

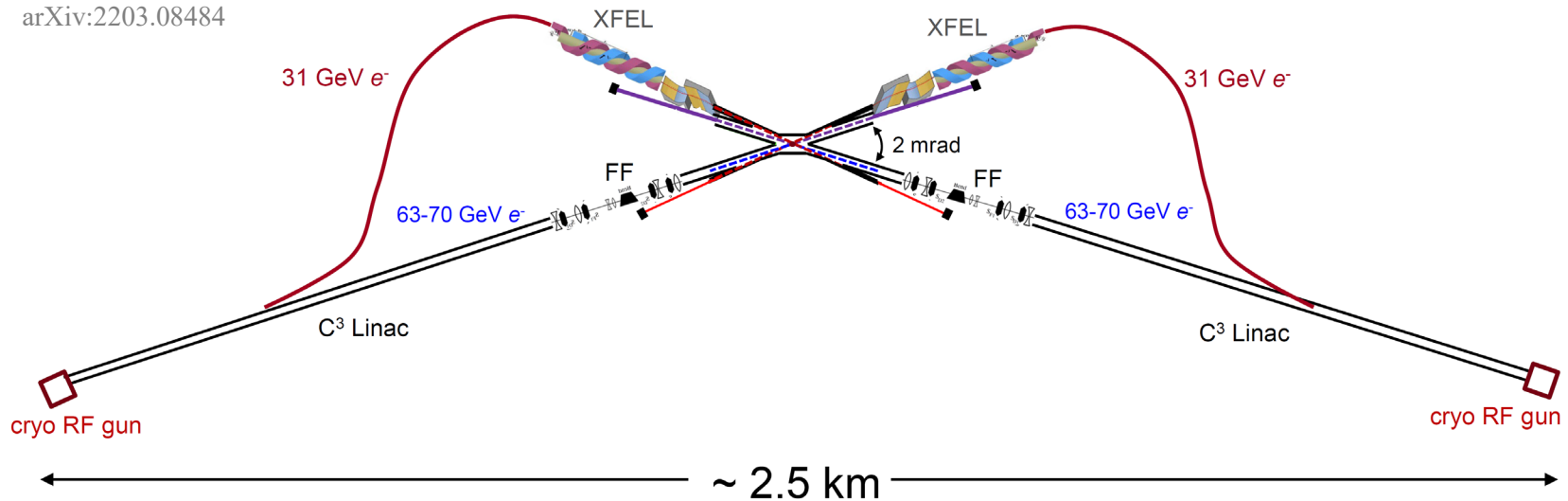
XCC Physics Case

Tim Barklow
SLACmass Summary Retreat
May 12, 2022



XCC – XFEL Compton Collider

arXiv:2203.08484



Run $\gamma\gamma \rightarrow H$ at $\sqrt{s_{\gamma\gamma}} = 125$ GeV 30% of the time
 and $e^- \gamma \rightarrow e^- H$ at $\sqrt{s_{e\gamma}} = 140$ GeV 70% of the time
 to calibrate the $\sigma \times BR$ measurements at $\sqrt{s_{\gamma\gamma}} = 125$ GeV.
 This produces model independent Higgs coupling
 measurements, just like the ILC.

$$\sigma_{ez} = 20 \mu\text{m} \quad \sigma_{\gamma z} = 20 \mu\text{m}$$

$$N_{e^-} = 1 \text{ nC} \quad \gamma\epsilon_{x,y} = 120 \text{ nm}$$

$$d_{cp} = 60 \mu\text{m} \quad P_{e^-} = 90\%$$

$$a_{\gamma FWHM} = 70 \text{ nm} \quad \xi_{\text{non-linear QED}}^2 = 0.10$$

$\gamma\gamma$ mode $\sqrt{s} = 125$ GeV		
Process	Luminosity ($10^{34} \text{ cm}^{-2} \text{ s}^{-1}$)	
	Total	$\sqrt{s} > 100$ GeV
$\gamma\gamma$	2.1	0.12
$e^- e^-$	0.23	0.18
$e^- \gamma + \gamma e^-$	2.5	0.42
$e^+ e^- + e^- e^+$	0.48	0.05
$e^+ \gamma + \gamma e^+$	0.47	0.01

The XCC is presented as a possible lower cost alternative to the ILC and C³ 250 GeV e^+e^- Higgs factories. It is being pursued because every e^+e^- linear collider proposal to date has been rejected due to its high cost. That said, it should be noted that strong synergies between XCC and the SLAC XFEL program also serve to motivate this concept.

