Bayesian inference and deterministic anisotropy for molecular geometry retrieval in gas-phase diffraction experiments

UED Opportunities for Dynamical Imaging of Materials 11/6/2023

Kareem Hegazy, Varun Makhija, Phil Bucksbaum, Jeff Corbett, James Cryan, Nick Hartmann, Markus Ilchen, Keith Jobe, Renkai Li, Igor Makasyuk, Xiaozhe Shen, Xijie Wang, Stephen Weathersby, Jie Yang, **Ryan Coffee**





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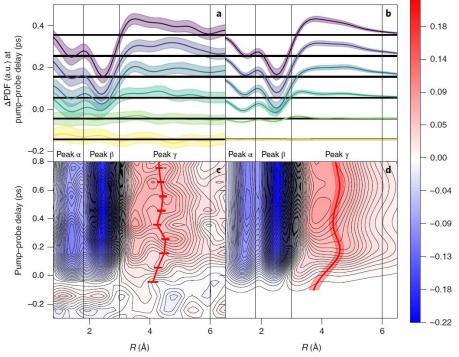
We introduce a data driven approach to rigorously retrieve molecular frame geometries while reducing simulation requirements.

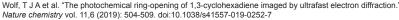
- 1. Current approaches and the curse of dimensionality
- 2. Anisotropy reveals the molecular frame (MF)
- 3. Bayesian Inference
 - a. Overcoming the curse of dimensionality
 - b. Metropolis Hastings Algorithm
 - c. Results
 - d. Systematic errors from single geometry assumptions
- 4. Future Application to Excited State Geometric Dynamics
- 5. Summary

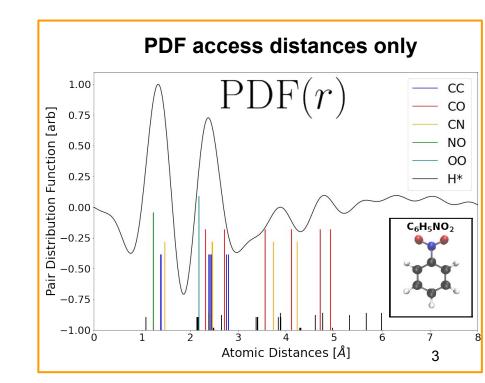
Current Molecular Geometry Retrieval Methods Traditional and New Methods



- Pair distribution function (PDF)
- Compare PDF, or diffraction, with theory to interpret the molecular geometry transience







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Current Molecular Geometry Retrieval Methods Traditional and New Methods



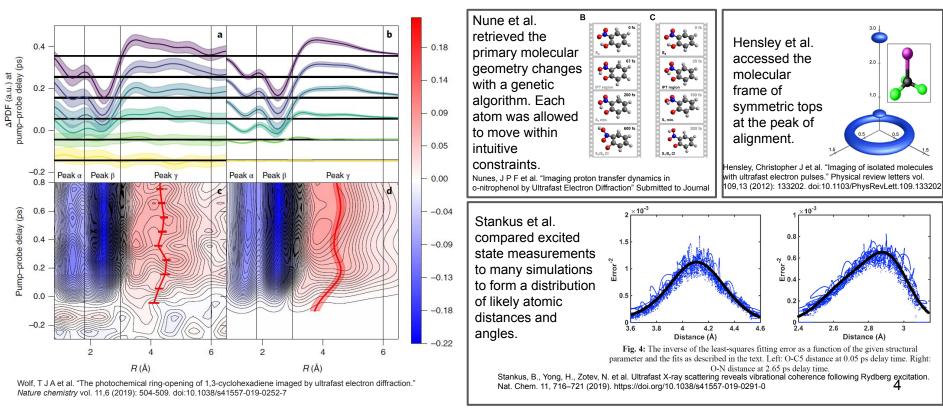
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Traditional Method

- Pair distribution function (PDF)
- Compare PDF, or diffraction, with theory to interpret the molecular geometry transience

Data Focused/Driven Method

- Use ML to optimize primary features
- Use molecular alignment to access the molecular frame of symmetric tops
- Use many simulated geometries to statistically improve precision



Current Molecular Geometry Retrieval Methods Curse of Dimensionality



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The number of geometries to sample (S) grows exponentially, for a grid search

 $S = R^{3N-6}$

R: Number of sample points per dimension N: Number of atoms V: Percent of volume

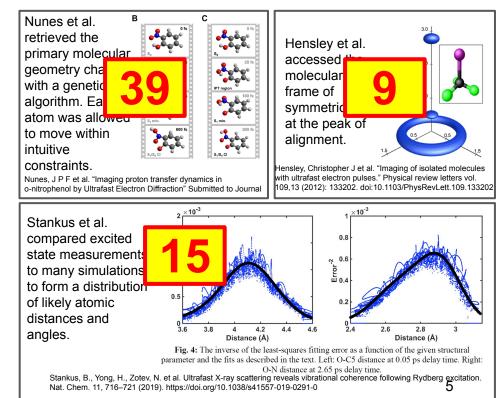


Nitrobenzene: 21 dimensions (ignoring H)

S	R	V R=10	V R=4
10 ⁶	1.93	10 ⁻¹⁴ %	2.3×10 ⁻⁵ %
10 ⁹	2.40	10 ⁻¹² %	2.3×10 ⁻³ %

Data Focused/Driven Method

- Use ML to optimize primary features
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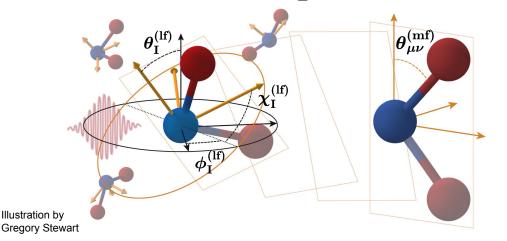
Anisotropy Reveals the Molecular Frame (MF)

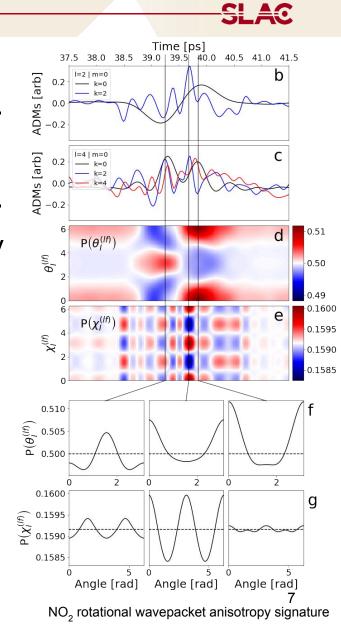
Anisotropy Reveals the Molecular Frame (MF)

- Anisotropy provides constraints on molecular frame (MF) degrees of freedom.
- Combining many measurements (constraints) allows one to retrieve the MF.

Application in photo-electron spectroscopy

- V. Makhija, et. al., (2016), arXiv:1611.06476 [physics.atom-ph]
- C. Marceau, et. al., Phys. Rev. Lett. 119, 083401 (2017)
- M. Gregory, et. al., (2020), arXiv:2012.04561 [physics.chem-ph]
- Use a stretched NO₂: an asymmetric top

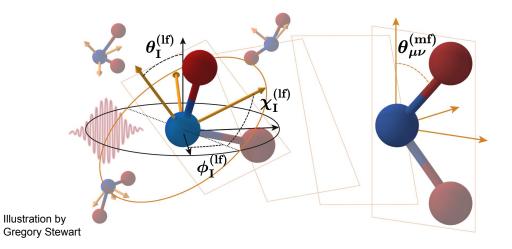




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$$\begin{split} \langle I(\mathbf{q},t) \rangle &= \mathcal{I}\bigg(\sum_{\mu} |f_{\mu}(q)|^{2} + \sum_{\mu,\nu:\mu\neq\nu} \operatorname{Re}\bigg\{f_{\mu}(q)f_{\nu}^{*}(q)\sum_{l} 4\pi i^{l} \\ &\times \sum_{m,k} (-1)^{k} \underbrace{Y_{l}^{m}\left(\theta_{q}^{(\mathrm{lf})},\phi_{q}^{(\mathrm{lf})}\right)}_{\mathrm{Lab \ Frame}} \langle \Psi(t)| \underbrace{D_{mk}^{l}\left(\phi_{\mathrm{I}}^{(\mathrm{lf})},\theta_{\mathrm{I}}^{(\mathrm{lf})},\chi_{\mathrm{I}}^{(\mathrm{lf})}\right)}_{\mathrm{Ensemble \ Anisotropy}} \underbrace{j_{l}(q\Delta R_{\mu\nu})Y_{l}^{-k}\left(\theta_{\mu\nu}^{(\mathrm{mf})},\phi_{\mu\nu}^{(\mathrm{mf})}\right)}_{\mathrm{Molecular \ Frame \ Structure}} |\Psi(t)\rangle\bigg\}\bigg) \end{split}$$

