 = 7e21 \text{ m}^{-3}$

16th stage



$n_0 = 2e21 \text{ m}^{-3}$

Goes as $n_0^{3/2}$

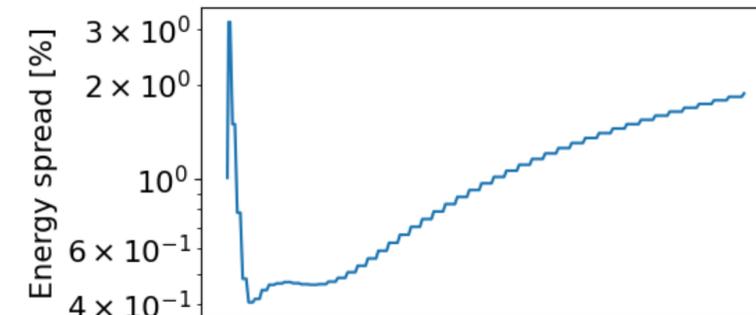
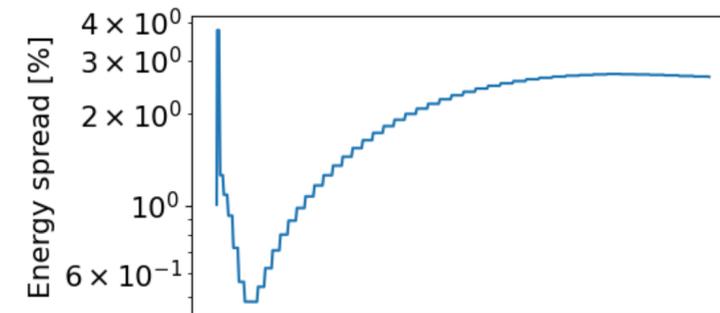
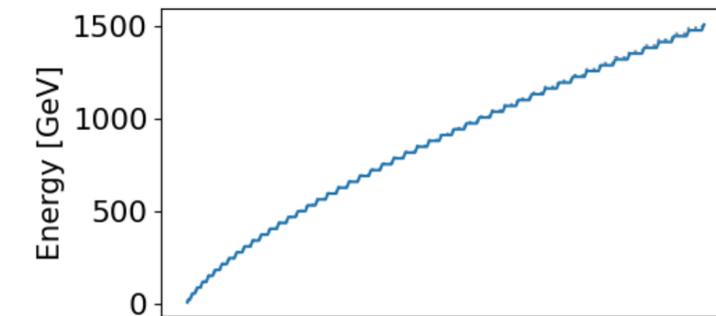
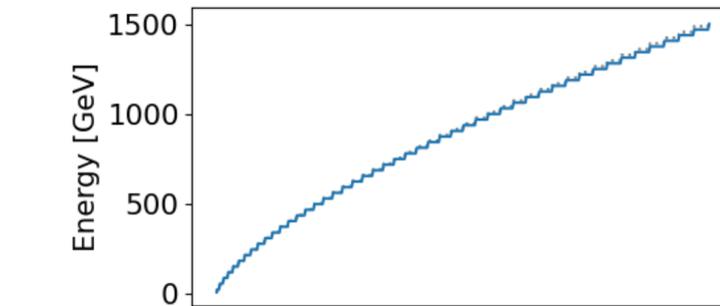
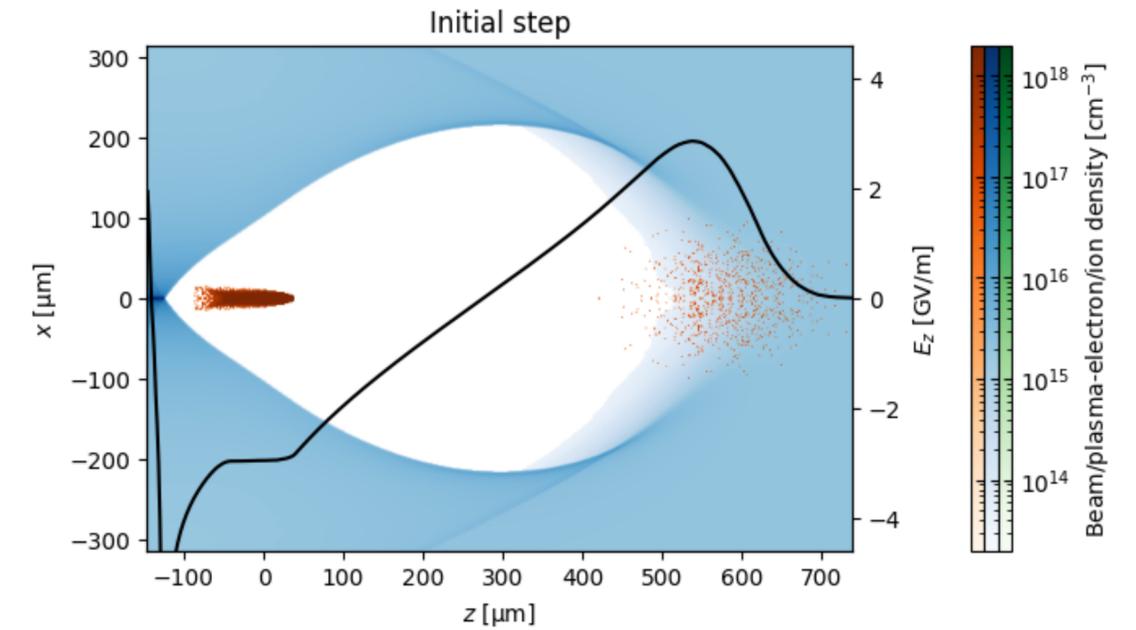
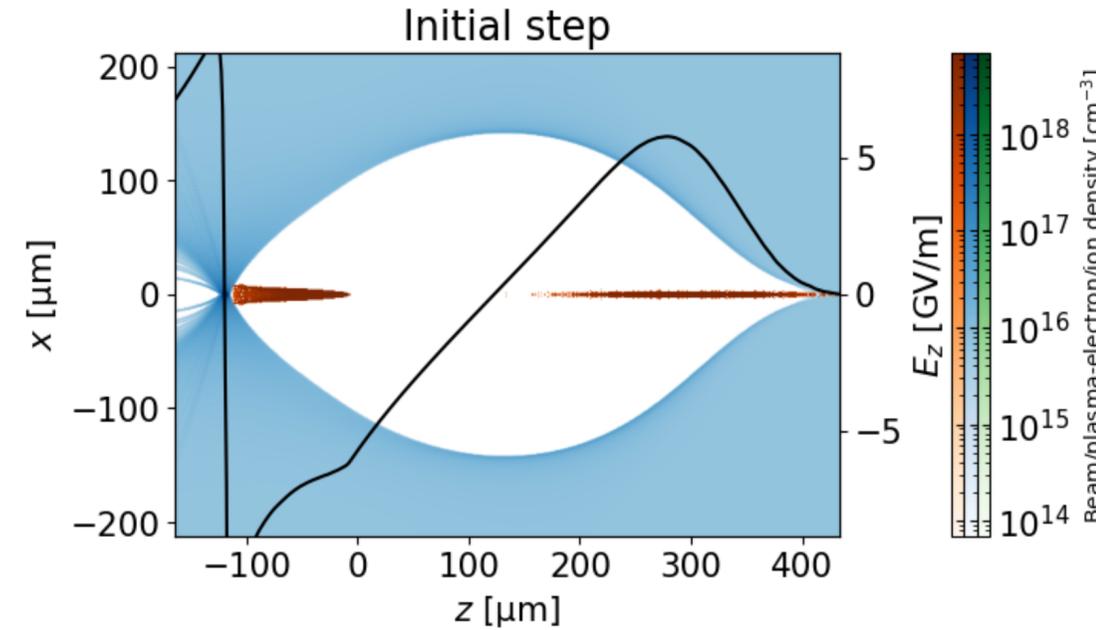
$$\sigma_{\bar{P}} = \frac{e^2 c}{24\pi\epsilon_0} (\sqrt{\gamma} k_p)^3 \sqrt{\epsilon_{nx}^2 + \epsilon_{ny}^2}$$

Reaching TeV energies

Can we take the self-correction further?