Potential Cost Savings with the XCC

C³ 250 GeV Capital Cost Estimate

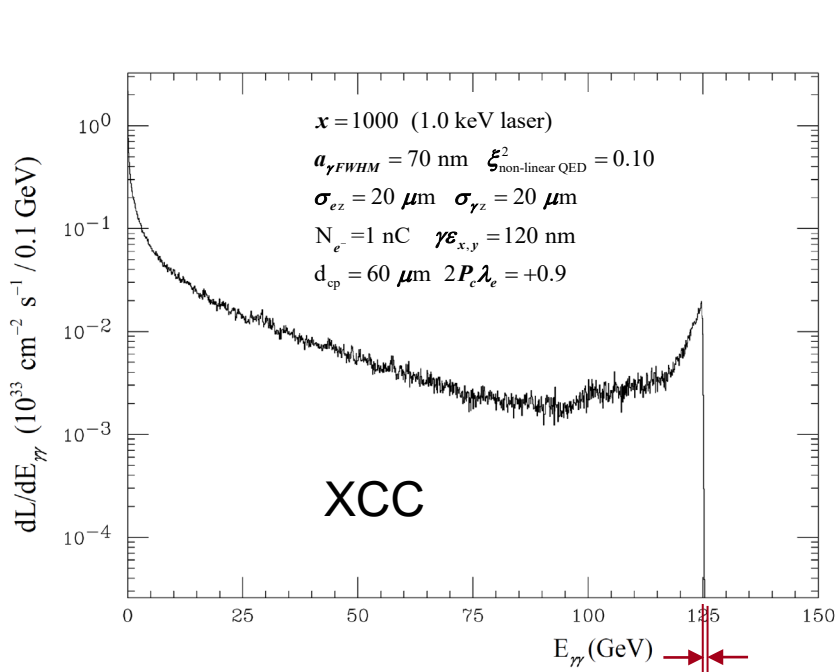
CCC	GeV	250		
	MeV/m	70		
	Sub-Domain	M\$	%	%
Sources	Injectors	301	8	35
	Damping Rings	461	12	
	Beam Transport	563	15	
Main Linac	Cryomodule	357	10	33
	C-band Klystron	871	23	
IP	Beam Delivery and FF	295	8	13
	IR	184	5	
Support Inf.	Civil Eng	204	5	19
	Common Facilities	396	11	
	Cryo-plant	101	3	
	Total	3733	100	

XCC 140 GeV Capital Cost Estimate

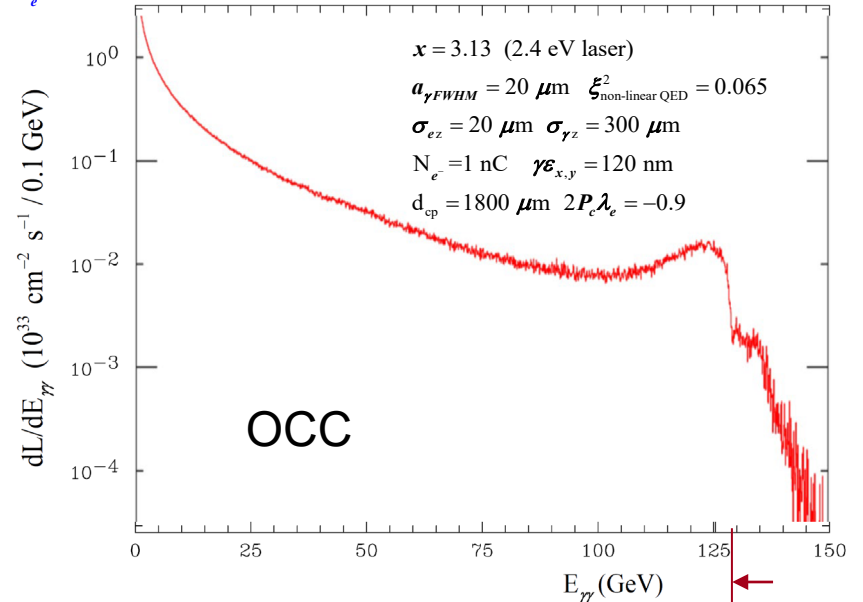
XCC	GeV	140		
	MeV/m	70		
	Sub-Domain	M\$	%	%
Sources	Injectors	200	9	26
	FEL	200	9	
	Beam Transport	197	9	
Main Linac	Cryomodule	200	9	30
	C-band Klystron	488	22	
IP	Beam Delivery and FF	148	7	15
	IR	184	8	
Support Inf.	Civil Eng	114	5	28
	Common Facilities	396	18	
	Cryo-plant	133	6	
	Total	2260	100	

With these estimates the XCC would be 60% of the cost of C³ 250 GeV. Given the very early stage of the XCC design and the many XFEL technical challenges, it is important that these tables are viewed as illustrative, providing insight into the *potential* cost savings of the XCC.