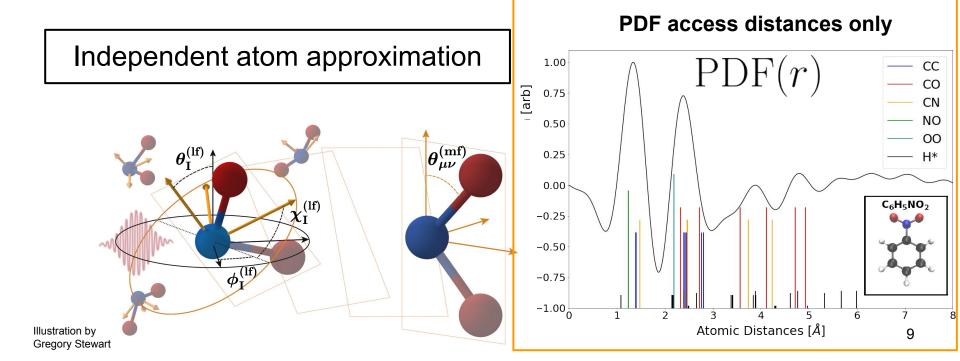
Independent atom approximation



$$\begin{split} \langle I(\mathbf{q},t) \rangle &= \mathcal{I}\bigg(\sum_{\mu} |f_{\mu}(q)|^{2} + \sum_{\mu,\nu:\mu\neq\nu} \operatorname{Re}\bigg\{f_{\mu}(q)f_{\nu}^{*}(q)\sum_{l} 4\pi i^{l} \\ &\times \sum_{m,k} (-1)^{k} \underbrace{Y_{l}^{m}\left(\theta_{q}^{(\mathrm{lf})},\phi_{q}^{(\mathrm{lf})}\right)}_{\mathrm{Lab \ Frame}} \langle \Psi(t)| \underbrace{D_{mk}^{l}\left(\phi_{\mathrm{I}}^{(\mathrm{lf})},\theta_{\mathrm{I}}^{(\mathrm{lf})},\chi_{\mathrm{I}}^{(\mathrm{lf})}\right)}_{\mathrm{Ensemble \ Anisotropy}} \underbrace{j_{l}(q \Delta R_{\mu\nu})Y_{l}^{-k}\left(\theta_{\mu\nu}^{(\mathrm{mf})},\phi_{\mu\nu}^{(\mathrm{mf})}\right)}_{\mathrm{Molecular \ Frame \ Structure}} |\Psi(t)\rangle\bigg\} \bigg)$$

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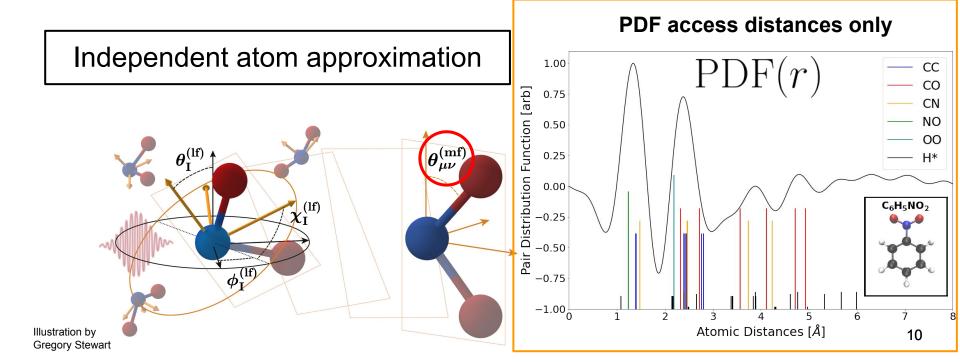
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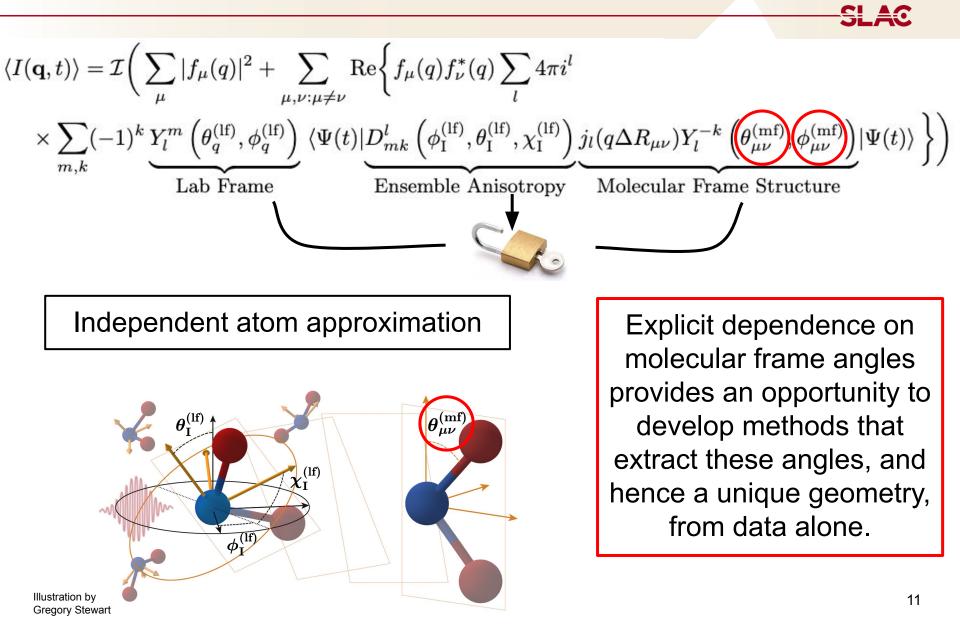


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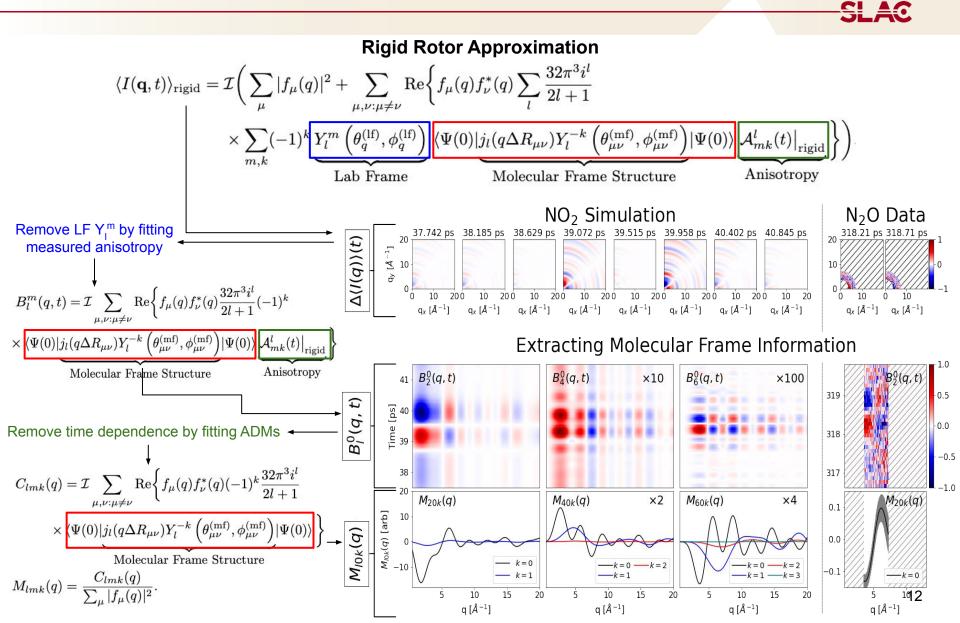
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Accessing the MF via Deterministic Anisotropy Rigid Rotor



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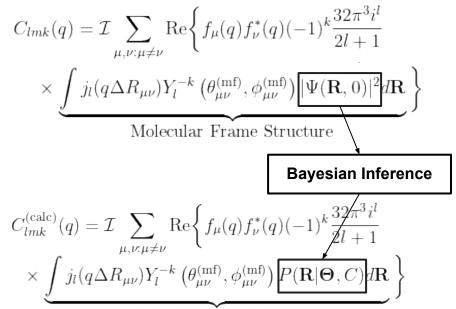
Bayesian Inference

Bayesian Inferencing Reframing the Problem

What we have: Molecular frame representations for the nuclear geometry with explicit dependence on pair-wise distances and angles

What we want: Invert the integral equation for $|\Psi(\mathbf{R}^{(mf)})|^2$

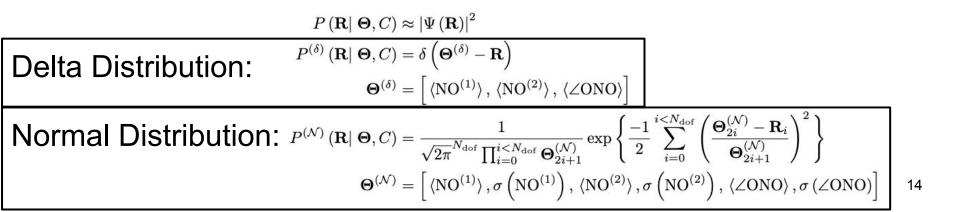
How we do it: Approximate $|\Psi(\mathbf{R}^{(mf)})|^2$ with a chosen distribution and solve for the model parameter ($\boldsymbol{\theta}$) distribution P($\boldsymbol{\theta}|$ C)



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Molecular Frame Structure

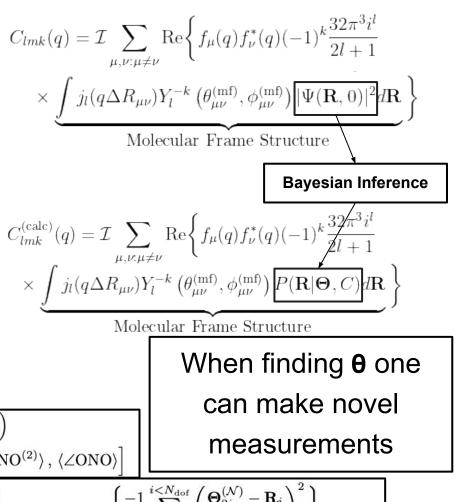


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$$\begin{array}{c}
\begin{array}{c}
P^{(\delta)}(\mathbf{R}|\,\Theta,C) = \delta\left(\Theta^{(\delta)} - \mathbf{R}\right) \\
\Theta^{(\delta)} = \left[\langle \mathrm{NO}^{(1)} \rangle, \langle \mathrm{NO}^{(2)} \rangle, \langle \angle \mathrm{ONO} \rangle\right]
\end{array}$$
Can make novel measurements measurements
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Normal Distribution:
$$P^{(\mathcal{N})}(\mathbf{R}|\,\Theta,C) = \frac{1}{\sqrt{2\pi}^{N_{\mathrm{dof}}}\prod_{i=0}^{i < N_{\mathrm{dof}}}\Theta^{(\mathcal{N})}_{2i+1}}\exp\left\{\frac{-1}{2}\sum_{i=0}^{i < N_{\mathrm{dof}}}\left(\frac{\Theta^{(\mathcal{N})}_{2i} - \mathbf{R}_{i}}{\Theta^{(\mathcal{N})}_{2i+1}}\right)^{2}\right\}\\
\end{array}\\
\left(\Theta^{(\mathcal{N})} = \left[\langle \mathrm{NO}^{(1)} \rangle, \sigma\left(\mathrm{NO}^{(1)}\right), \langle \mathrm{NO}^{(2)} \rangle, \sigma\left(\mathrm{NO}^{(2)}\right), \langle \angle \mathrm{ONO} \rangle, \sigma\left(\angle \mathrm{ONO} \rangle\right)\right]\right] 15$$

