SLAC MeV UED Experience

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UED Opportunities for

Dynamical Imaging of Materials

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- Brief Introduction to MeV UED
- SLAC MeV UED Experimental Setup and Science Highlights
- Perspectives and Comments

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Probing the ultrasmall and ultrafast

Atomic spatial-temporal scales

- Angstrom (Å, 10⁻¹⁰ m)
- femtosecond (fs, 10⁻¹⁵ s)

occurs on 0.1 s time scale







The 1st step of vision process happens over ~100 fs time scale

Probing the ultrasmall and ultrafast



- Atomic structure and phase (function) correlation
- Dynamical process can lead to exotic properties not accessible in thermal equilibrium
- Need **fs-Å** probe to study the correlation

VO₂ insulator-metal phase transition



A. Sood et al., Science. 373, 352–355 (2021).

Room temperature superconductor does exist! – in a transient state over ps-ns scale for K_3C_{60}



Ultrafast probes – X-ray and Electron

Unique features of electron compared to x-ray

- 10⁴ 10⁶ times larger scattering cross sections
- shorter wavelength, higher spatial resolution
- 10³ times less radiation damage







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Real space vs diffractive (reciprocal space) imaging



Real space imaging

- Good at resolving local structure, such as defects, dislocations, etc.
- Can reach Å spatial resolution
- Not yet reaching fs temporal resolution at the same time

Diffractive (reciprocal space) imaging



- Conjugated to real space imaging
- Ensemble average imaging
- Can reach Å-fs spatial-temporal resolution

Molecule ensemble





real space

reciprocal space



Development of keV UED



Prof. Gérard Albert Mourou recipient of 2018 Nobel Prize for physic "for his invention of chirped pulse amplification"



FIG. 1. Schematic of picowecond electron-diffuscion apparatus. A stroak-curners tabe (deflection plates removed) is used to produce the electron public. The 25-keV electron public passes through the Al appelinen and produces a diffusc-tion pattern of the structure with a 20% exposure.





Prof. Ahmed Hassan Zewail recipient of 1999 Nobel Prize in Chemistry, "for showing that it is possible with rapid laser technique to study in slow motion how atoms in a molecule move



gas phase UED femtochemistry



Prof. R. J. Dwayne Miller Chemistry and Physics, University of Toronto



Prof. Jianming Cao Department of Physics Florida State University



Prof. Bradley J. Siwick Dept. of Phys. and Chem. McGill University



M. Chergui and A. H. Zewail, ChemPhysChem 10, 28-43 (2009)



Prof. Chong-Yu Ruan Dept. of Phys. and Astron. Michigan State University



Prof. Ralph Ernstorfer Dept. of Phys. Chem. Fritz Haber Institute









Prof. Peter Baum Department of Physics Universität Konstanz



Prof. Claus Ropers Max Planck Institute for Multidisciplinary Sciences



Prof. David Flannigan Chemical Engineering and Materials Science University of Minnesota

and many other researchers...

Space-charge forces suppression with relativistic electrons

- shorter bunch ⇒ higher time resolution
- more electrons in a bunch \Rightarrow single shot measurement

Negligible pump-probe velocity mismatch

• Δt_{vm} <10 fs for 3 MeV e beam passing 150 μ m gas target









Really flat Ewald sphere

Ewald sphere radius $\propto \frac{1}{\lambda} = \frac{p}{h} = \frac{\beta \gamma m_0 c}{h}$

- k >> G \rightarrow relativley flat Ewald-sphere
- extended reciprocal lattice "points" (finite sample thickness)
- Bragg-condition only approximately fulfilled.



•mismatch s needs to be considered



• Bragg-conditions (almost) exactly fulfilled.





Shorter wavelength

- higher momentum transfer ⇒ higher spatial resolution
- smaller diffraction angle → longer sample-to-detector distance → flexible environment for "dirty" samples



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SLAC MeV UED Milestones





- User Run03 solid UED
- User Run04 gas UED













MeV UED instrument overview



X. Shen, et al., Struct. Dyn. 6, 054305 (2019)

Multifunctional platform for ultrafast science



Ultrafast science enabled by MeV UED



Phase switch with a single flash of light (**Sci. Adv.** 4, eaau5501 (2018)).



Operando characterization of a quantum electronic device (**Science** 373, 352–355 (2021)).