Higgs Rate and Background for XCC vs Optical $\gamma\gamma$ Collider (OCC) & ILC



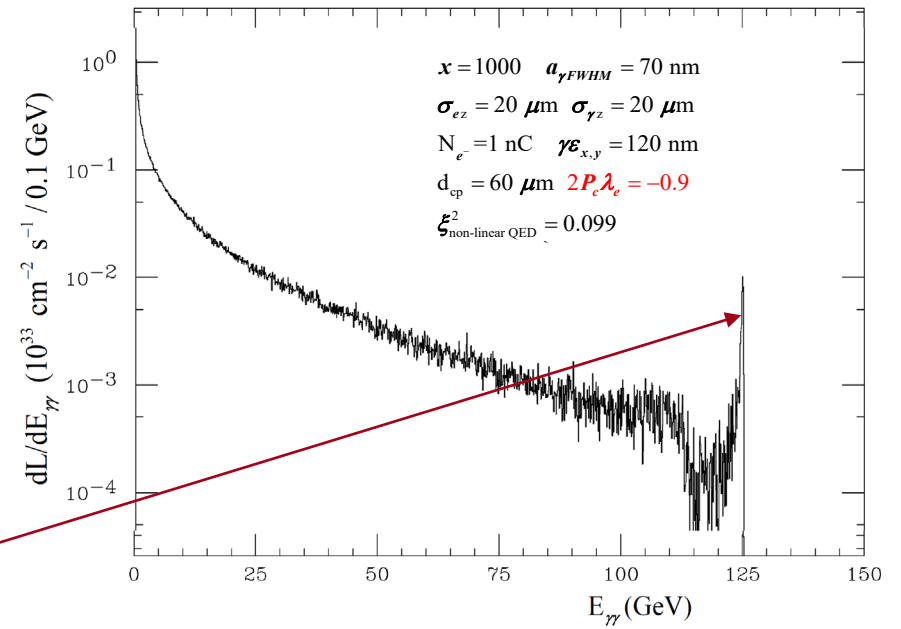
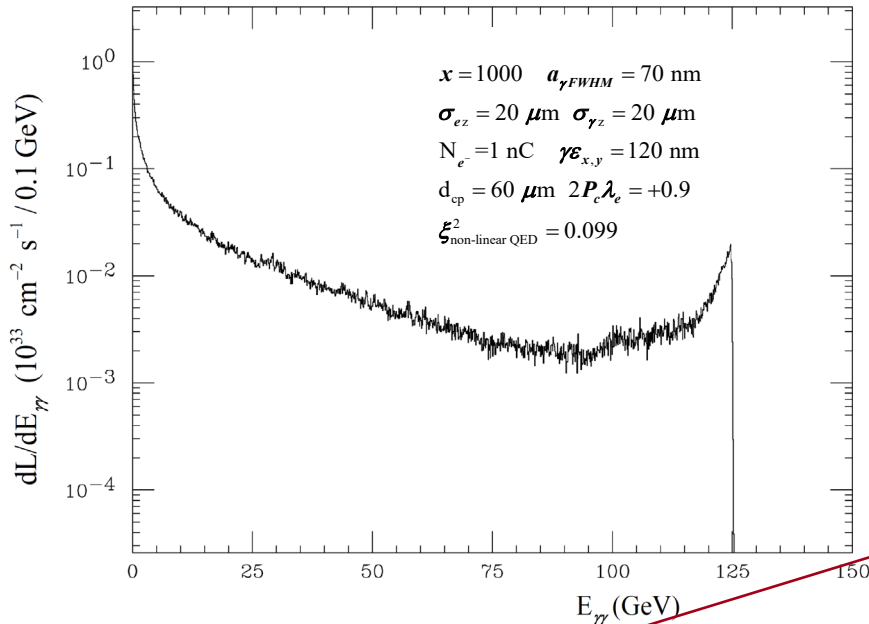
$$2 \max(E_\gamma^{\text{non-linear QED}}) - 2 \max(E_\gamma^{\text{linear}}) = \frac{2E_e}{x+1} = 0.13 \text{ GeV}$$



$$2 \max(E_\gamma^{\text{non-linear QED}}) - 2 \max(E_\gamma^{\text{linear}}) = \frac{2E_e}{x+1} = 42 \text{ GeV}$$

Machine	E_{e^-} (GeV)	N_{e^-} (nC)	Polarization	N_H/yr	N_{Hadronic}/N_H	$N_{\text{minbias}}/\text{BX}$
XCC	62.8	1.0	90% e^-	34,000	170	9.5
OCC	86.5	1.0	90% e^-	30,000	540	50
ILC	125	3.2	-80% e^- +30% e^+	42,000	140	1.3
ILC	125	3.2	+80% e^- -30% e^+	28,000	60	1.3

Alternative Polarization for Scanning Higgs Resonance



Narrower leading edge width (45 MeV) at cost of lower Higgs rate. Width dominated by 0.05% e^- energy spread.