 $P(\mathbf{R} \mid \boldsymbol{\Theta}, C) \approx |\Psi(\mathbf{R})|^2$

Modeling $|\Psi(\mathbf{R}^{(mf)})|^2$ and Search Parameters

 $P(\mathbf{R}|\boldsymbol{\Theta}, C) \approx |\Psi(\mathbf{R})|^2$

 $\Theta^{(\mathcal{N})} =$

 $P^{(\delta)}(\mathbf{R}|\boldsymbol{\Theta}, C) = \delta\left(\boldsymbol{\Theta}^{(\delta)} - \mathbf{R}\right)$

 $\boldsymbol{\Theta}^{(\delta)} = \left| \left< \mathrm{NO}^{(1)} \right>, \left< \mathrm{NO}^{(2)} \right>, \left< \angle \mathrm{ONO} \right> \right.$

 $\sqrt{2\pi}^{N_{\mathrm{dof}}}\prod_{i=0}^{i< N_{\mathrm{dof}}} \mathbf{\Theta}_{2i+1}^{(\mathcal{N})}$

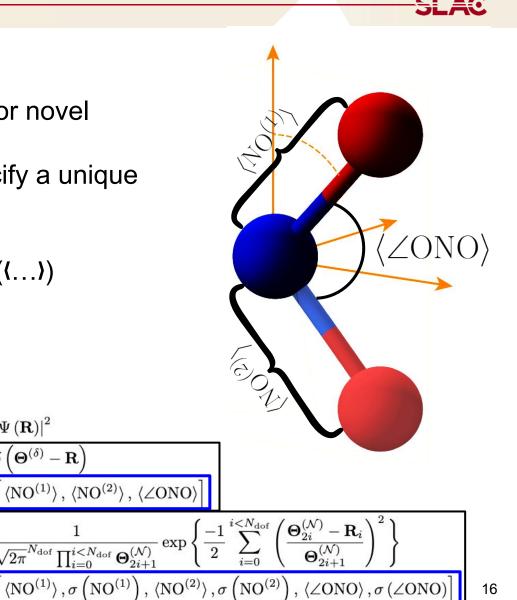
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- The **θ** parameterization allows for novel measurements
 - Degrees of freedom to specify a unique Ο geometry
 - Bond distances and angles Ο

Normal Distribution: $P^{(N)}(\mathbf{R}|\Theta, C) = -$

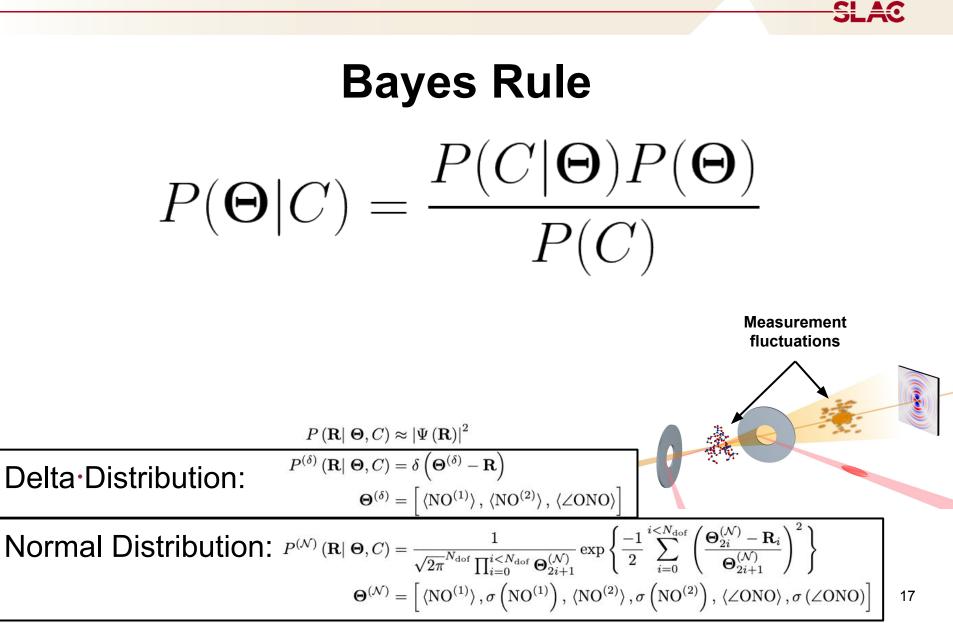
Delta Distribution:

Width of the wave packet $\sigma(\ldots)$ Ο



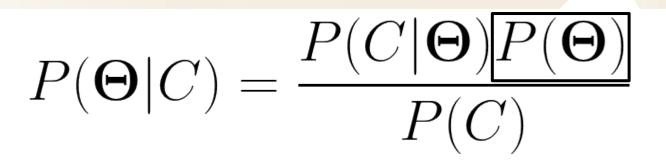
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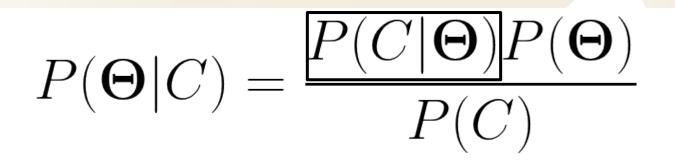




Prior Distribution $P(\Theta) = ext{constant}$

• The probability of observing these **Θ** parameters

- Relies on a priori chemical knowledge and chemical intuition
- A constant prior is chosen to have an unbiased **O** search
 - For unphysical values the prior is 0 e.g., negative bond distance

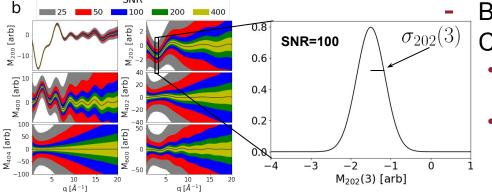


Likelihood Distribution

$$P(C|\mathbf{\Theta})$$

$$= \left[\prod_{lmk,q} \frac{1}{\sigma_{lmk}(q)\sqrt{2\pi}}\right] \exp\left\{\frac{-1}{2} \sum_{lmk,q} \left(\frac{C_{lmk}(q) - C_{lmk}^{(\text{calc})}(q, \Theta)}{\sigma_{lmk}(q)}\right)^{2}\right]$$

• Each measured $C_{Imk}(q)$ is it's own Gaussian probability distribution



- By the central limit theorem we expect $C_{lmk}(q)$ to be Gaussian
 - Each C_{Imk}(q) is a distribution of weighted sums of the measured pixels
 - $C_{lmk}(q)$ distribution has > 100 entries





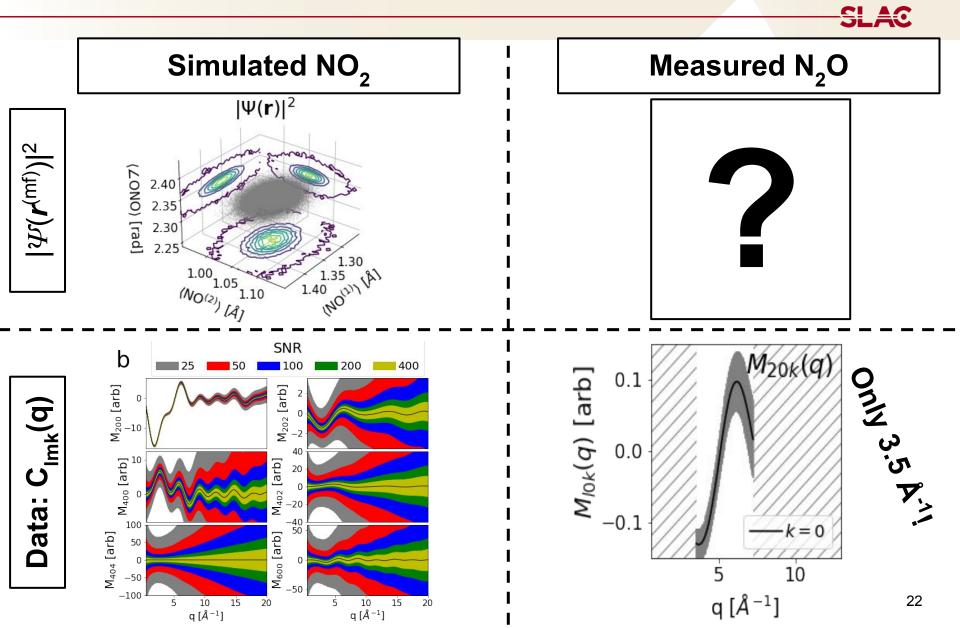
Metropolis Hastings Algorithm

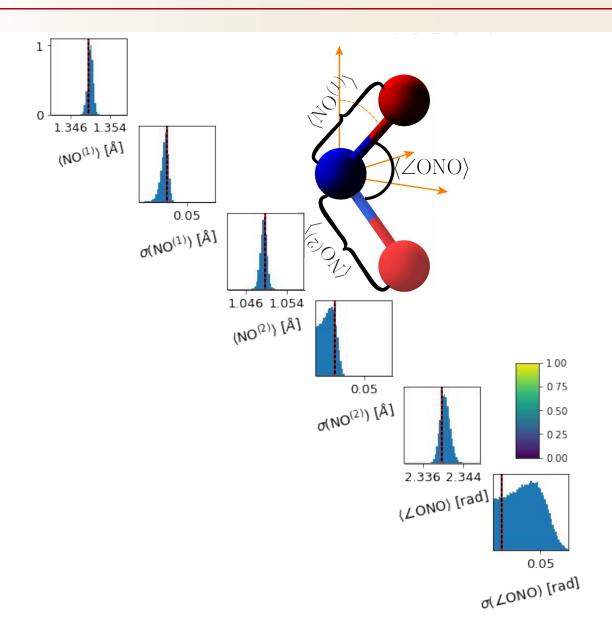
- Unbiased search for Θ
- Θ are preferentially selected based on their agreement with data
- Spend more time searching regions of high likelihood
 - $\circ~$ If Θ is twice as likely as Θ' we spend twice as much time sampling the area around Θ