Resolving ultrafast phase transitions (**Science** 360 1451–1455 (2018))



Ultrafast topological switch by Time-varying shear strain (*Nature* 565, 61-77 (2019))



Bond-breaking & nuclear wavepacket passing through conical intersections (*Science* 361 64–67 (2018))



Light induced CDW (*Nat. Phys.* 16, pages159– 163(2020)).



Ring-opening & ground state dynamics (*Nat. Chem.* 11, 504–509 (2019)).



Imaging conformer-specific photochemistry(*Science*. 374, 178–182 (2021).).



Simultaneous observation of nuclear and electronic dynamics (**Science.** 368, 885 (2020)).



Ultrafast hydrogen bond strengthening in liquid water(*Nature*. 596, 531(2021)).

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Spin-mediated shear oscillators in a van der Waals antiferromagnet (*Nature*. 620, 988–993 (2023))).



Phonon-mediated heat transfer in a van der Waals heterostructure (*Nat. Nanotechnol.* 18, 29–35 (2023)).



Incipient plasticity in dynamically compressed matter (*Nat Commun*. 13, 1055 (2022)).



Rehybridization dynmaics into a pericyclic minimum state (*Nat. Commu*. 14, 2795 (2023)).

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- Temporal compression using RF bunching cavity
 - achieved $\tau_{electron} = 10$ fs (rms)=23.5 fs (fwhm) bunch length
 - Introduced large τ_{TOA} > 100 fs



- achieved $\tau_{electron}$ = 39 fs (rms)=91.65 fs (fwhm)
- reduced τ_{TOA} by a factor of 2 to 31 fs (rms) = 72.9 fs (fhwm)
- Low THz laser conversion efficiency (~0.1%)
- Introduced large emittance growth due to the broadband nature the THz pulse



E. C. Snively, et al. Phys. Rev. Lett. 124, 054801 (2020).

Bunch compression and timing jitter reduction by double bend achromat (DBA)

- Isochronous beamline with DBA to minimize au_{TOA}
- Use space charge effect to induce correct chirp for $au_{electron}$ compression
- Achieved ~30 fs $\tau_{electron}$ and τ_{TOA}
- Fixed energy, fixed high charge, large emittance growth



H. W. Kim, et al. Nat. Photonics 14, 245–249 (2020)



F. Qi, et al. Phys. Rev. Lett. 124, 134803

THz-based arrival time monitor



2.0 **e**

1.5

0.5

0.0

-0.5

-1.0

-1.5

1.5

(i 1.5 1.0 **g**

0.0

-3

-2 -1

arb 0.5

BO -0.5

-3

-2 -1 0

pointing stability

200 400 600 800 1000

3.0 µrad rms

number of shots

0 20 40 60 0 20 40 60 0 20 40 60 0 20 40 60

20 10 **f**

-10

-20_____0

(pe

y_o' (µra

streaking speed

2 3 4

1 2 f(THz)

4

D = 7.4 µrad/fs

10-3

i

t (ps)

t (ps)

0

roid y_{o}^{\prime} (mrad) 1.0

-0.60 ps

3.1 MeV electrons, 50 um gap, 100 um thick slit

Parallel-plate waveguide for enhanced streaking



On reciprocal space resolution (spatial coherence)





Moire patterns from twisted bilayer graphene

Higher spatial coherence \Rightarrow narrower Bragg peak

Ultracold Electrons via Near-Threshold Photoemission from Single-Crystal Cu(100)

Siddharth Karkare^{®*} Physics Department, Arizona State University, Tempe, Arizona 85282, USA

Gowri Adhikario and W. Andreas Schroeder Department of Physics, University of Illinois at Chicago, Chicago, Illinois 60607, USA

J. Kevin Nangoi, Tomas Arias, and Jared Maxson Department of Physics, Cornell University, Ithaca, New York 14853, USA

Howard Padmore[®] Lawrence Berkeley National Lab, Berkeley, California 94720, USA



On electron beam brightness

E beam transverse brightness dependence on photoelectron gun gradient

 $B_{\perp}^{pancake} \propto E$, $B_{\perp}^{cigar} \propto E^{3/2}$



FIG. 1. Comparison of 1.6-cell and 1.4-cell guns for (a) the output kinetic energy and energy spread (10-ps-long, low-charge electron beam) and (b) the acceleration gradient seen by the beam.

R. K. Li, P. Musumeci, Phys. Rev. Applied. 2, 024003 (2014).

High gradient comes with high dark current – exponential dependence!