A Higgs resonance scan can measure the total width to 4.5 MeV, i.e., a 112% measurement if the width has the SM value.

A total Higgs width $\gg 4$ MeV is not ruled out by LHC data, but would require, for example, a conspiracy of a universal

kappa scale factor $\kappa_0 > 1$ combined with $B_{BSM} \approx 1 - \kappa_0^{-2}$ assuming LHC kappa ratios $\rightarrow 1$ with ever greater precision:

κ_0	B_{BSM} (%)	Γ_H (MeV)
1.4	49	15
1.2	31	8
1.0	0	4

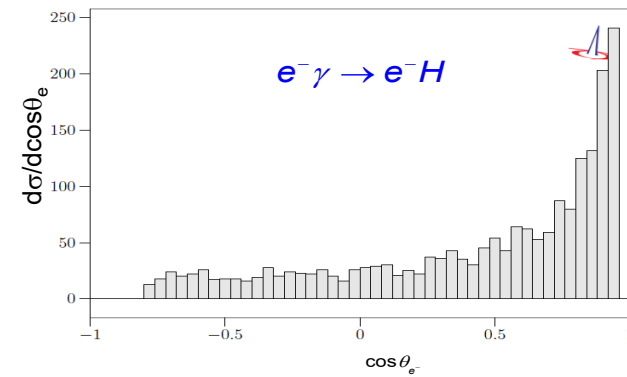
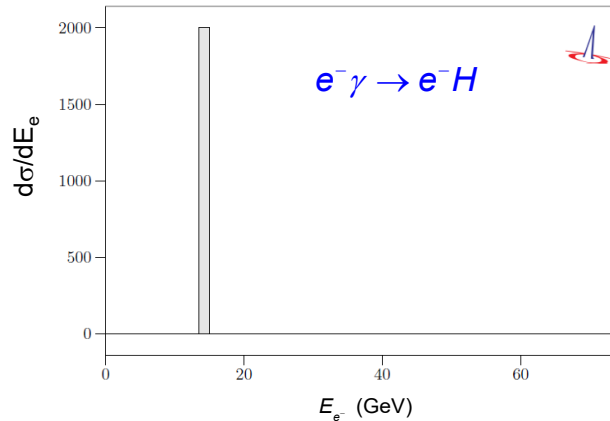
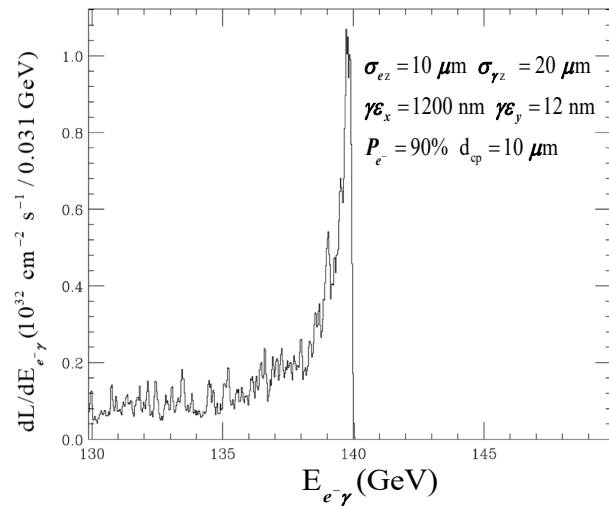
$$\Leftarrow \Gamma_H \text{ assuming universal } \kappa_0, B_{BSM} = 1 - \kappa_0^{-2}, \frac{\kappa_g \kappa_Z}{\kappa_H} = \frac{\kappa_W}{\kappa_Z} = \frac{\kappa_\gamma}{\kappa_Z} = \frac{\kappa_g}{\kappa_Z} = \dots = 1$$

(also LHC off-shell $H^* \rightarrow ZZ$ measurements would have to be addressed)

Another possible application of the alternative polarization is improved signal-to-background. Detailed studies are required to determine if this improvement can compensate for the loss in signal statistics.