$$P(\Theta'|\Theta) = \frac{P(\Theta'|C)}{P(\Theta|C)}$$
$$= \min\left[\frac{P(C|\Theta')P(\Theta')}{P(C|\Theta)P(\Theta)}, 1\right]$$

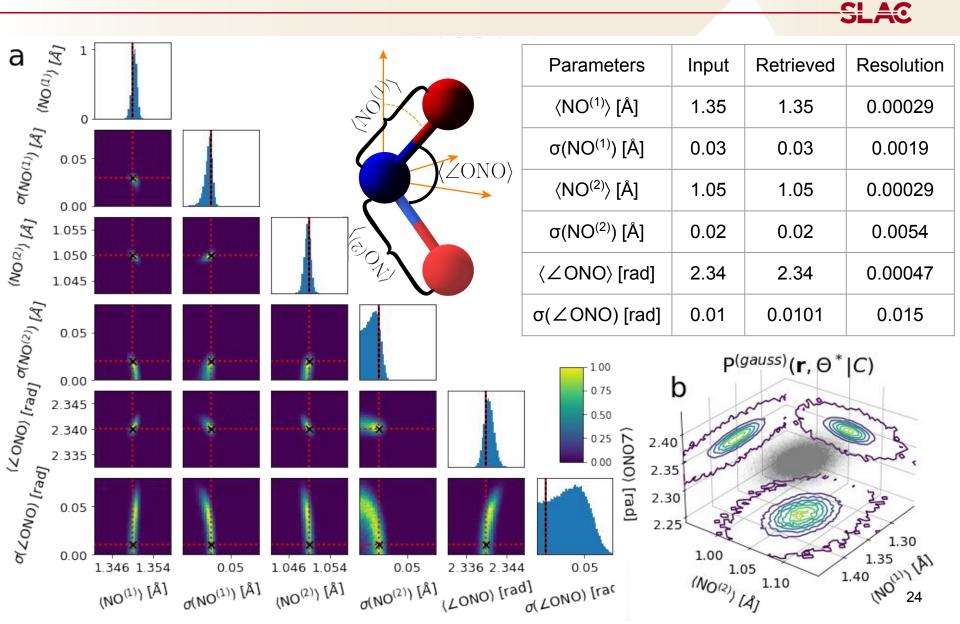
Application to simulated NO_2 and measured N_2O

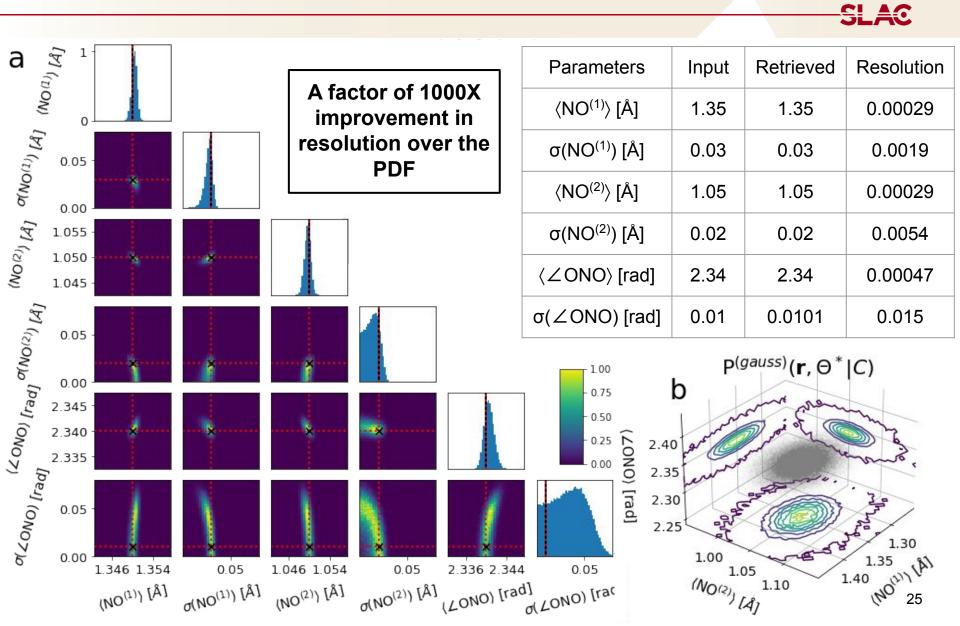
Retrieving the MF Geometry Probability Distribution U SE Algorithm Input

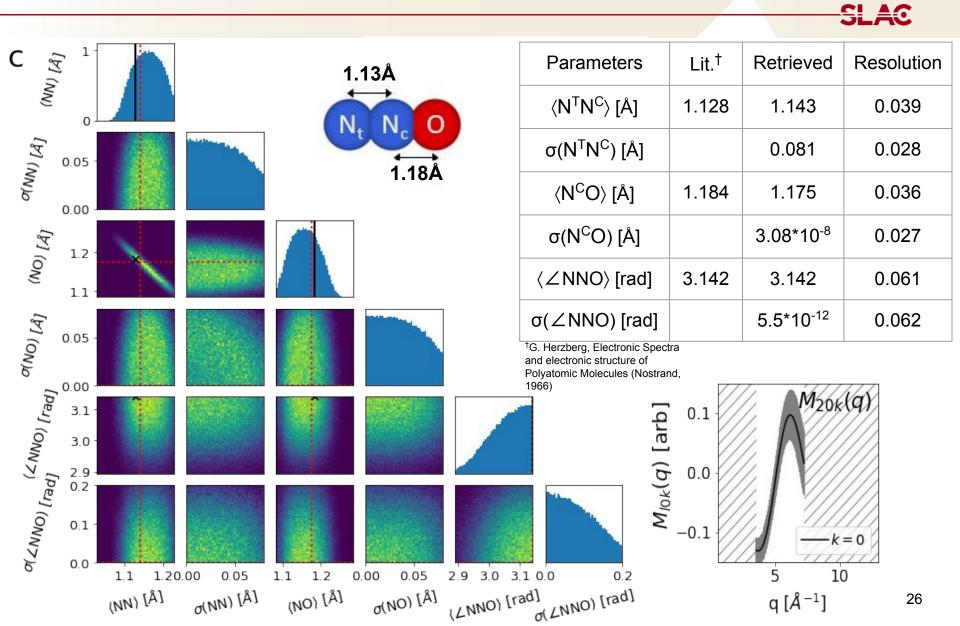


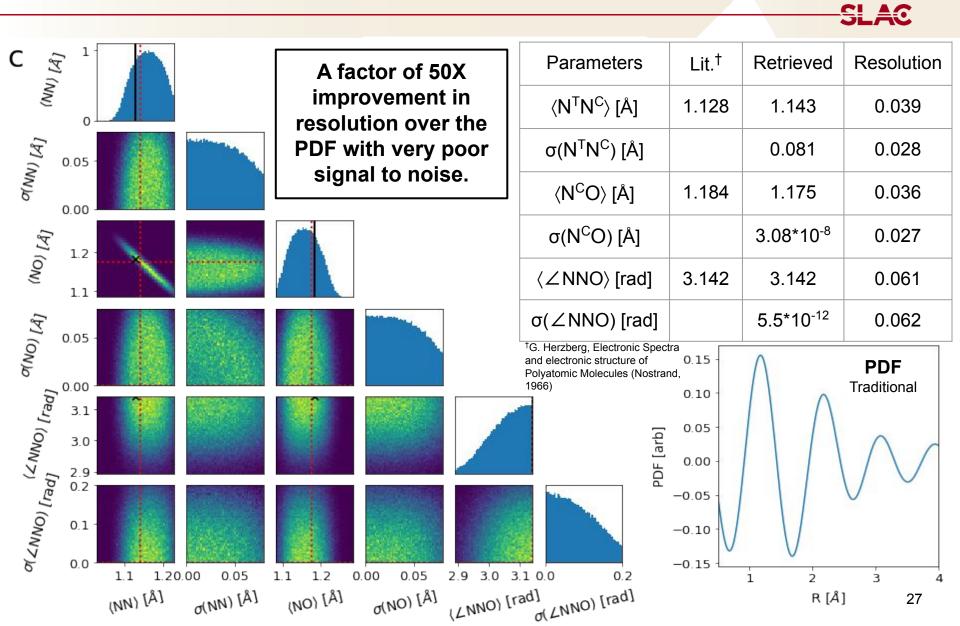


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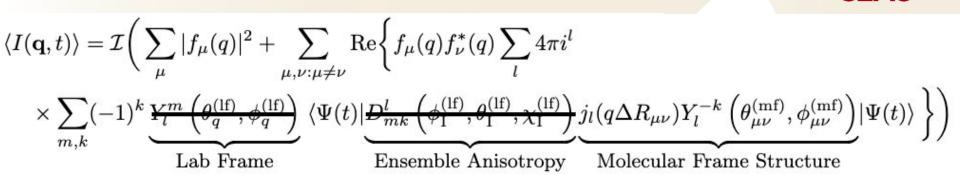