> The energy spread becomes too large.

> The beam reaches the end of the wake.

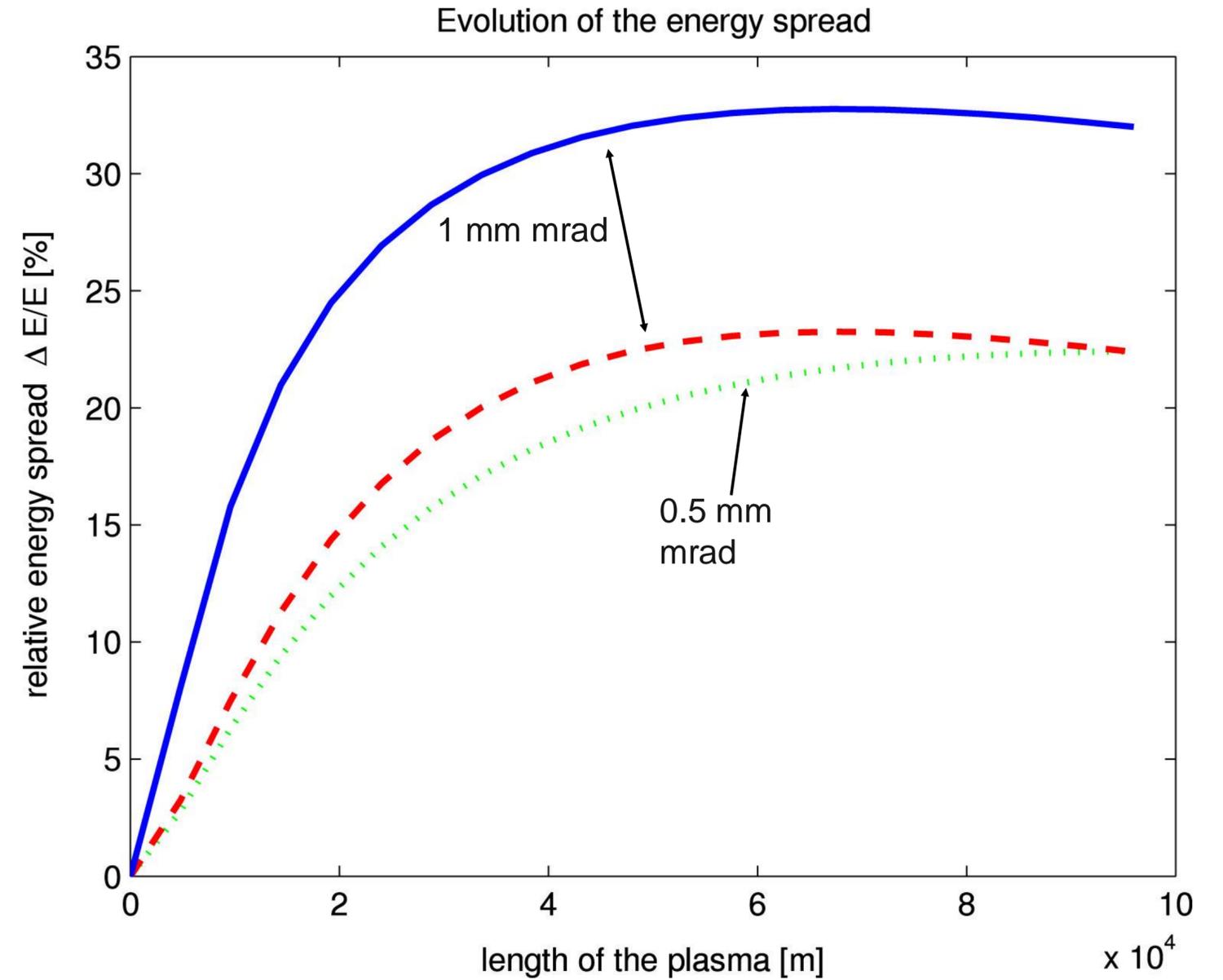
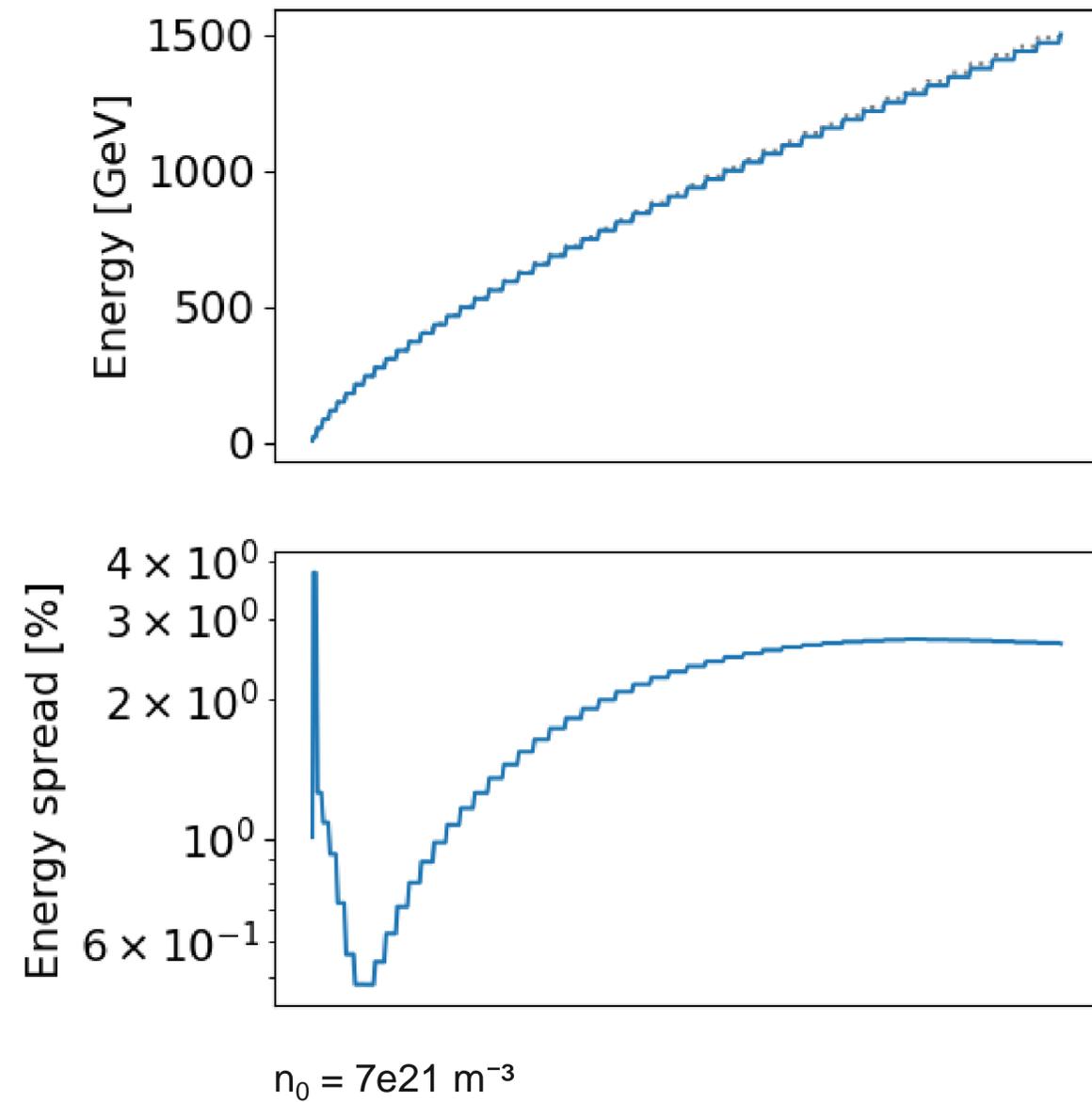