Measurement of Γ_γ using $e^- \gamma \rightarrow e^- H$ at $E_{\text{cm}}=140$ GeV

If, as is likely, a direct 5 MeV measurement of the Higgs width corresponds to a large fractional error, then individual Higgs partial widths and the total Higgs width will have to be extracted at XCC by measuring Γ_γ through $e^- \gamma \rightarrow e^- H$ at $\sqrt{s} = 140$ GeV. This is the XCC analog of the e^+e^- Higgs factory measurement of Γ_Z through Higgs recoil in $e^+e^- \rightarrow ZH$.



The signal is a monochromatic 14.2 GeV electron, predominantly in the forward direction. In order to achieve model independent ILC-like precision for Higgs couplings and the total Higgs width, about 1 $e^- \gamma \rightarrow e^- H$ event must be detected at 140 GeV per 125 $\gamma\gamma \rightarrow H$ events collected at 125 GeV. $\sigma(e^- \gamma \rightarrow e^- H) = 4.1$ fb at 140 GeV assuming forward detector coverage down to $\theta > 3$ mrad (there is no Compton scatter background on this side of the IP in $e^- \gamma$ collisions).

With the current $e^- \gamma$ collider design, the yearly luminosity with $\sqrt{\hat{s}}$ within 1% of the 140 GeV peak is 32 fb^{-1} \Rightarrow for every year collecting Higgs events at $\sqrt{s} = 125$ GeV, two years must be spent producing $e^- \gamma \rightarrow e^- H$ at 140 GeV.

XCC Coupling Errors Using EFT Higgs Program

coupling a	ILC Δa (%)	XCC Δa (%)
HZZ	0.57	1.2
HWW	0.55	1.2
Hbb	1.0	1.4
$H\tau\tau$	1.2	1.4
Hgg	1.6	1.7
Hcc	1.8	1.8
$H\gamma\gamma$	1.1	0.77
$H\gamma Z$	9.1	10.0
$H\mu\mu$	4.0	3.8
Γ_{tot}	2.4	3.8
$\Gamma_{\text{inv}}^\dagger$	0.36	—
$\Gamma_{\text{other}}^\dagger$	1.6	2.7

† 95% C.L. limit

ILC: $0.5 \times 10^6 e^+e^- \rightarrow ZH$ events
 full $2 \text{ ab}^{-1} \sqrt{s} = 250 \text{ GeV}$
 10 year program

XCC: $0.5 \times 10^6 \gamma\gamma \rightarrow H$ events
 4000 $e^- \gamma \rightarrow e^- H$ events

4 years $\gamma\gamma \rightarrow H @ \sqrt{s} = 125 \text{ GeV}$

8 years $e^- \gamma \rightarrow e^- H @ \sqrt{s} = 140 \text{ GeV}$
 assuming $n_{\text{bunch}} = 76 \rightarrow 290$

Use ILC σ_{XBR} measurement errors
 for XCC:

-80% e^- , +30% e^+ polarization:						
	250 GeV		350 GeV		500 GeV	
	Zh	$\nu\bar{\nu}h$	Zh	$\nu\bar{\nu}h$	Zh	$\nu\bar{\nu}h$
σ	2.0		1.8		4.2	
$h \rightarrow \text{invis.}$	0.86		1.4		3.4	
$h \rightarrow b\bar{b}$	1.3	8.1	1.5	1.8	2.5	0.93
$h \rightarrow c\bar{c}$	8.3		11	19	18	8.8
$h \rightarrow g\bar{g}$	7.0		8.4	7.7	15	5.8
$h \rightarrow WW$	4.6		5.6*	5.7*	7.7	3.4
$h \rightarrow \tau\tau$	3.2		4.0*	16*	6.1	9.8
$h \rightarrow ZZ$	18		25*	20*	35*	12*
$h \rightarrow \gamma\gamma$	34*		39*	45*	47	27
$h \rightarrow \mu\mu$	72		87*	160*	120	100
a	7.6		2.7*		4.0	
b	2.7		0.69*		0.70	
$\rho(a,b)$	-99.17		-95.6*		-84.8	