Retrieving Excited State Dynamics (Isotropic)



Use only the isotropic component

- Lose explicit dependence on MF angles
 I=0, m=0, k=0
- Can apply to current datasets
- One may expect similar results
 - Resolution: Order 1 10mÅ

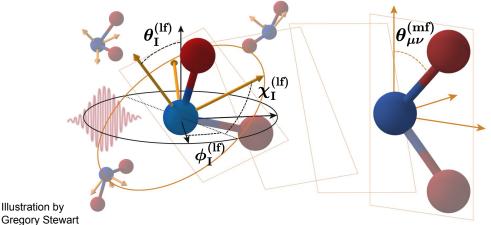
$$C_{lmk}^{(\text{calc})}(q) = \mathcal{I} \sum_{\mu,\nu:\mu\neq\nu} \operatorname{Re}\left\{ f_{\mu}(q) f_{\nu}^{*}(q) (-1)^{k} \frac{32\pi^{3} i^{l}}{2l+1} \times \underbrace{\int j_{l}(q\Delta R_{\mu\nu}) Y_{l}^{-k} \left(\theta_{\mu\nu}^{(\text{mf})}, \phi_{\mu\nu}^{(\text{mf})}\right) P(\mathbf{R}|\Theta, C) d\mathbf{R}}_{\text{Molecular Frame Structure}} \right\}$$

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Retrieving Excited State Dynamics (Anisotropic)

$$\langle I(\mathbf{q}) \rangle_{\text{sep}}^{(2)}(t,\tau) \approx \mathcal{I}\left(\sum_{\mu} |f_{\mu}(q)|^{2} + \sum_{\mu,\nu:\mu\neq\nu} \operatorname{Re}\left\{f_{\mu}(q)f_{\nu}^{*}(q)\sum_{l} \frac{32\pi^{3}i^{l}}{2l+1}\sum_{m_{1},m_{2}}(-1)^{m_{1}}Y_{l}^{m_{2}}\left(\theta_{q}^{(\text{lf})},\phi_{q}^{(\text{lf})}\right) \right. \\ \times \sum_{n,n'} \tilde{\mathcal{A}}_{m_{2}m_{1}}^{(2)l}(n,n'\cdot\tau) \left(\psi_{\text{el-vib}}^{n'}(t)\left|j_{l}(q\Delta R_{\mu\nu})Y_{l}^{-m_{1}}\left(\theta_{\mu\nu}^{(\text{mf})},\phi_{\mu\nu}^{(\text{mf})}\right)\right|\psi_{\text{el-vib}}^{n}(t)\right)\right\} \right).$$
• Separation of rotational and vibrational time scales
• Independent atom approximation
• Measurement: Lab frame anisotropy

- Simulation: Ensemble anisotropy
- Result: Molecular frame geometry
 - Explicit dependence on MF geometric angles
 - Each pair-wise contribution is labelled



a tew to 10 picoseconds and vibrational isomerization or dynamics of interest are often on the femtosecond timescale. sufficient rotation occurs lf outside of the ground state geometry one can use $C_{000}(q,t)$ independent which is of anisotropy but does not have an explicit dependence on the molecular frame angles.

JΥ





- Directly probe $|\Psi(\mathbf{R}^{(mf)})|^2$ in a high dimensional space
 - Retrieve distributions of geometry parameters in high dimensions
 - MHA: efficient, unconstrained, and unbiased search
- Generally applicable to current experiments
 - Choice of model (Normal, ...) allows novel measurements like width
 - Can use induced anisotropy from excitation dipole and C₀₀₀(q,t)
- High resolution (roughly 100-1000 times better than PDF)
 - Simulation resolutions distances/angles: Order 1 0.1 pm / ~1 mrad
 - Data with LMK=[2,0,0] and q=[3.5,7] Å⁻¹ distances/angles: ~4 pm / 6 mrad
 - High signal to noise is more important than q range
- Potentially turn ultrafast gas phase diffraction into a discovery oriented technique without requiring excited state simulations
 - This analysis only requires knowledge of the ground state geometry and ensemble anisotropy simulations which are much more tractible than excited state dynamics









Applying Bayesian Inference and deterministic anisotropy to retrieve the molecular structure $|\Psi(\textbf{R})|^2$ distribution from gas-phase diffraction experiments

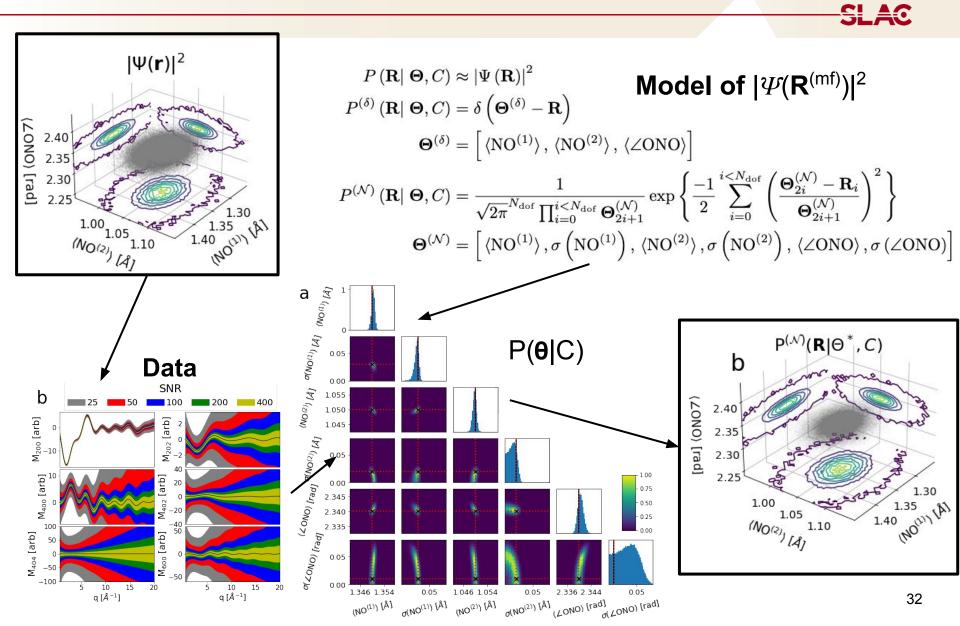
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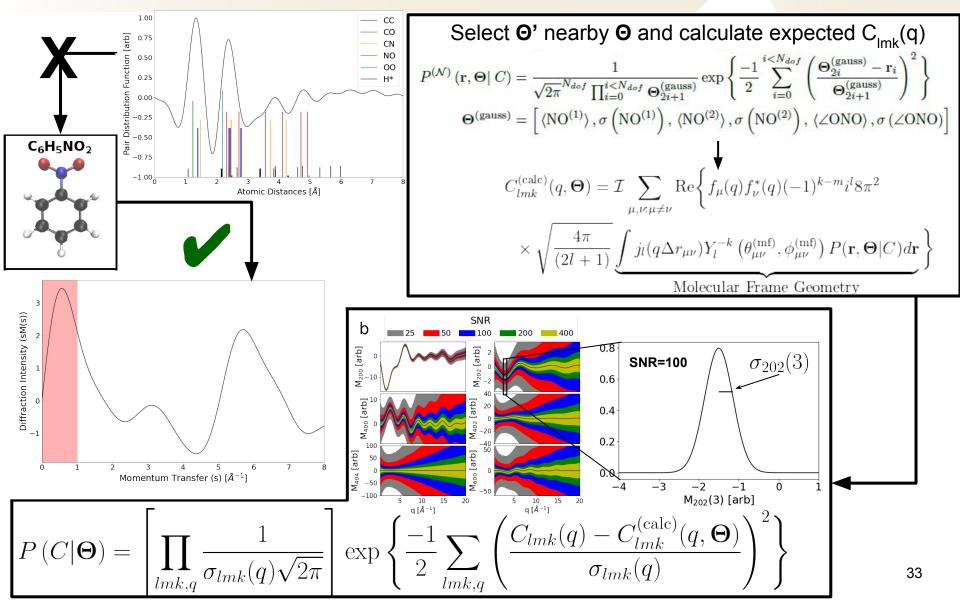
Kareem Hegazy© ^{1,233}, Varun Makhija³, Phil Buckstaum⊚^{1,2,4}, Jeff Corbett⁵, James Cryan², Nick Hartmann⁶, Markus <mark>lichen^{2,7,8}, Keith Jobe⁵,</mark> Renkai Lie[®], Igor Makasyuk⁵, Xiaozhe Shen⊚⁵, Xije Wang⊚⁵, Stephen Weathersby⁵, Jie Yang¹⁰ & Ryan Coffee⊚ ^{2,638}

Backup Slides

Method Overview



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Bayesian Inferencing Reframing the Problem

What we have: Molecular frame $C_{lmk}(q) = \mathcal{I} \sum_{\mu,\nu:\mu\neq\nu} \operatorname{Re} \left\{ f_{\mu}(q) f_{\nu}^{*}(q) (-1)^{k} \frac{32\pi^{3} i^{l}}{2l+1} \right\}$ representations for the nuclear geometry with explicit dependence on pair-wise distances and angles