$n_0 = 7e21 \text{ m}^{-3}$

$n_0 = 2e21 \text{ m}^{-3}$

Maximum energy-spread

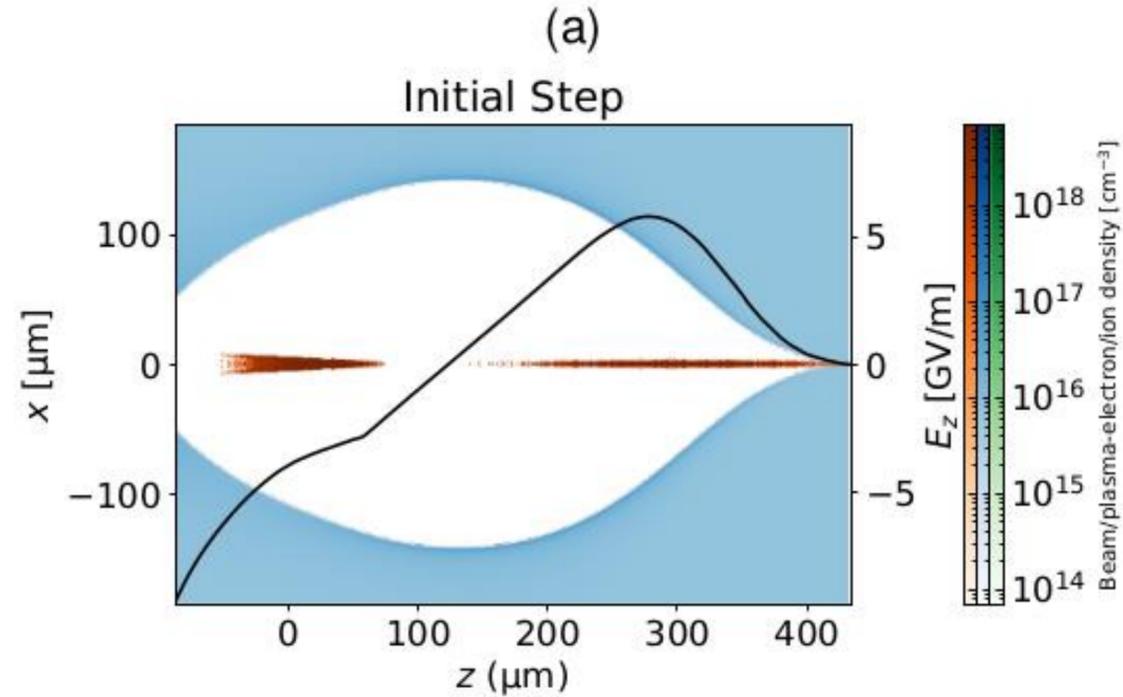
Shpakov et al. Nuclear Instruments and Methods in Physics Research Section A: Accelerators, Spectrometers, Detectors and Associated Equipment 909 (2018)



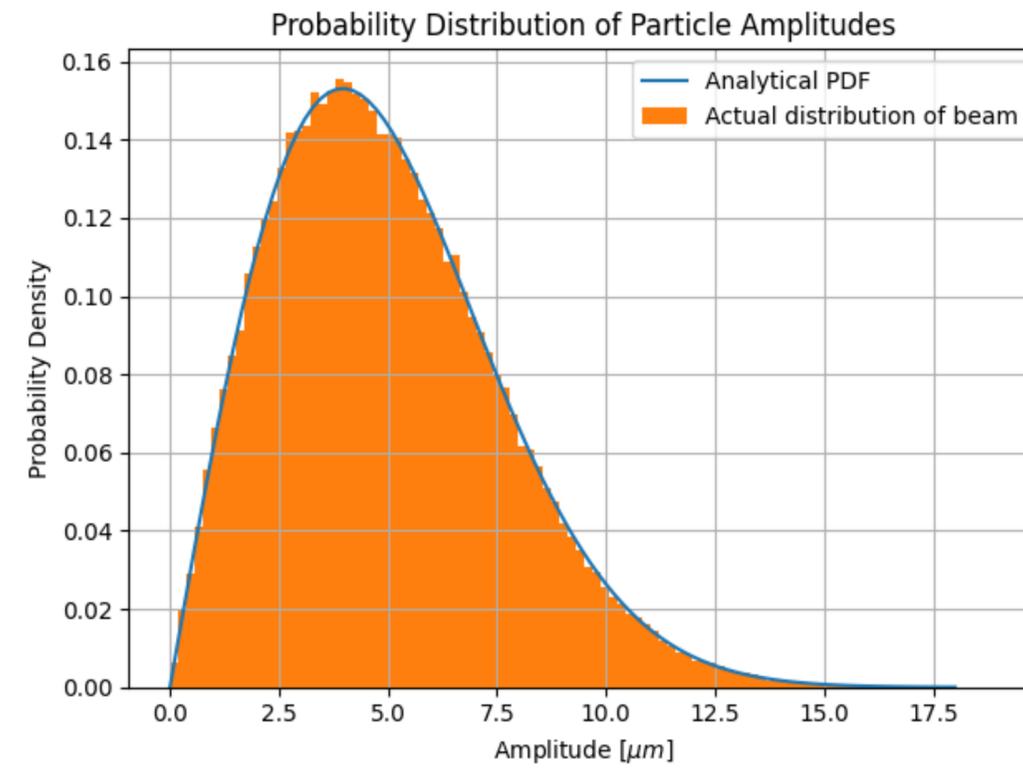
$n_0 = 1e22 \text{ m}^{-3}$ (blue/red), $n_0 = 5e21 \text{ m}^{-3}$ (green). no self-correction

Reaching TeV energies

- > There is now a strong correlation between longitudinal position and betatron amplitude!