The $e^- \gamma$ Luminosity Problem

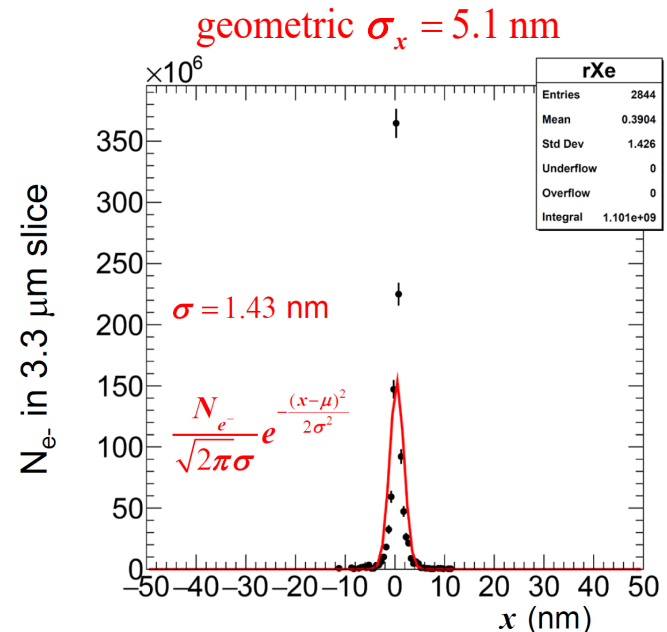
With the current $e^- \gamma$ collider design, the Higgs rate in $e^- \gamma$ collisions at $\sqrt{s} = 140$ GeV is 0.8% of the rate in $\gamma \gamma$ collisions at $\sqrt{s} = 125$ GeV. This is an unsatisfactory situation as 2/3 of the running time is spent waiting for $e^- \gamma \rightarrow e^- H$ events to dribble in at $\sqrt{s} = 140$ GeV. Another related issue is the factor of 3.8 increase in the number of bunches per train required to achieve ILC-like Higgs precision over 12 years. Only a factor of 2 increase in the number of bunches per train at XCC is required to match the ILC's count of 0.5×10^6 Higgs bosons over a decade (note that ILC also assumes a 2x luminosity upgrade).

The $e^- e^-$ geometric luminosity for $e^- \gamma$ collisions is $11 \times 10^{34} \text{ cm}^{-2} \text{ s}^{-1}$ and yet the $e^- \gamma$ luminosity within 1% of the 140 GeV peak is only $0.09 \times 10^{34} \text{ cm}^{-2} \text{ s}^{-1}$ for symmetric x and y emittiances.

Large coherent $e^+ e^-$ pair production leads to pinching of the opposite e^- beam which further increases the E_{field} leading to more positron production and pinching in a feedback manner (new effect discovered in XCC study).

$$\gamma \epsilon_x = 120 \text{ nm} \quad \gamma \epsilon_y = 120 \text{ nm}$$

$e^- \gamma$ mode $\sqrt{s} = 140$ GeV		
Process	Luminosity ($10^{34} \text{ cm}^{-2} \text{ s}^{-1}$)	
	Total	$\sqrt{s} > 139$ GeV
$e^- \gamma$	139	0.09
$\gamma \gamma$	292	-
$e^- e^+$	173	-
γe^+	171	-
γe^-	4.5	-
$e^- e^-$	0.9	-
$e^+ \gamma$	3.0	-
$e^+ e^-$	0.06	-



The $e^- \gamma$ Luminosity Problem

Solution for now is to go to asymmetric emittances

$$\gamma \epsilon_x = 120 \text{ nm} \quad \gamma \epsilon_y = 120 \text{ nm}$$

$e^- \gamma$ mode $\sqrt{s} = 140 \text{ GeV}$		
Process	Luminosity ($10^{34} \text{ cm}^{-2} \text{ s}^{-1}$)	
	Total	$\sqrt{\hat{s}} > 139 \text{ GeV}$
$e^- \gamma$	139	0.09
$\gamma \gamma$	292	-
$e^- e^+$	173	-
γe^+	171	-
γe^-	4.5	-
$e^- e^-$	0.9	-
$e^+ \gamma$	3.0	-
$e^+ e^-$	0.06	-



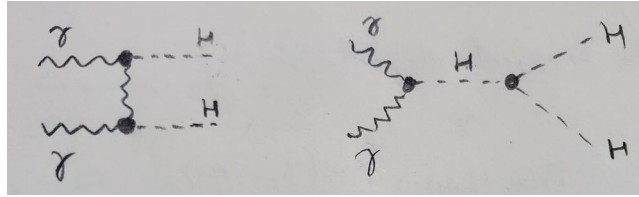
$$\gamma \epsilon_x = 1200 \text{ nm} \quad \gamma \epsilon_y = 12 \text{ nm}$$

$e^- \gamma$ mode $\sqrt{s} = 140 \text{ GeV}$		
Process	Luminosity ($10^{34} \text{ cm}^{-2} \text{ s}^{-1}$)	
	Total	$\sqrt{\hat{s}} > 139 \text{ GeV}$
$e^- \gamma$	11.5	0.32
$\gamma \gamma$	14.5	-
$e^- e^+$	13.4	-
γe^+	11.3	-
γe^-	0.92	-
$e^- e^-$	0.43	0.07
$e^+ \gamma$	0.09	-
$e^+ e^-$	0.01	-

Ultimately we would like to suppress the E_{field}

- Introduce a plasma to neutralize the IP (suggestion by F. Zimmerman to reduce the anti-pinch in $e^- e^-$ collisions)
- Studies using CAIN indicate that the introduction of an additional 10 GeV e^- beam with suitable timing and location could deflect the Compton-scattered beam just enough to significantly suppress beamstrahlung and coherent $e^+ e^-$ pair-production.