What we want: Invert the integral equation for $|\Psi(\mathbf{R}^{(mf)})|^2$

How we do it: Approximate $|\Psi(\mathbf{R}^{(mf)})|^2$ with a chosen distribution and solve

for the model pa distribution P(**θ**|

Curse of Dimensionality

 $\times \int j_l(q\Delta R_{\mu\nu})Y_l^{-k}\left(\theta_{\mu\nu}^{(\mathrm{mf})},\phi_{\mu\nu}^{(\mathrm{mf})}\right) \left[\Psi(\mathbf{R},0)\right]^2 d\mathbf{R}$

Molecular Frame Structure

 $\left(\theta_{\mu\nu}^{(\mathrm{mf})}, \phi_{\mu\nu}^{(\mathrm{mf})}\right) P(\mathbf{R}|\mathbf{\Theta}, C) d\mathbf{R}$

When finding **θ** one

can make novel

measurements

r Frame Structure

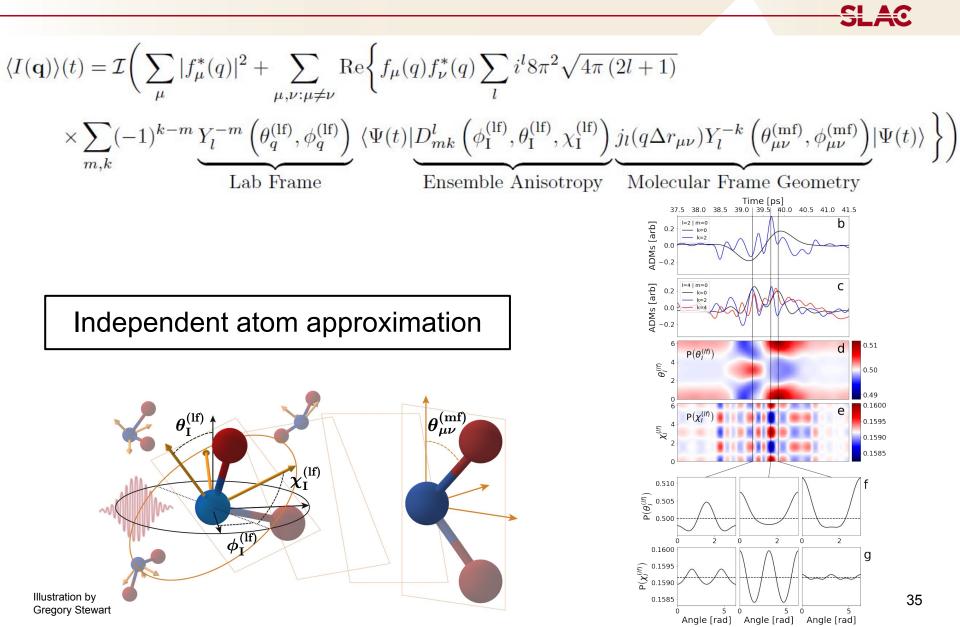
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Bayesian Inference

- Evaluate ~20 Equations
- Order 100 terms
 - Embedded in a k*(3N-6)

Delta Distributio

Normal Distribution: $P^{(\mathcal{N})}(\mathbf{R}|\Theta,C) = \frac{1}{\sqrt{2\pi^{N_{dof}}\prod_{i=0}^{i< N_{dof}}\Theta_{2i+1}^{(\mathcal{N})}}} \exp\left\{\frac{-1}{2}\sum_{i=0}^{i< N_{dof}}\left(\frac{\Theta_{2i}^{(\mathcal{N})} - \mathbf{R}_{i}}{\Theta_{2i+1}^{(\mathcal{N})}}\right)^{2}\right\}$ $\boldsymbol{\Theta}^{\left(\mathcal{N}\right)} = \left[\left< \mathrm{NO}^{\left(1\right)} \right>, \sigma\left(\mathrm{NO}^{\left(1\right)} \right), \left< \mathrm{NO}^{\left(2\right)} \right>, \sigma\left(\mathrm{NO}^{\left(2\right)} \right), \left< \angle \mathrm{ONO} \right>, \sigma\left(\angle \mathrm{ONO} \right) \right]$ 34

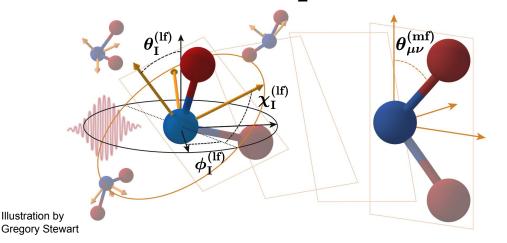


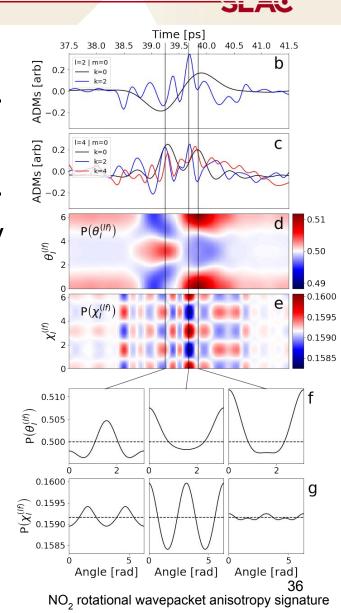
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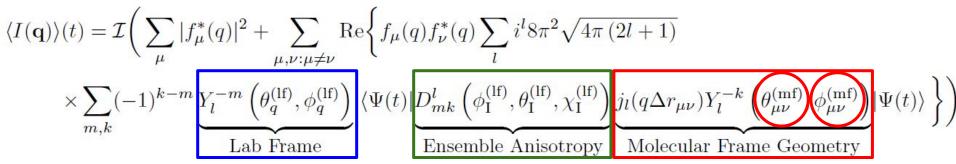
Application in photo-electron spectroscopy

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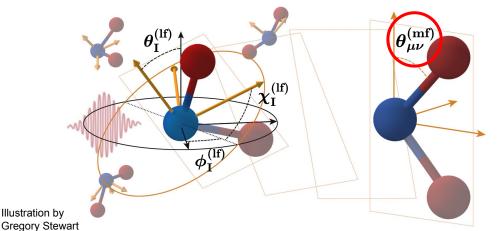




Accessing the MF via Deterministic Anisotropy



- Independent atom approximation
- Measurement: Lab frame anisotropy
- Simulation: Ensemble anisotropy
- Result: Molecular frame geometry
 - Explicit dependence on MF geometric angles
 - Each pair-wise contribution is labelled



Explicit dependence on molecular frame angles provides an opportunity to develop methods that extract these angles, and hence a unique geometry, from data alone.

Retrieving the MF Geometry Probability Distribution Delta Distribution and Systematic Error

Delta Distribution Systematic Error

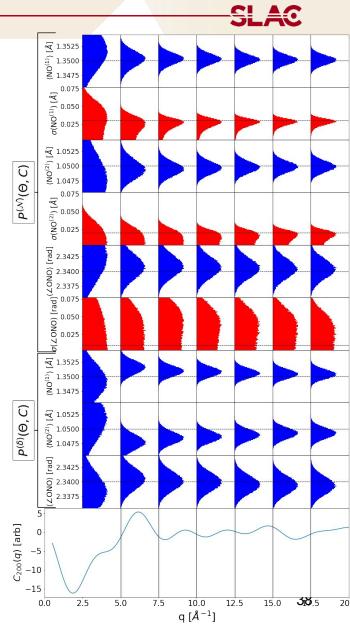
- Assumes signal from ensembles and a single geometry are comparable
- Accuracy is roughly 1000 times worse
- At ≤picometer resolution, the correct value can be 2-3 standard deviations of P(θ|C)

Normal Distribution Mitigates this Systematic

- Normal distribution P(**θ**|C) distributions are closely centered around the expected value
- The 1d mode does not change with q range

Why use the delta distribution?

- It is of order 100 times faster
- Provides sufficient accuracy for debugging analysis



Convergence: Autocorrelation Time (τ)



Criteria for each chain (1000 chains are used)

- At least 100 τ in length
- $\Delta \tau / \tau < 0.01$

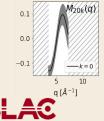
Autocorrelation Time

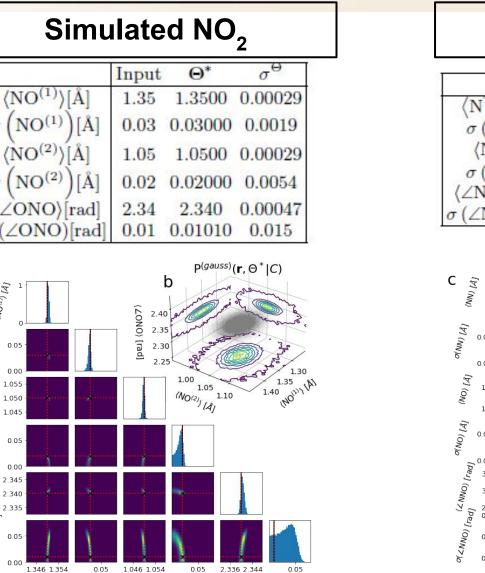
• Number of steps needed for geometries to be uncorrelated

• Expectations:
$$E_{p(\theta)}[f(\theta)] \approx \frac{1}{N} \sum_{n}^{N} f\left(\theta^{(n)}\right)$$
 SEM = $\sqrt{\frac{\operatorname{Var}_{p(\theta)}[f(\theta)]}{N}}$

• Correlation Effects: SEM =
$$\sqrt{\frac{\tau}{N}} \operatorname{Var}_{p(\theta)}[f(\theta)]$$

Retrieving the MF Geometry Probability Distribution Results





σ(NO⁽²⁾⁾ [Å] (LONO) [rad] σ(LONO) [rad]

 σ

 σ

 σ

MOR) [4] MUNIJ [4] (NOU) [4]

(2010) [rad] o(NO(2)) [A]

