

On the influence of electrostatic barriers that would modify brighter-fatter corrections for PSF estimations Andy Rasmussen & Alex Broughton – ISPA24C41



- We use a generic Si drift simulator (<u>https://github.com/arasmssn/bi\_ccd\_pixpart</u>), operated under cold carrier approximation & validated against all available model-constraining observables to compute pixel boundary response to fixed-pattern & dynamic charge configurations within the pixel. *Linked list* implementation for E-field source management is appropriate for this problem: 10-20 pixel electrostatic field range & trajectory starting point bisection algorithm tolerance (<10<sup>-5</sup> pixel). (cf. A.Rasmussen 2014 *JINST* 9 C04027).
- This utilizes a non-unique solution to Poisson's equation within the photosensitive volume (depleted of carriers) and uses superposition of electric field sources that are solutions to Gauss's law, method of images are used to approximate planes of symmetry (and equipotential surfaces) e.g. the conductive backside window and polysilicon gate structure in these CCDs.
- Derived solutions for the time being are specific to one flavor of *LSST* focal plane sensor model (Teledyne/e2v-CCD250) operated under design conditions (-70V BSS) & integrating clock swing configuration (cf. Guyonnet et al. 2015).
- Many, many other datasets exist for the 205 devices/3216 amp channels distributed across sensor models & optimized operating conditions (cf. Y. Utsumi, contrib.42). There's a lot to learn from a single operating case, single signal level.



### Reconciling observables for carrier drift to validate E-field contributions (figs from R14)









Previous slide demonstrated the intrinsic "fingerprint" E-field basis vectors for this sensor



The implementation of this model is represented as a sketchup cartoon (image charge configurations not shown). cf. A.Rasmussen 2015 *JINST* **10** C05028.

- 2x2 pixel region near the channel shown.
- Serial addresses vary along the red coordinate.
- Polysilicon gates form equipotentials extending along this axis.
- Parallel addresses vary along the green coordinate.
- Channel stop barriers form (periodic/infinite) p+ negative charge distributions extending along this axis.
- Non-unique solution to Poisson equation represented by E<sub>BD</sub>.
- Carriers drift toward the channel (green ovals) from the backside window and partition into pixels. For marginal cases near boundaries, saddle condition is approached. Saddle condition is defined as:

 $\vec{\nabla}\phi = 0$  and  $\partial^2\phi\partial z^2 < 0$  and  $(\partial^2\phi/\partial x^2 > 0 \text{ or } \partial^2\phi/\partial y^2 > 0)$ 





#### What does this have to do with the brighter-fatter effect?





#### How are the simulations evaluated and compared to observables?









NB#0: These area distortions are between x400 larger than what is seen in flat fields! <sqrt(var)> ~ 240e<sup>-</sup>.

NB#1: area distortions ( $\Delta a_{01}$  etc.) are observable in flat field correlations, while the corrections ( $M_{0100}$  etc.) are not observable. Recovering them was a primary focus of P.Antilogus et al. 2014 JINST 9 C03048, §5.

Pixel boundary formalism (right hand figures) permits analysis of overlaps as corrections e.g. sparse matrix with 4 indices indicating source & destination pairs: From: (i,j) To: (i',j') and expressed as M<sub>iii</sub>. In this extremely distorted case\*\*,

3×10-

$$M_{0100} \sim 2.22 \times \Delta a_{01} M_{1000} \sim 5.22 \times \Delta a_{10} M_{1000} \sim 4.20 \times \Delta a_{10}$$

M<sub>1110</sub> ~ 1.36 × Δa<sub>11</sub> M<sub>1101</sub> ~ 1.30 × Δa<sub>11</sub>

 $M_{1100} \sim 0.27 \times \Delta a_{11}$  (scales quadratically, but outside scope of Coulton [2018] algorithm.)

9.159



Alex's careful work shows C18 (deterministic, no free parameters) algorithm systematically **under-corrects BFE** along the parallel direction. Better performance along serial direction might be coincidental. **NB:** Higher order moment problems may be present.

#### .. what gives !? 7



.. And what did other papers, exploring the limits of high fidelity correction have to say?