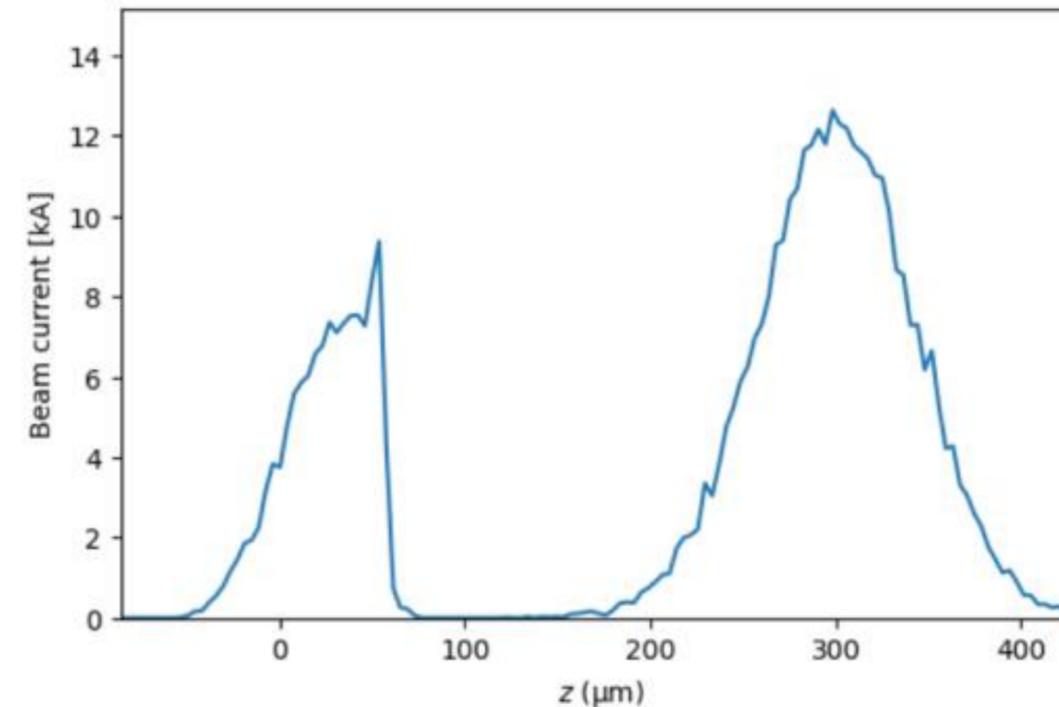


$$n_0 = 7e21 \text{ m}^{-3}$$



(b)