2.335 (Deal) (ONOZ)o

0.05

0.00

1.055

1.050

1.045

0.05

0.00 2.345 2.340

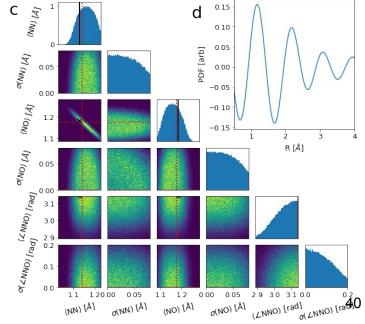
σ(NO⁽¹⁾) [Å]

(NO⁽¹⁾) [Å]

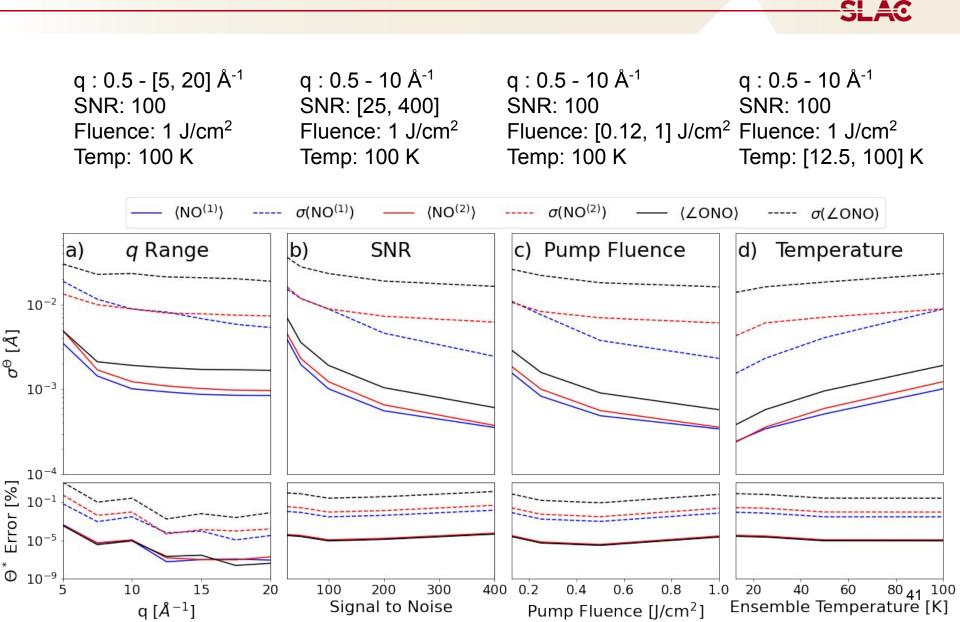
(NO⁽²⁾) [Å]

Measured N₂O

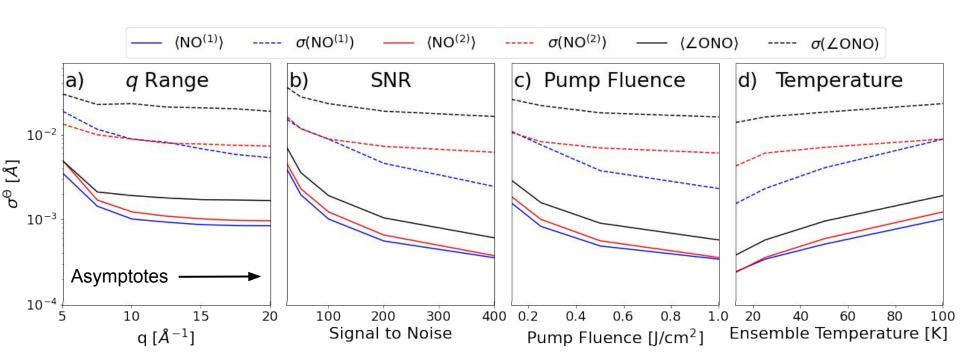
9 04 64 95 6000 10	$\Theta^*_{\text{Literature}}$	Θ^*	σ^{Θ}
$\langle N^T N^C \rangle$ [Å]	1.128	1.142	0.039
σ (NN) [Å]		0.081	0.028
(NO) [Å]	1.184	1.175	0.036
σ (NO) [Å]		3.08×10^{-8}	0.027
(∠NNO) [rad]	3.142	3.142	0.061
σ (\angle NNO) [rad]		5.5×10^{-12}	0.062



Retrieving the MF Geometry Probability Distribution Experimental Parameters

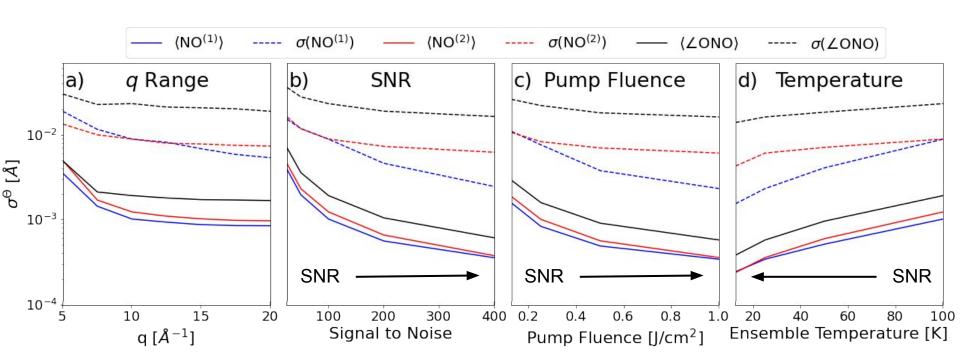


Retrieving the MF Geometry Probability Distribution Experimental Parameters



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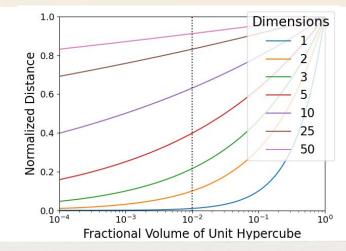
Retrieving the MF Geometry Probability Distribution Experimental Parameters



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Curse of Dimensionality





Assumptions

- 1. Grid Search
- 2. Ignore hydrogens
- 3. Know atom pair-wise distances within 1 Å

Degrees of freedom: $N_{dof} = 3N_{atoms} - 6$ Number of samples: N_s Number of sample per dimension: N_s/D Sampled volume given SS step size: V_s |SS

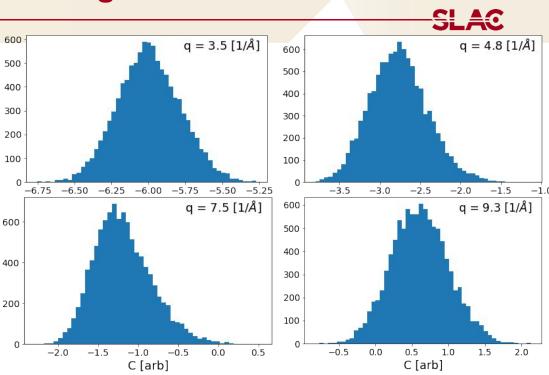
NO ₂			Cyclohexadiene				
0 N	N _{dof}	= 3		L	N N	_{dof} = 12	
Ns	N _s /D	1	N _s	N _s /D	V _s 0.1	V _s 0.25	
10 ⁶	100		10 ⁶	3.16	10 ⁻⁵ %	5.96%	
10 ⁷	215		10 ⁷	3.83	10 ⁻⁴ %	59.6%	
10 ⁸	464		10 ⁸	4.64	10 ⁻³ %	>100%	

 $N_{dof} = 21$ N_s/D N_s V_s|0.1 V_s|0.25 10⁶ 10-14% 2.3×10⁻⁵% 1.93 10-13% 10⁷ 2.3×10⁻⁴% 2.15 10⁸ 10-12% 2.3×10⁻³% 2.40

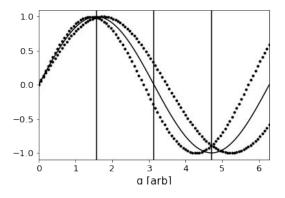
Nitrobenzene

Systematic Errors in Retrieving Geometric Parameters

- Systematic errors are caused by non-gaussian C_{Imk}(q) distributions
- Gaussian distributions of geometries lead to non-gaussian distributions due to sinusoids.

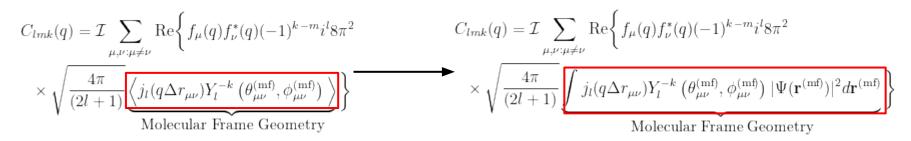


$$P(C|\mathbf{r}^{(mf)}) = \prod_{lmk,q} \frac{1}{\sigma_{lmk}(q)\sqrt{2\pi}} \exp\left\{\frac{-1}{2} \left(\frac{C_{lmk}^{(data)}(q) - C_{lmk}^{(calc)}(q, \mathbf{r}^{(mf)})}{\sigma_{lmk}(q)}\right)^2\right\}$$



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What we have: Molecular frame representations for the nuclear geometry with explicit dependence on pair-wise distances and angles