Energy upgrade to $E_{cm}=280$ GeV for Higgs Self Coupling Study



2.8 km footprint

assuming gradient 70 MeV/m \rightarrow 120 MeV/m
(C³ uses assumes this gradient upgrade to get from $E_{cm}=250 \rightarrow 500$ GeV)

2012 Study

A feasibility study of the measurement of Higgs pair creation at a Photon Linear Collider

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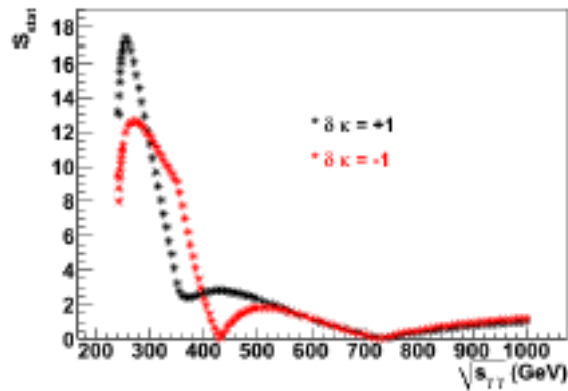
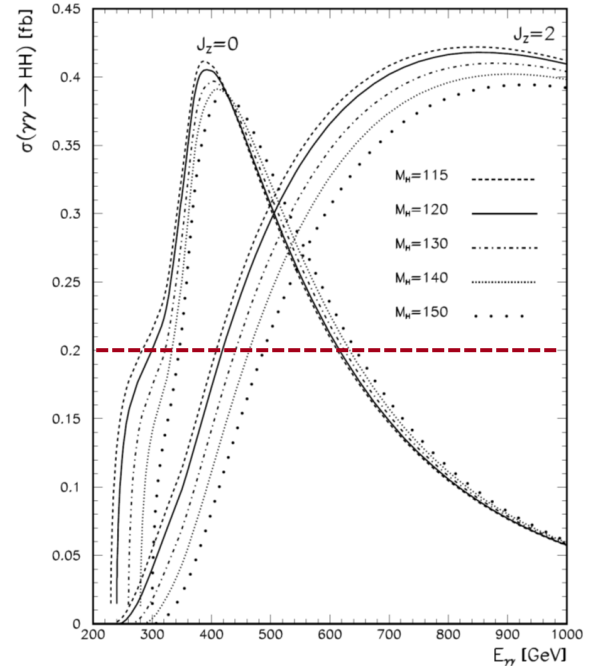


FIG. 3. Statistical sensitivity (S_{stat}) as a function of $\gamma\gamma$ collision energy. Black and red dots show the $\delta\kappa = +1$ and $\delta\kappa = -1$ cases.



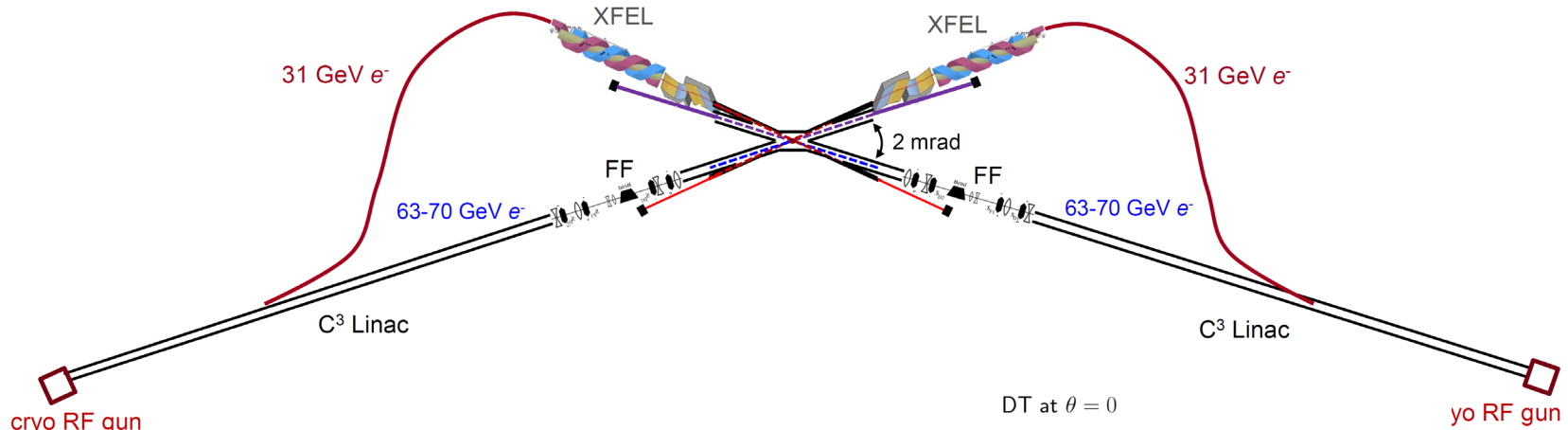
$$\sigma(\gamma\gamma \rightarrow HH) @ \sqrt{s} = 280 \text{ GeV} \approx \sigma(e^+e^- \rightarrow ZHH) @ \sqrt{s} = 500 \text{ GeV}$$