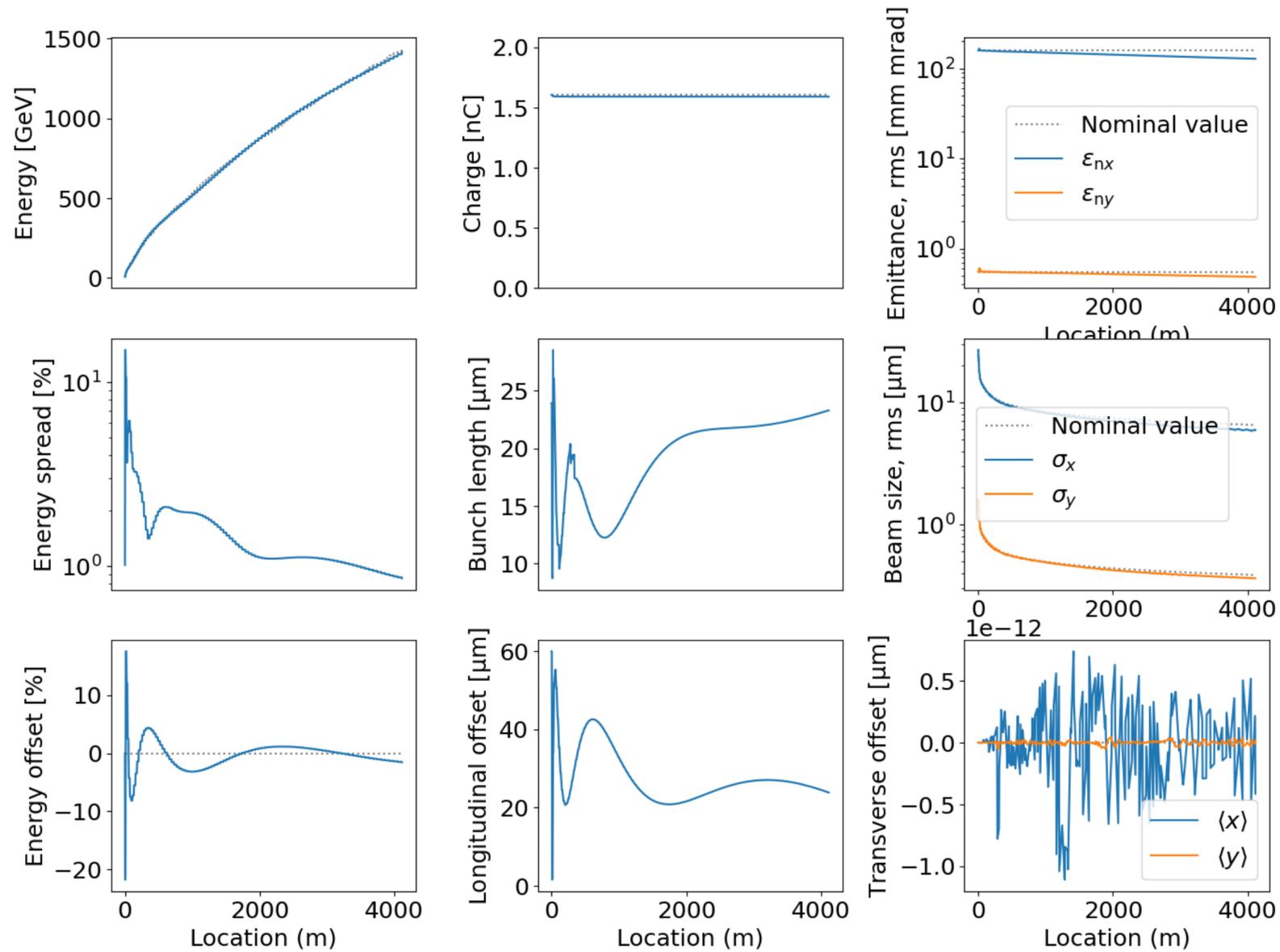
Current Profiles



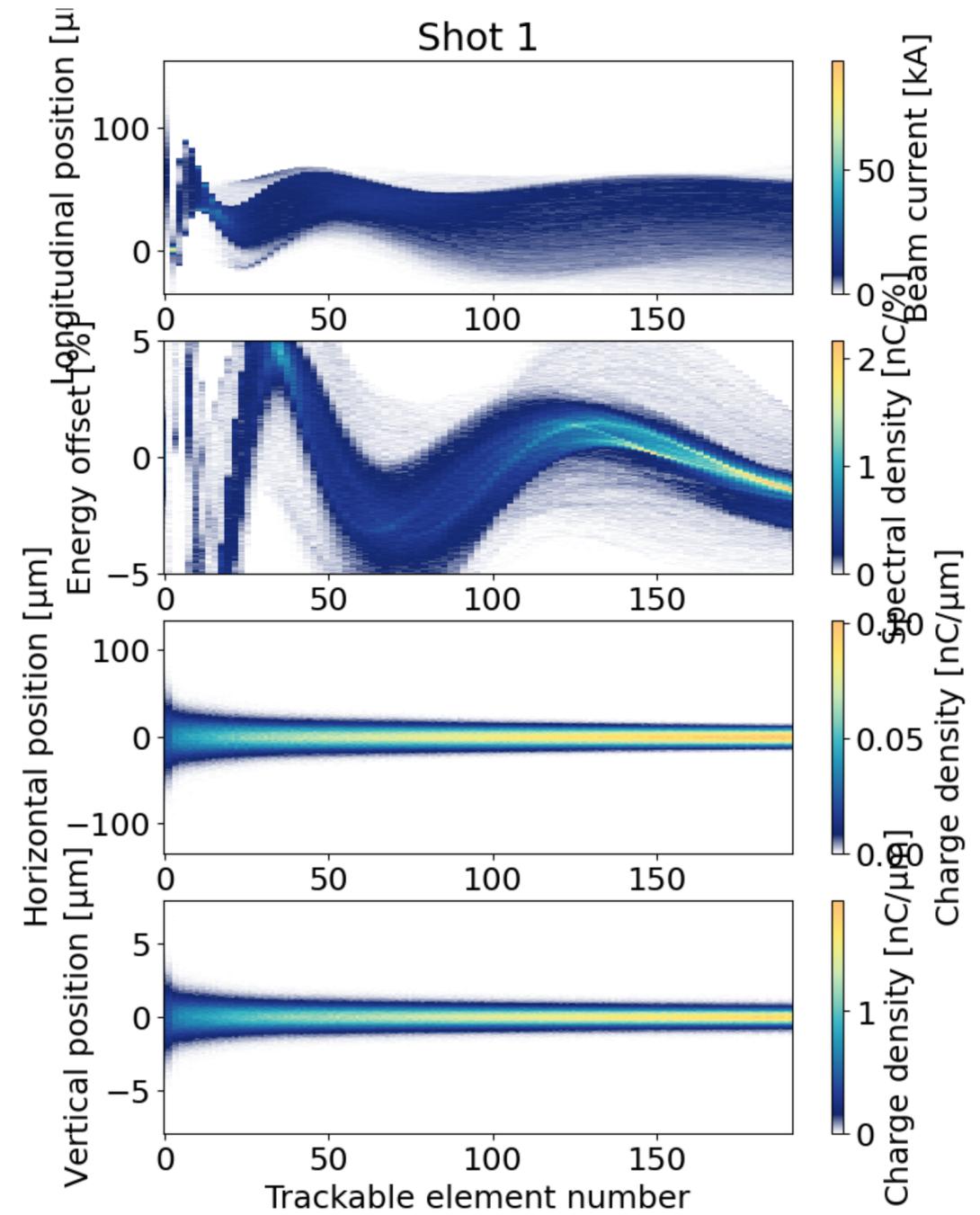
Reaching TeV energies

Adjusting the parameters

Evolution of Beam Parameters



$$n_0 = 7e21 \text{ m}^{-3}$$



Reaching TeV energies

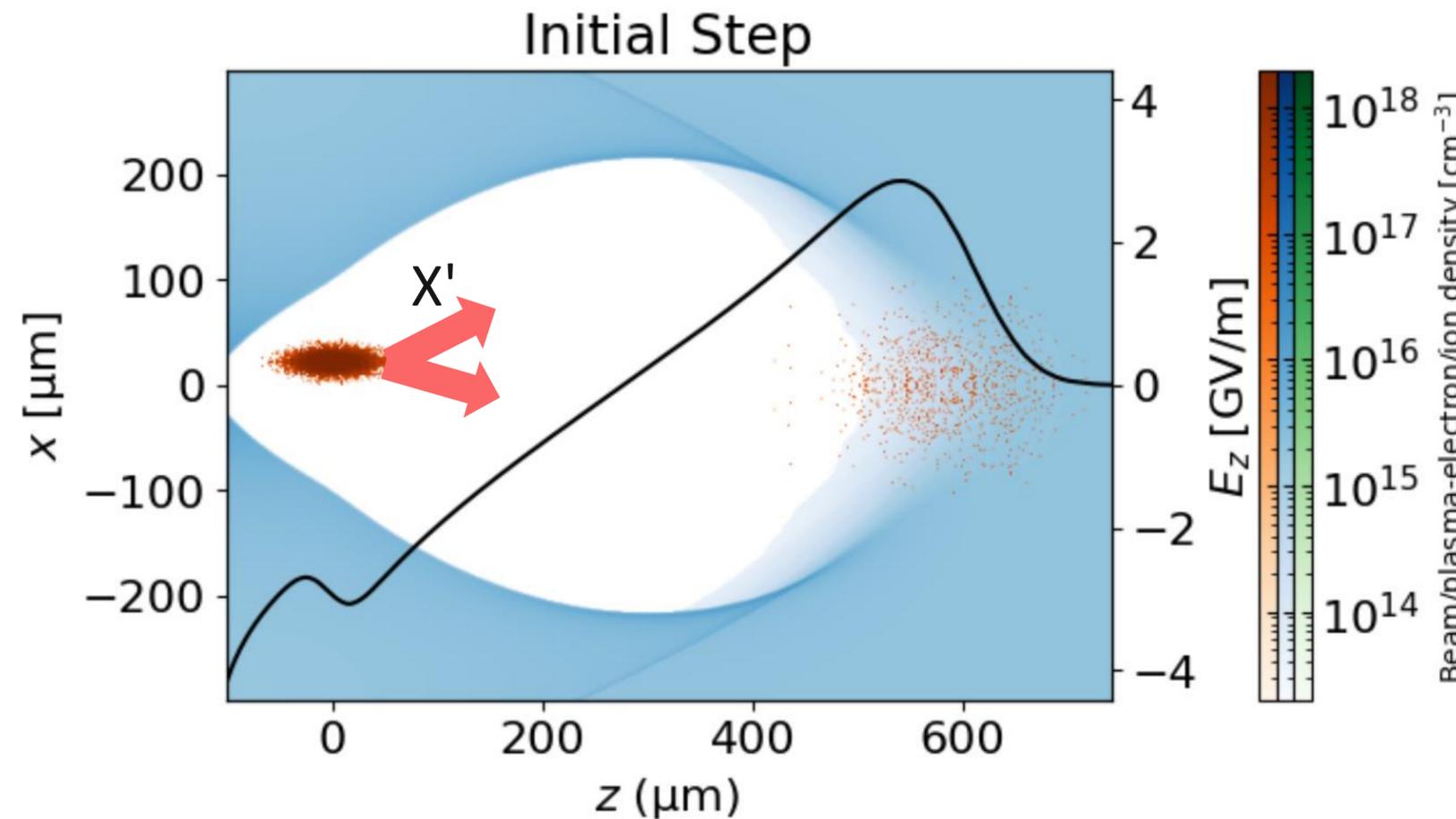
The self-correction can't be applied forever. The beam size will increase in the magnetic chicanes.

What About Jitter?