A selection of statements from recent publications on the state-of-the-art:

- Coulton '18: ... We can see that we meet the requirements for the Final HSC dataset, but that we need to *improve our correction by a factor of two or more* in order to reach the required levels for LSST.
- Lage '19 (arXiv:1911.09567v1): ... We also show, assuming sufficient care is paid to calculating the correction kernel, that the Coulton algorithm does in fact correct 90% or more of the BF effect on measured spots images. ... Correcting 90% of the effect should get us down near m = 0.006 To achieve the desired levels of m ≈ 0.001 0.003, *we need to do a factor of 2-5 times better*. ... Algorithm improvement will continue as more data becomes available from a larger sample of sensors.
- Broughton et al. '23: (arXiv:2312.03115v1) Our findings also motivate a detailed study on more realistic PSF stars and how measurement errors from BF could ultimately impact cosmology and other science goals. Ultimately, it is important to characterize the sensitivity of cosmological parameters to observables biased with BFEs. Even with state-of-the-art correction techniques, *the residual effects could represent a significant component* of the systematics error budget for cosmological analyses of LSST observations.



Let's take the overlap formalism with drift calculation as *ground truth* and see how Coulton (2018) performs



- Recall  $\Delta a_{ij}$  are **observable** via correlations,  $M_{iji'j'}$  are not.
- How does the C18 algorithm reproduce  $M_{iji'j'}$ ?? **Test** by integrating  $\Delta a_{ij}$  twice to obtain the kernel *K* according to:  $(\nabla^2 K + \Delta a)_{ij} = 0$  via successive over relaxation (SOR). Compare  $(\nabla K)_{ij}$  with corresponding elements of  $M_{iji'j'}$ .
- Use BFE source levels corresponding to the flat field statistical fluctuations (Δa<sub>00</sub>~5e-4 pix<sup>2</sup>) instead of source levels near full well. Those have not been observed.
- On the following slides, the following are compared:
  - Measured  $\Delta a_{ij}$  are compared to  $\Delta a_{ij}$  based on the computed kernel  $(-\nabla^2 K)$ .
  - Measured overlaps  $M_{iji'j'}$  are compared to corresponding element of the kernel-based boundary displacement  $\overrightarrow{\nabla}K$ .