Figure 7: Dark charge vs. the peak cathode field.

On nanoprobe

Require low emittance and high brightness e beam Mechanical vibration and other environmental noises needs to be well diminished





SLAC 4D (x, y, q_x, q_y) Scanning Transmission Electron Microscopy

On electron beam detector

Every electron counts! Direct electron detectors







© Nobel Media. III. N. Elmehed Jacques Dubochet Prize share: 1/3

© Nobel Media, III, N. Elmehed Joachim Frank Prize share: 1/3 Prize share: 1/3

Critical components enabling full capability of cryoEM



Atomic resolution single particle imaging enabled by direct electron detector



Dr. Peter Denes, LBNL





T. Vecchione et al., Review of Scientific Instruments. 88, 033702 (2017).



Comments





Beam energy: a few MeV

Time resolution: science dependence (fs-ps)

Rep rate: sample dependence

Emittance: the lower the better with reasonable flux

Sample: ~100 nm thin film

Vacuum: gun 1e-10 torr, sample <1e-4 torr, enviornmental TEM possible.



FIG 3. Select zoology of ultrafast photoinduced dynamics and their approximate timescales. Boundary conditions become increasingly important to energy dissipation with increasing time owing to dynamics encompassing the entire specimen (*i.e.*, initiating with charge-carrier dynamics and increasing in scale to whole-crystal/specimen motion and thermal diffusion). Relatively long-lived states and dissipation channels will limit the UEM operational f_{rep} if full recovery is desired, which can take microseconds or longer. For comparison, the f_{rep} (and the associated time between packets) needed to capture full recovery of the discrete examples is shown in relation to the overall chain of events, further illustrating the need for $f_{rep} < MHz$ for full recovery of the broadest range of chemical and materials phenomena.

Flannigan, D. J., The Journal of Chemical Physics 157, 180903 (2022).

Comments

Multi-modal experiment

Less sample damage

- less energy deposition
- pulsed operation, lower dose rate, less damage to dose rate sensitive matter
- gentle probe for multi-modal

2. Comment on e-beam thermal effects on the UED data.

Response: We thank the referee for the question on electron beam induced thermal effects on the UED data. We carried out the following calculations to evaluate this question. We quote the stopping power of ~4 MeV electron in graphite as $\frac{1}{\rho} \frac{dE}{dx} = 1.697$ MeV cm²/g from the ESTAR database of National Institute of Standards Technology and (https://physics.nist.gov/PhysRefData/Star/Text/method.html). The mass density of graphite is $\rho = 1.7$ g/cm (to first order should be also applicable for graphene). When a 4-MeV electron pulse, with number of electrons $n = 2.8 \times 10^4$ and FWHM spot size $d = 100 \ \mu m$, passes through a graphite lattice with a thickness of $t \approx 0.34$ nm (the graphite interlayer spacing), the average deposited electron beam fluence is calculated as

$$F_{e} = \frac{\Delta E}{A} = \frac{\left(\frac{1}{\rho}\frac{dE}{dx}\right)\rho tn}{d^{2}} = \frac{1.697 \left[MeV\frac{cm^{2}}{g}\right] \cdot 1.7 \left[\frac{g}{cm^{3}}\right] \cdot 0.34 \cdot 10^{-7} [cm] \cdot 2.84 \cdot 10^{4}}{(100 \cdot 10^{-4} \ [cm])^{2}} = 27.5 \left[\frac{MeV}{cm^{2}}\right]$$

On the other hand, the lowest pump laser fluence incident on the sample in the experiment was ~ 3 mJ/cm². With a typical value 2.3% absorption of a graphene layer, the deposited laser fluence in a monolayer graphene is around F_i =6.9 x 10⁻² mJ/cm², which is 7 order of magnitude larger than the deposited electron fluence F_e . Thus, the laser-induced heating effect is dominant.





Comments

In-situ experiments

Flexible sample environment

- Bragg's law $m\lambda = 2d \sin \theta \leftrightarrow \sin \theta = \frac{m\lambda}{2d}$
- For MeV electrons, $\lambda \sim 0.001$ Å, $\theta \sim$ mrad
- large space chamber, long distance separation between parts

Sample
$$q = \frac{4\pi}{\lambda}\sin\theta$$





All-solid electron-transparent implementation





Imaging Diffraction Spectroscopy





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