Radiation due to small offsets in staging

- > When the beam enters the plasma, it will have a small random offset (jitter). In x, this means the new energy loss will be

$$\langle \bar{P} \rangle_x = \Lambda \left(1 - \frac{1}{2} \frac{\Delta x^2}{\sigma_x^2} \right) \left(2\sigma_x^2 + 2\Delta x^2 + \frac{15}{16} \frac{\Delta x^4}{\sigma_x^2} + \frac{5}{8} \frac{\Delta x^6}{\sigma_x^4} \right)$$

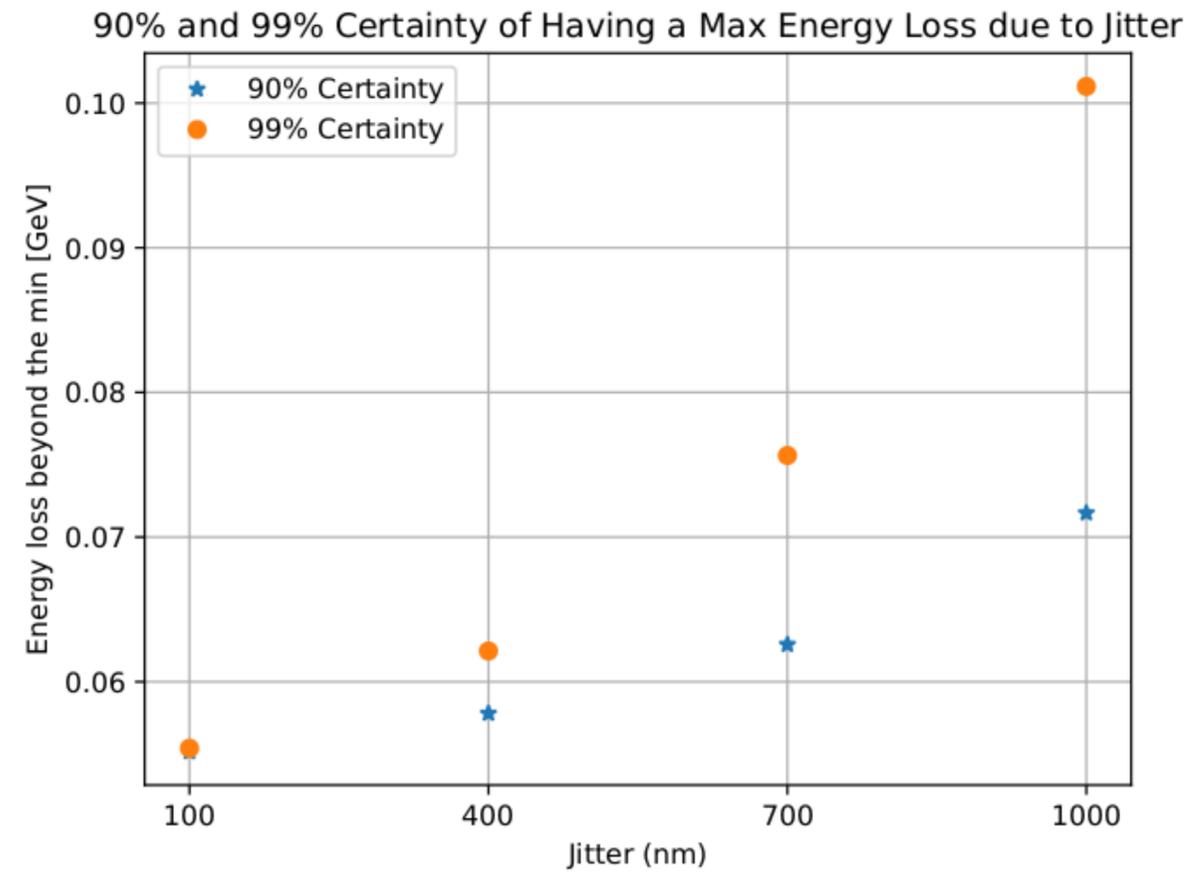


$$\Lambda = \frac{e^2 c}{12\pi\epsilon_0} (\gamma k_\beta)^4$$

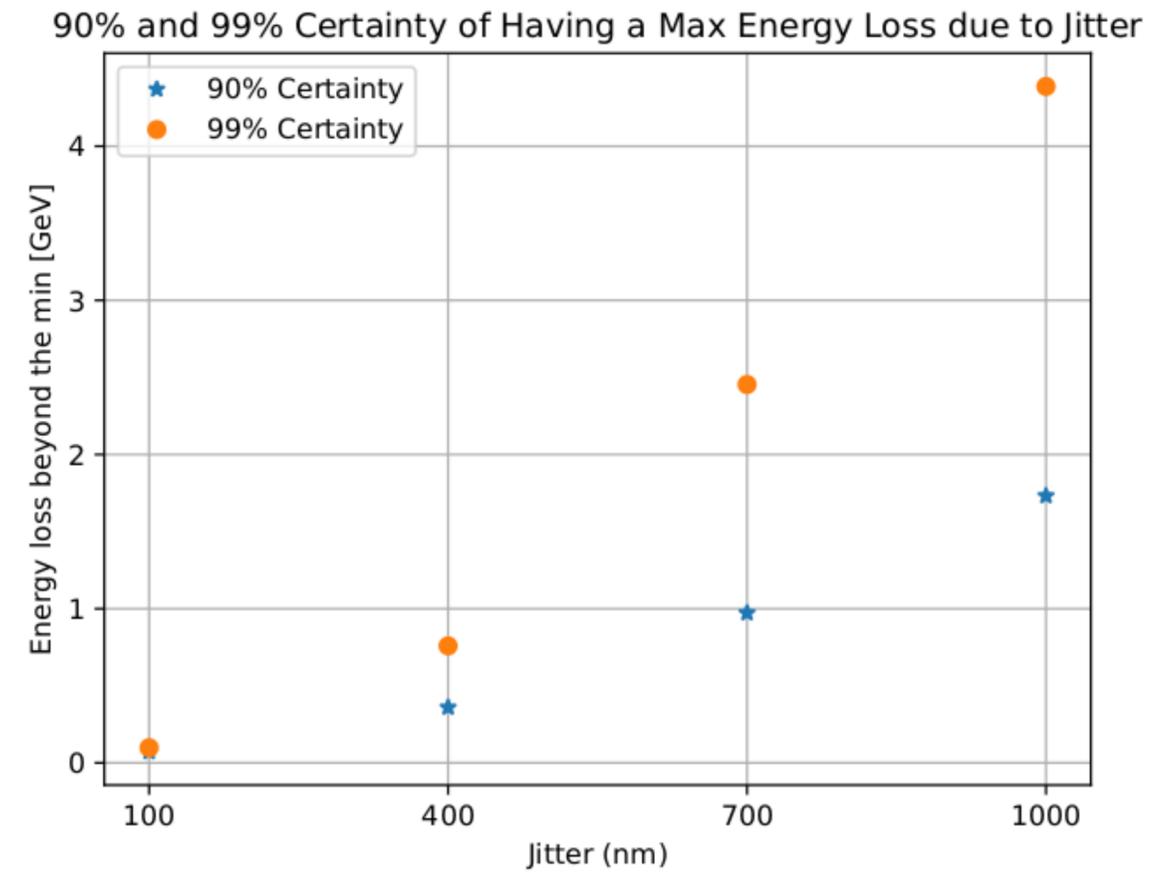
What about jitter?