### Let's take the overlap formalism with drift calculation as *ground truth* and see how Coulton (2018) performs (2)



#### Subset of calculation results and SOR derived kernel. Calculation extends out to +/- 15 pixels in both axes.

da: (pix <sup>2</sup> )	
+3.15088e-07 +4.16141e-07 +5.38447e-07 +6.67701e-07 +7.73036e-07 +8.14609e-07 +7.73036e-07 +6.67701e-07 +5.38447e	07 +4.16141e-07 +3.15088e-07
+4.15933e-07 +5.82952e-07 +8.0997e-07 +1.09104e-06 +1.35794e-06 +1.4745e-06 +1.35794e-06 +1.09104e-06 +8.0997e	07 +5.82952e-07 +4.15933e-07
+5.35508e-07 +8.06878e-07 +1.2457e-06 +1.9283e-06 +2.77118e-06 +3.21868e-06 +2.77118e-06 +1.9283e-06 +1.2457e-06	06 +8.06878e-07 +5.35508e-07
+6.61191e-07 +1.07824e-06 +1.90459e-06 +3.66419e-06 +7.07441e-06 <u>+9.92344e-06</u> +7.07441e-06 +3.66419e-06 +1.90459e	06 +1.07824e-06 +6.61191e-07
+7.62878e-07 +1.33013e-06 +2.67477e-06 +6.65653e-06 +1.91584e-05 +6.43842e-05 +1.91584e-05 +6.65653e-06 +2.67477e	06 +1.33013e-06 +7.62878e-07
+8.02762e-07 +1.43811e-06 +3.06201e-06 +8.68985e-06 +2.07967e-05 -0.000530519 +2.07967e-05 +8.68985e-06 +3.06201e	06 +1.43811e-06 +8.02762e-07
+7.62878e-07 +1.33013e-06 +2.67477e-06 +6.65653e-06 +1.91584e-05 +6.43842e-05 +1.91584e-05 +6.65653e-06 +2.67477e	06 +1.33013e-06 +7.62878e-07
+6.61191e-07 +1.07824e-06 +1.90459e-06 +3.66419e-06 +7.07441e-06 +9.92344e-06 +7.07441e-06 +3.66419e-06 +1.90459e	06 +1.07824e-06 +6.61191e-07
+5.35508e-07 +8.06878e-07 +1.2457e-06 +1.9283e-06 +2.77118e-06 +3.21868e-06 +2.77118e-06 +1.9283e-06 +1.2457e-06	06 +8.06878e-07 +5.35508e-07
+4.15933e-07 +5.82952e-07 +8.0997e-07 +1.09104e-06 +1.35794e-06 +1.4745e-06 +1.35794e-06 +1.09104e-06 +8.0997e-	07 +5.82952e-07 +4.15933e-07
+3.15088e-07 +4.16141e-07 +5.38447e-07 +6.67701e-07 +7.73036e-07 +8.14609e-07 +7.73036e-07 +6.67701e-07 +5.38447e	07 +4.16141e-07 +3.15088e-07
kernel: (pjx <sup>4</sup> )	
-1.45285e-06 -2.21647e-06 -3.08762e-06 -3.96849e-06 -4.67875e-06 -4.97416e-06 -4.68952e-06 -3.98976e-06 -3.11885e	06 -2.25687e-06 -1.50139e-06
-2.29881e-06 -3.48221e-06 -4.91856e-06 -6.48624e-06 -7.87442e-06 -8.52155e-06 -7.88567e-06 -6.50846e-06 -4.9512e	06 -3.52444e-06 -2.34955e-06
-3.30695e-06 -5.09516e-06 -7.44567e-06 -1.02922e-05 -1.31868e-05 -1.48441e-05 -1.31985e-05 -1.03152e-05 -7.47948e	06 -5.1389e-06 -3.35952e-06
-4.37826e-06 -6.97057e-06 -1.07407e-05 -1.59969e-05 -2.25261e-05 <u>-2.77065e-05 -2.25381e-05</u> -1.60205e-05 -1.07754e	05 -7.01552e-06 -4.43226e-06
-5.28796e-06 -8.76493e-06 -1.4473e-05 -2.41116e-05 -4.03076e-05 -6.08601e-05 -4.03198e-05 -2.41357e-05 -1.45084e	05 -8.81073e-06 -5.34299e-06
-5.67973e-06 -9.67723e-06 -1.69686e-05 -3.23451e-05 -7.29104e-05 -0.00019951 -7.29227e-05 -3.23694e-05 -1.70044e	05 -9.72353e-06 -5.73536e-06
-5.27632e-06 -8.75293e-06 -1.44607e-05 -2.40992e-05 -4.02951e-05 -6.08477e-05 -4.03075e-05 -2.41237e-05 -1.44966e	05 -8.79937e-06 -5.33213e-06
-4.35529e-06 -6.94689e-06 -1.07165e-05 -1.59723e-05 -2.25014e-05 -2.76819e-05 -2.25138e-05 -1.59967e-05 -1.07522e	05 -6.9931e-06 -4.41081e-06
-3.27322e-06 -5.06038e-06 -7.41012e-06 -1.02562e-05 -1.31506e-05 -1.4808e-05 -1.31628e-05 -1.02802e-05 -7.44538e	06 -5.10599e-06 -3.32802e-06
-2.25519e-06 -3.43723e-06 -4.87259e-06 -6.43966e-06 -7.8276e-06 -8.47488e-06 -7.8395e-06 -6.46315e-06 -4.90709e	06 -3.48186e-06 -2.30881e-06

Serial direction  $\rightarrow$ 



### Let's take the overlap formalism with drift calculation as *ground truth* and see how Coulton (2018) performs (3)



Pixel area distortions are in orange boxes; overlaps (M<sub>iii'i</sub>) are in between. 240 e<sup>-</sup> source at 00..