What we want: Invert the integral equation for $|\Psi(\mathbf{r}^{(mf)})|^2$

How we do it: Approximate $|\Psi(\mathbf{r}^{(mf)})|^2$ with a chosen distribution and solve for the model parameter ($\mathbf{\Theta}$) distribution P($\mathbf{\Theta}|$ C)

$$P(\mathbf{r}, \Theta | C) \approx |\Psi(\mathbf{r})|^{2}$$

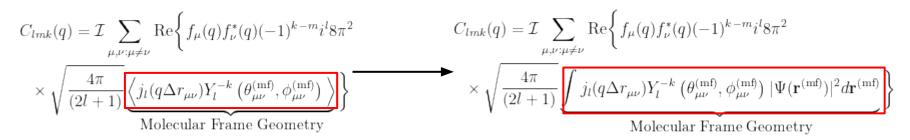
$$P(\Theta | C) = \int P(\mathbf{r}, \Theta | C) d\mathbf{r}$$
Delta Distribution:
$$P^{(\delta)}(\mathbf{r}, \Theta | C) = \delta \left(\Theta^{(\text{delta})} - \mathbf{r} \right)$$

$$\Theta^{(\text{delta})} = \left[\langle \text{NO}^{(1)} \rangle, \langle \text{NO}^{(2)} \rangle, \langle \angle \text{ONO} \rangle \right]$$
Normal Distribution:
$$P^{(\mathcal{N})}(\mathbf{r}, \Theta | C) = \frac{1}{\sqrt{2\pi}^{N_{dof}} \prod_{i=0}^{i < N_{dof}} \Theta_{2i+1}^{(\text{gauss})}} \exp \left\{ \frac{-1}{2} \sum_{i=0}^{i < N_{dof}} \left(\frac{\Theta_{2i}^{(\text{gauss})} - \mathbf{r}_{i}}{\Theta_{2i+1}} \right)^{2} \right\}$$

$$\Theta^{(\text{gauss})} = \left[\langle \text{NO}^{(1)} \rangle, \sigma \left(\text{NO}^{(1)} \right), \langle \text{NO}^{(2)} \rangle, \sigma \left(\text{NO}^{(2)} \right), \langle \angle \text{ONO} \rangle \right]$$
46

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What we have: Molecular frame representations for the nuclear geometry with explicit dependence on pair-wise distances and angles



What we want: Invert the integral equation for $|\Psi(\mathbf{r}^{(mf)})|^2$

How we do it: Approximate $|\Psi(\mathbf{r}^{(mf)})|^2$ with a chosen distribution and solve for the model parameter ($\mathbf{\Theta}$) distribution P($\mathbf{\Theta}|$ C) When finding $\mathbf{\Theta}$ one

$$P(\mathbf{r},\Theta|C) \approx |\Psi(\mathbf{r})|^{2}$$

$$P(\Theta|C) = \int P(\mathbf{r},\Theta|C) d\mathbf{r}$$

$$P^{(\delta)}(\mathbf{r},\Theta|C) = \delta\left(\Theta^{(\text{delta})} - \mathbf{r}\right)$$

$$\Theta^{(\text{delta})} = \left[\langle \text{NO}^{(1)} \rangle, \langle \text{NO}^{(2)} \rangle, \langle \angle \text{ONO} \rangle\right]$$

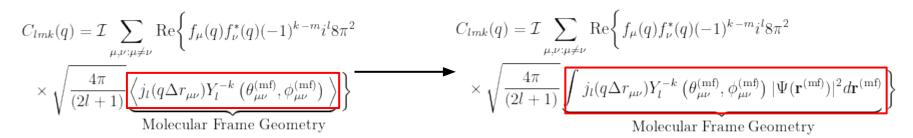
$$\text{Normal Distribution:} P^{(\mathcal{N})}(\mathbf{r},\Theta|C) = \frac{1}{\sqrt{2\pi}^{N_{dof}} \prod_{i=0}^{i < N_{dof}} \Theta^{(\text{gauss})}_{2i+1}} \exp\left\{\frac{-1}{2}\sum_{i=0}^{i < N_{dof}} \left(\frac{\Theta^{(\text{gauss})}_{2i} - \mathbf{r}_{i}}{\Theta^{(\text{gauss})}_{2i+1}}\right)^{2}\right\}$$

$$\Theta^{(\text{gauss})} = \left[\langle \text{NO}^{(1)} \rangle, \sigma\left(\text{NO}^{(1)}\right), \langle \text{NO}^{(2)} \rangle, \sigma\left(\text{NO}^{(2)}\right), \langle \angle \text{ONO} \rangle\right]$$

$$47$$

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What we have: Molecular frame representations for the nuclear geometry with explicit dependence on pair-wise distances and angles

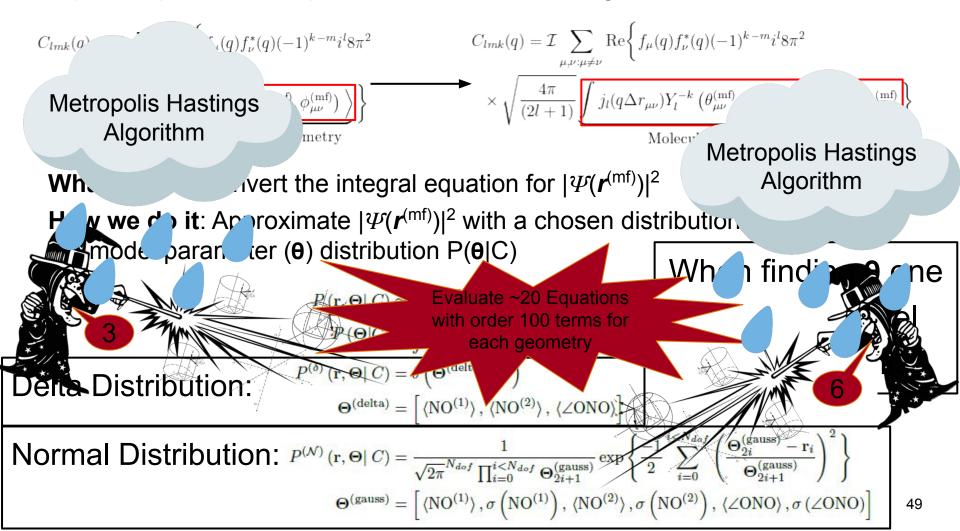


What we want: Invert the integral equation for $|\Psi(\mathbf{r}^{(mf)})|^2$

How we do it: Approximate $|\Psi(\mathbf{r}^{(mf)})|^2$ with a chosen distribution and solve for model parameter (**θ**) distribution P(**θ**|C) When finding one $\mathbb{R}(\mathbf{r},\Theta) \longrightarrow \mathbb{E}$ Evaluate ~20 Equations with order 100 terms for PIOID each geometry $(\mathbf{r}, \Theta | C) = \mathbf{O}^{(\text{delt})}$ Délta Distribution: $\Theta^{(\text{delta})} = \left[\langle \text{NO}^{(1)} \rangle, \langle \text{NO}^{(2)} \rangle, \langle \angle \text{ONO} \rangle \right]$ $\Theta_{2i}^{(gauss)}$ Normal Distribution: $P^{(\mathcal{N})}(\mathbf{r}, \Theta | C) = \frac{1}{\sqrt{2\pi^{N_{dof}} \prod_{i=0}^{i < N_{dof}} \Theta_{2i+1}^{(gauss)}}} exp$ $\Theta_{2i+1}^{(\text{gauss})}$ $\boldsymbol{\Theta}^{(\text{gauss})} = \left[\left< \text{NO}^{(1)} \right>, \sigma \left(\text{NO}^{(1)} \right), \left< \text{NO}^{(2)} \right>, \sigma \left(\text{NO}^{(2)} \right), \left< \angle \text{ONO} \right>, \sigma \left(\angle \text{ONO} \right) \right] \right]$ 48

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What we have: Molecular frame representations for the nuclear geometry with explicit dependence on pair-wise distances and angles



Bayesian Inferencing Metropolis Hastings Algorithm

- Inverts the system of equations to solve for the joint P(θ|C) distribution
- Unbiased sampling method designed for high dimensional spaces

$$C_{lmk}(q) = \mathcal{I} \sum_{\mu,\nu:\mu\neq\nu} \operatorname{Re} \left\{ f_{\mu}(q) f_{\nu}^{*}(q) (-1)^{k-m} i^{l} 8\pi^{2} \\ \times \sqrt{\frac{4\pi}{(2l+1)}} \underbrace{\int j_{l}(q\Delta r_{\mu\nu}) Y_{l}^{-k} \left(\theta_{\mu\nu}^{(\mathrm{mf})}, \phi_{\mu\nu}^{(\mathrm{mf})}\right) |\Psi(\mathbf{r})|^{2} d\mathbf{r}} \right\} \\ \text{Molecular Frame Geometry} \\ \mathbf{Bayesian Inferencing} \\ C_{lmk}^{(\mathrm{calc})}(q, \mathbf{\Theta}) = \mathcal{I} \sum_{\mu,\nu:\mu\neq\nu} \operatorname{Re} \left\{ f_{\mu}(q) f_{\nu}^{*}(q) (-1)^{k-m} i^{l} 8\pi^{2} / (2l+1) \int j_{l}(q\Delta r_{\mu\nu}) Y_{l}^{-k} \left(\theta_{\mu\nu}^{(\mathrm{mf})}, \phi_{\mu\nu}^{(\mathrm{mf})}\right) \underbrace{P(\mathbf{r}, \mathbf{\Theta}|C)} d\mathbf{r} \right\}$$

Molecular Frame Geometry

Method

- 1. Select neighboring **0**
- 2. Add new $\boldsymbol{\theta}$ to P($\boldsymbol{\theta}|C$) with probability P(C| $\boldsymbol{\theta}$)_{New}/P(C| $\boldsymbol{\theta}$)_{Prev}

$$P(C|\Theta) = e^{L(\Theta)} \left[\prod_{lmk,q} \frac{1}{\sigma_{lmk}(q)\sqrt{2\pi}} \right] \exp\left\{ \frac{-1}{2} \sum_{lmk,q} \left(\frac{C_{lmk}(q) - C_{lmk}^{(\text{calc})}(q,\Theta)}{\sigma_{lmk}(q)} \right)^2 \right\}$$

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