Extra energy loss (beyond what the beam would lose anyway) due to jitter

500 GeV



5 TeV



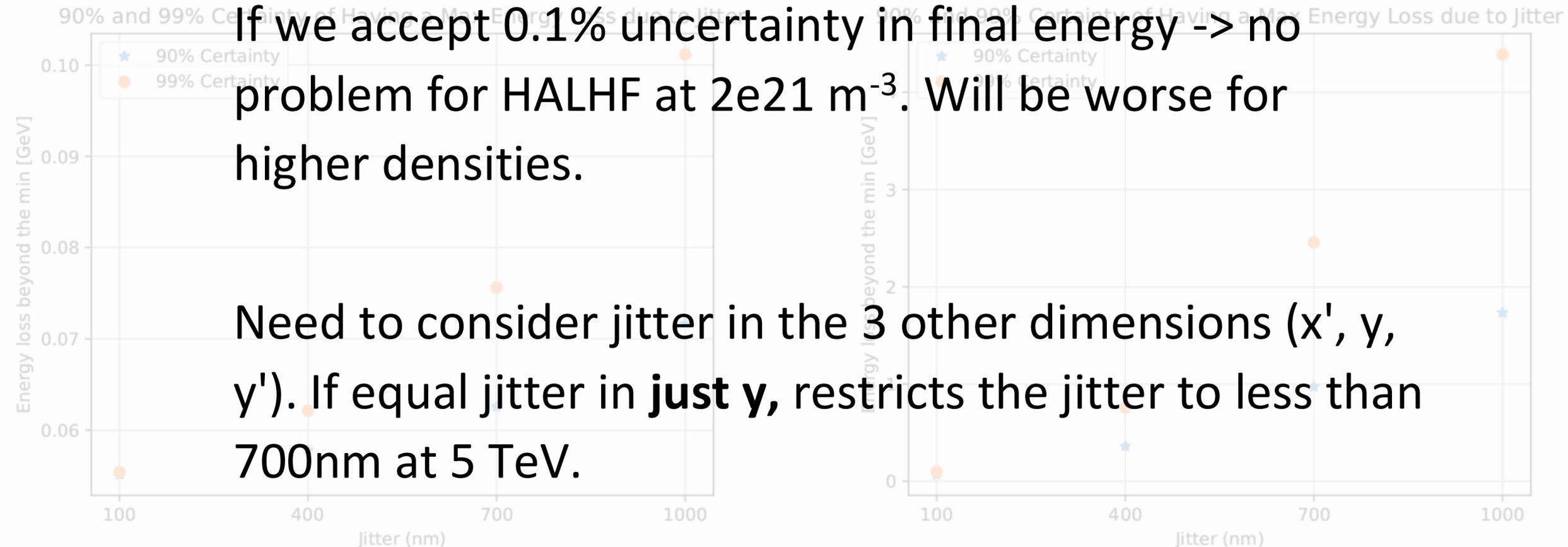
> Maximum energy loss with 90/99 % certainty.

What About Jitter?

Extra energy loss (beyond what the beam would lose anyway) due to jitter

500 GeV

5 TeV



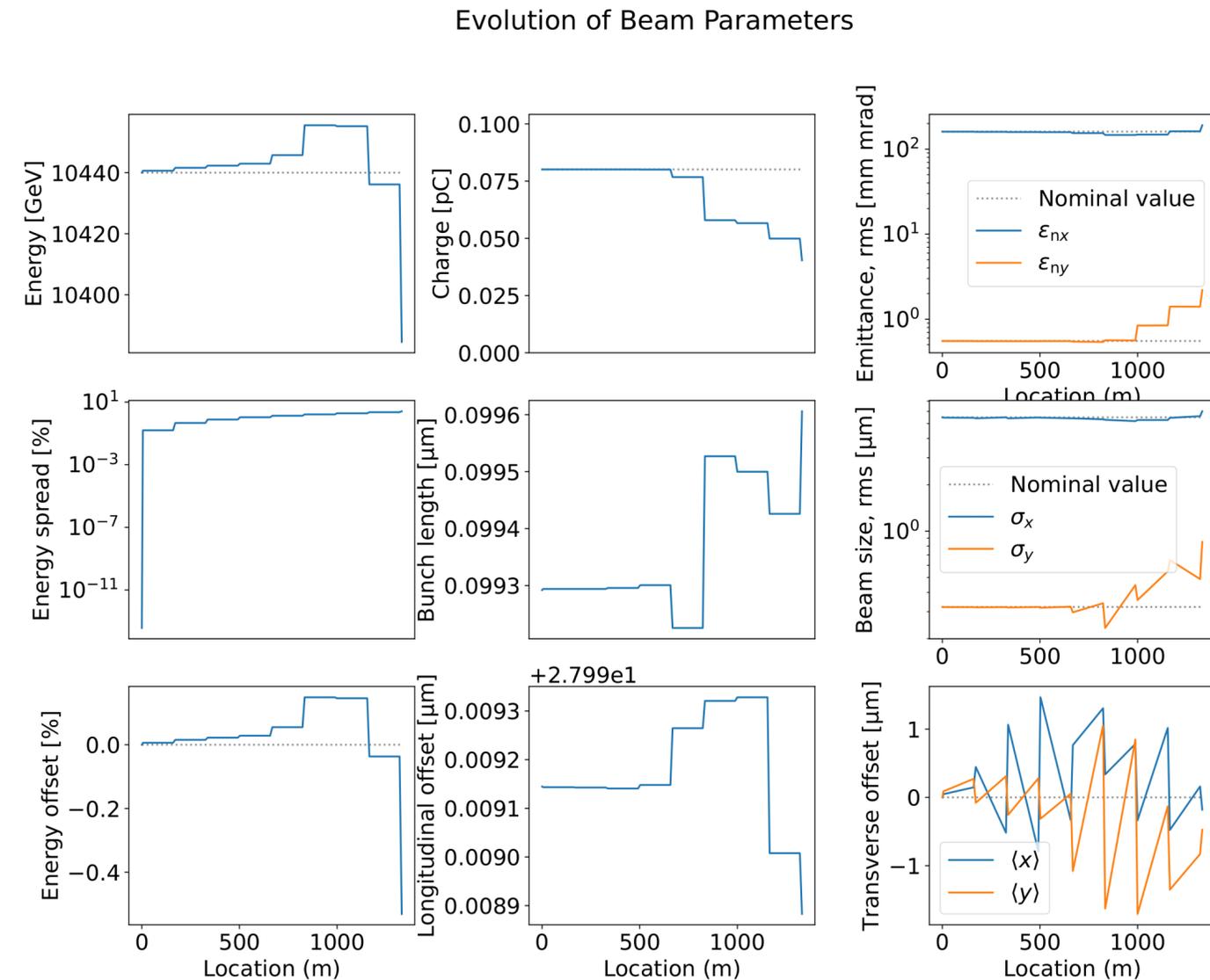
If we accept 0.1% uncertainty in final energy -> no problem for HALHF at $2e21 \text{ m}^{-3}$. Will be worse for higher densities.

Need to consider jitter in the 3 other dimensions (x' , y , y'). If equal jitter in **just y**, restricts the jitter to less than 700nm at 5 TeV.

What about jitter?

Maximum achievable energy

- > With staging, there is a maximum achievable energy.
- > This maximum is probabilistic, and depends on the offset.



Jitter = 250 nm in x and y.