Direction

Serial

Direction

Parallel



Let's take the overlap formalism with drift calculation as *ground truth* and see how Coulton (2018) performs (4)

#### Similar trend, more pronounced for a ~36ke<sup>-</sup> BFE source (x148 of source on previous slide)



VERA C. RUBIN OBSERVATORY



# Let's take the overlap formalism with drift calculation as *ground truth* and see how Coulton (2018) performs (5)

#### But cleans up substantially when we artificially dial in x100 weaker barriers



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Direction

Serial

Parallel Direction



## Let's take the overlap formalism with drift calculation as *ground truth* and see how Coulton (2018) performs (6)



Fool around further by altering one of the model parameters (channel stop barrier)

- Persistent underestimation of M<sub>0100</sub> and overcorrection of M<sub>1000</sub> using C18 method.
- Noted improvement in C18 method when electrostatic barrier strengths are dialed down to 1% level.
- Tried scaling one of the parameters to investigate effect on the C18 errors in recovering  $M_{0100}$  etc. and tabulated observables (e.g.  $\Delta a_{01} / \Delta a_{10}$ ).
- Result: C18 works best when either:  $\Delta a_{01} / \Delta a_{10} \sim 1;$

barriers are very weak.

 We could have asked vendors to provide sensors with Δa<sub>01</sub> / Δa<sub>10</sub> ~ 1 as a requirement !!





Sources of nonlinearity that cause problems with C18 restriction of fixed *K*, ranked



Mis-match of gradient-of-derived-scalar-potential  $(\vec{\nabla}K)$  wrt computed boundary shifts tops the short list of mechanisms that would affect *mapping of correlation measurements into boundary shifts*:

- Shape of pixels when barrier strengths are not equal, i.e. C01/C10 ~ 3-5 (correction error +/- 10%);
- Attraction & shift of channel potential well toward polysilicon gates as channel fills (correction error +/- 6-7% at 36ke<sup>-</sup> for N<sub>d</sub>~5e15 cm<sup>-3</sup>);
- 3. Presence of recorded flat field level per pixel when BFE is being characterized in flat pairs: (correction error +/- 4% at 36ke<sup>-</sup>);
- Extended distribution of collected conversions (distributed into 3x3 grid) at channel vs. modeled (point-like) spatial distribution (correction error +/- 0.06%). cf. Lage et al. '21;



![](_page_15_Picture_2.jpeg)

- We've used a first principles-based drift simulator tool, tuned to reproduce various observable signatures/artifacts seen in LSST sensors to compute how charge is partitioned and ultimately recorded. It permits us to evaluate how current baseline correction algorithms (C18) perform and to qualitatively understand their limits.
- C18 can't be tweaked; it features a signal-independent kernel *K* derived directly from observables and is a scalar quantity. We've shown systematic deviations between ground truth pixel boundary shifts and C18 corrections that can partially explain residual BFE terms observed by Broughton et al.
- Time to revive old/altérnate BFE correction strategies?
  - Use the sparse "pixel overlap matrix" to correct pixel-by-pixel (a la Astier et al.)
  - Compute previously proposed book-keeping terms (pixel areas, astrometric shifts, 2<sup>nd</sup> moments), either for uncorrected data or representing we C18 apparently under- or over-corrects. *One option* could be for no C18 correction in ISR (but carry these pixel terms based on recorded image esp. for high <u>S</u>/N stars).
  - Work out a way to incorporate properties of **K** such that  $\nabla \times \nabla K \neq 0$  (and allow it to be signal dependent while we're at it, we see this in drift calculations).

![](_page_16_Picture_0.jpeg)

![](_page_16_Picture_1.jpeg)

![](_page_16_Picture_2.jpeg)

![](_page_16_Picture_3.jpeg)

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![](_page_17_Picture_0.jpeg)

![](_page_17_Picture_2.jpeg)

Iteratively. Maybe one step is good enough? Similar to other work, e.g. substitute  $M_{ijij}$  for  $\nabla K$  wherever it appears in C18. Otherwise, keep in mind that the  $M_{ijij}$  overlap matrix elements have more indices because they account for all area displacements between nominal and recorded pixels. 5<sup>th</sup> index may be introduced to allow overlap coefficients to evolve with signal level:

Evaluate local flux at overlap centroids using recorded image  $R_{mn}$ :  $S_{p,q,p+\Delta i,q+\Delta j} \equiv \frac{\sum_{mn} R_{mn} (\delta_{p,m} \delta_{q,n} + \delta_{p+\Delta i,m} \delta_{q+\Delta j,n})}{\sum_{mn} (\delta_{p,m} \delta_{q,n} + \delta_{p+\Delta i,m} \delta_{q+\Delta j,n})}$ 1<sup>st</sup> order displaced signal estimate using local flux S & recorded flux R<sub>mn</sub>:  $D_{p,q,p+\Delta i,q+\Delta j}^{(0)} \equiv S_{p,q,p+\Delta i,q+\Delta j} \cdot \left[ \mathbf{M}_{i,j,i',j'k} R_{mn} \right] \cdot \left( \frac{1}{2} \right) \cdot \mathbf{W}_{p,q,\Delta i,\Delta j,i,j,i',j',m,n,k}$ 1<sup>st</sup> order estimate for incident flux F:  $F_{mn}^{(0)} \equiv R_{mn} + \sum_{i,j} \left( D_{m,n,i,j}^{(0)} - D_{i,j,m,n}^{(0)} \right)$ Refine displaced signal map estimate D  $D_{p,q,p+\Delta i,q+\Delta j}^{(0)} \equiv \sum_{\alpha}^{I-1} D_{p,q,p+\Delta i,q+\Delta j,\tau}^{(0)}$ evaluated for timeslices:  $\equiv S_{p,q,p+\Delta i,q+\Delta j} \cdot \sum_{\tau=0}^{T-1} \left[ \mathbf{M}_{i,j,i',j'k} R_{mn} \right] \cdot \left( \frac{\tau + \frac{1}{2}}{T} \right) \cdot \mathbf{\Omega}_{p,q,\Delta i,\Delta j,i,j,i',j',m,n,k,\tau,T}.$ Iterate incident flux F estimate and while displaced signal D  $F_{mn}^{(iter.+1)} \equiv R_{mn} + \sum_{i=1}^{n} \left( D_{m,n,i,j}^{(iter.+1)} - D_{i,j,m,n}^{(iter.+1)} \right)$ by resolving in time slices. End with convergence condition:  $D_{p,q,p+\Delta i,q+\Delta j,\tau}^{(iter,+1)} \equiv S_{p,q,p+\Delta i,q+\Delta j} \cdot \mathbf{M}_{i,j,i',j'k} \cdot \left[ F_{mn}^{(iter)} \left( \frac{\tau+\frac{1}{2}}{T} \right) - \sum_{a,-a}^{\tau-1} \left( D_{m,n,i,j,\theta}^{(iter)} - D_{i,j,m,n,\theta}^{(iter)} \right) \right]$  $\left[ \left( \frac{\tau+1}{T} \right) \cdot R_{m,n} \right] \longrightarrow \left[ \left( \frac{\tau+1}{T} \right) \cdot F_{mn}^{(iter.)} - \sum_{s=s}^{\tau-1} \left( D_{m,n,i,j,\theta}^{(iter.)} - D_{i,j,m,n,\theta}^{(iter.)} \right) \right]$  $\cdot \Omega_{p,q,\Delta i,\Delta j,i,j,i',j',m,n,k,\tau,T}$ 

![](_page_18_Picture_0.jpeg)

![](_page_18_Picture_2.jpeg)

In the previous expressions, these matrices are meant to simplify computation in nested loops over the signal- and time-resolved slices of correction. Basically, they are an amalgamation of Kronecker and Heaviside functions:

$$\mathbf{W}_{p,q,\Delta i,\Delta j,i,j,i',j',m,n,k} \equiv \delta_{m+i,p} \delta_{n+j,q} \delta_{m+i',p+\Delta i} \delta_{n+j',q+\Delta j} \cdot \Theta \left( U_k - R_{mn} \right) \cdot \Theta \left( R_{mn} - L_k \right)$$

$$\Omega_{p,q,\Delta i,\Delta j,i,j,i',j',m,n,k,\tau,T} \equiv \delta_{m+i,p}\delta_{n+j,q}\delta_{m+i',p+\Delta i}\delta_{n+j',q+\Delta j} \\ \cdot \Theta\left(U_k - \left(\frac{\tau+1}{T}\right)R_{mn}\right) \cdot \Theta\left(\left(\frac{\tau+1}{T}\right)R_{mn} - L_k\right).$$

![](_page_19_Picture_0.jpeg)

### Pixel distortion bookkeeping as an alternate BFE correction strategy appropriate for moments (R18)

![](_page_19_Picture_2.jpeg)

![](_page_19_Figure_3.jpeg)