Summary

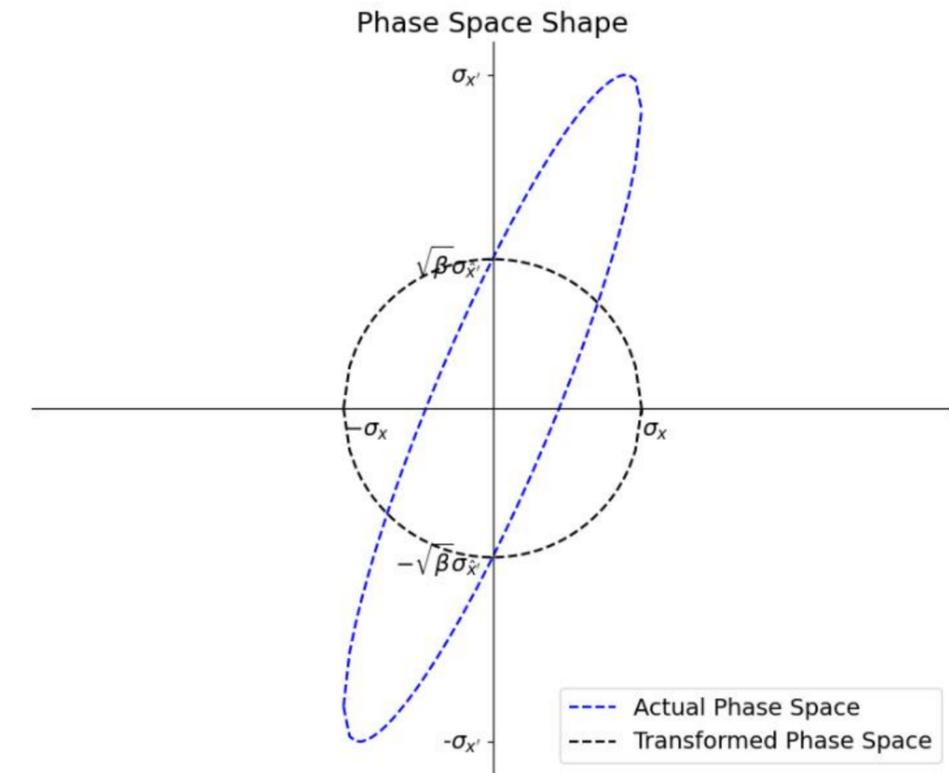
- > The radiation reaction causes energy loss and energy spread. The self-correction mechanism helps correct for the energy spread, but will it work in practice at 5 TeV?
- > These simulations show that TeV acceleration is possible with regards to energy spread. Still need to examine the radiation reaction in the correcting chicanes.
- > Jitter contributes to energy loss, and random energy offsets. Seems not to be a problem for HALHF but seems to impose restrictions at 5 TeV.

Extra

Matching and Betatron Amplitude

> Normalized phase space (from ellipse eq.)

$$x^2 + \beta \hat{x}^2 = \varepsilon \beta \quad \text{where} \quad \hat{x}' = \alpha x + \sqrt{\beta} x'$$



> We also have $x^2 + x'^2 = r_x^2 (\cos^2(k_\beta ct + \phi) + k_\beta^2 \sin^2(k_\beta ct + \phi))$

Which means $x^2 + \bar{x}'^2 = r_x^2$ where $\bar{x}' = \frac{1}{k_\beta} x'$

Caption

Extra

Matching and Betatron Amplitude

> For a matched beam; $\alpha=0$, and the single particle trajectories must follow the shape of the phase space.

> I.e $x^2 + \bar{x}'^2 = r_x^2$ must match the shape of $x^2 + \beta \hat{x}^2 = \varepsilon \beta$, or the energy spread will cause emittance increase.

> We therefore know that
$$r_x = \sqrt{x^2 + (x' \beta_m)^2}$$

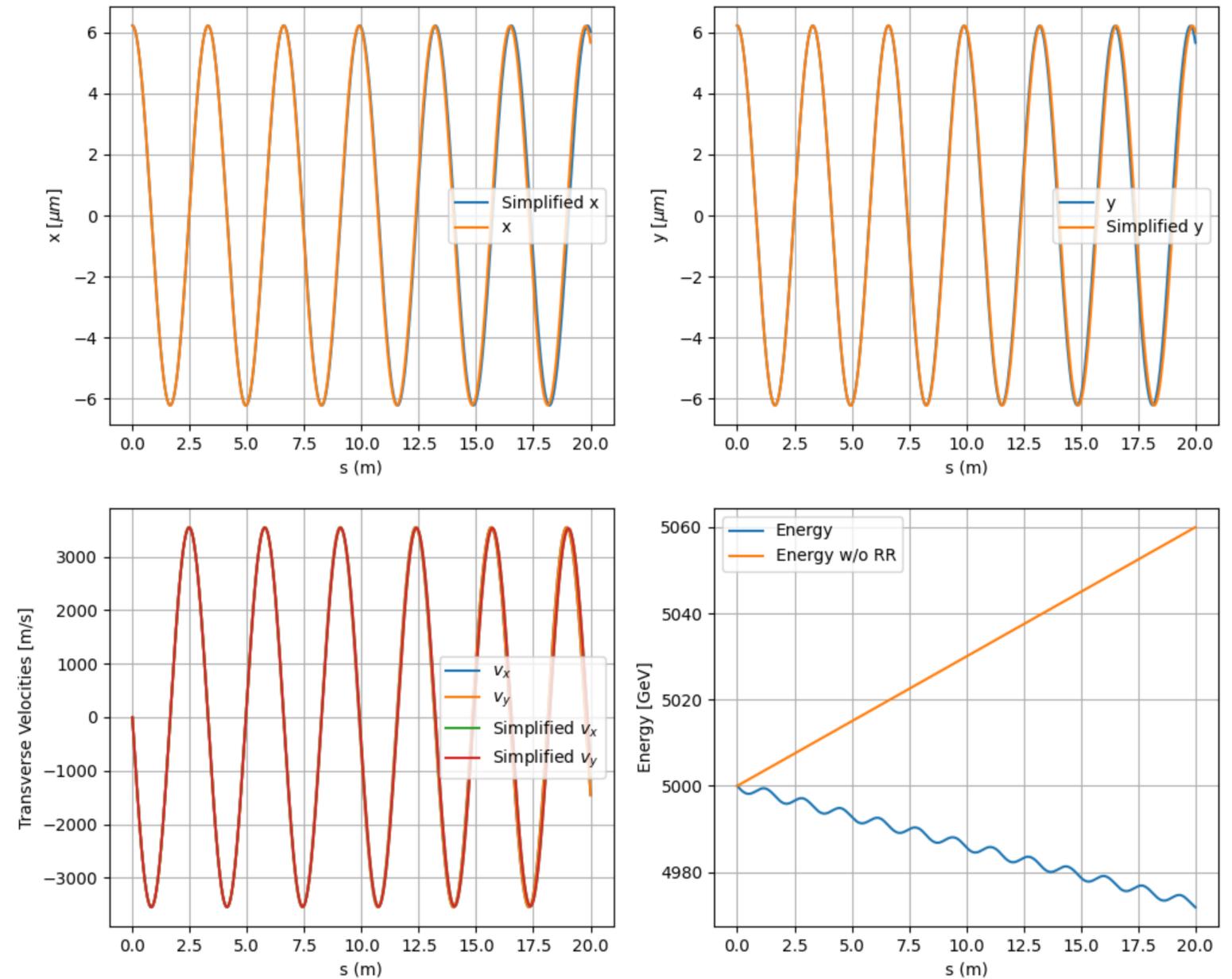
Caption

Extra

Betatron Motion at 5 TeV

- > Even at 5 TeV, the motion w/wo radiation reaction is almost the same

Evolution of Beam Parameters



Extra

Jitter energy loss at 7.5 TeV

> The plot that led me to believe my initial expression for the jitter energy loss was correct