Need to redo the KEK $\gamma\gamma \rightarrow HH$ study with the XCC $\gamma\gamma$ spectrum.

Optimum sensitivity at $\sqrt{s_{\gamma\gamma}} = 280$ GeV

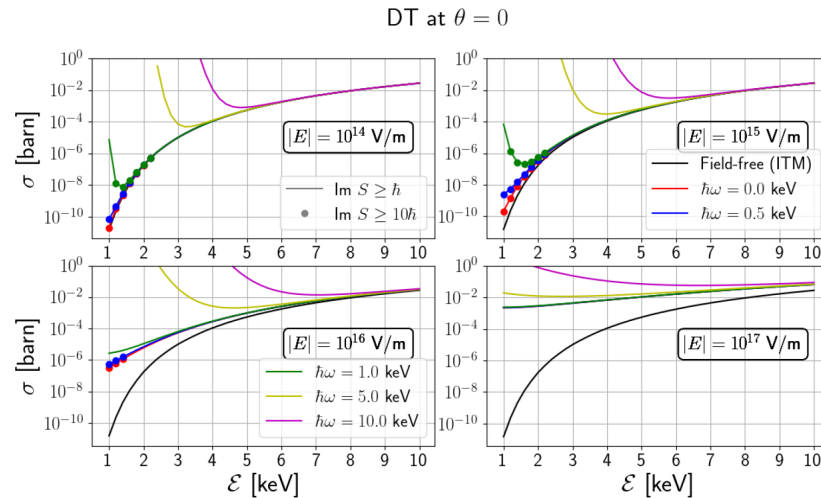
Used optical laser for Compton scattering

Photon Science at XCC - HEDS Fusion Example



$\sigma(^2\text{H}^3\text{H} \rightarrow ^4\text{H}_e\text{n})$ enhanced in presence of $E_{\text{field}} > 10^{16}$ V/m

“Applicability of semiclassical methods for modeling laser-enhanced fusion rates in a realistic setting”
Phys. Rev. C **105**, 054001 – 10 May 2022



$$E_{\text{field}} \propto E_{\text{pulse}}^{\frac{1}{2}} a_{\gamma FWHM}^{-1} \sigma_{\gamma z}^{-\frac{1}{2}} \quad E_{\text{pulse}} = 0.7 \text{ J}$$

$2\sigma_{\gamma z}$ (μm)	$2\sigma_{\gamma z}$ (fs)	$a_{\gamma FWHM}$ (nm)	E_{field} (V/m)	I_{peak} (W cm^{-2})	P (TW)	Validity
10	33	71	1.6×10^{15}	7.1×10^{23}	21	Soft X-rays
10	33	12	1.0×10^{16}	2.8×10^{25}	21	Hard X-rays

XCC Physics Summary

- The XCC at $E_{\text{cm}}=125\text{-}140$ GeV can measure absolute Higgs couplings in a model independent manner with an accuracy of order 1% . This is pretty close to the ILC precision. To fully match or exceed the ILC Higgs coupling accuracy, a way must be found to increase the top 1% $e\text{-}\gamma$ luminosity at $E_{\text{cm}}=140$ GeV.
- The Higgs self coupling can be studied via $\gamma\gamma\rightarrow\text{HH}$ if the XCC energy is upgraded to $E_{\text{cm}}=280$ GeV. Given that $\sigma(\gamma\gamma\rightarrow\text{HH}) \sim \sigma(e^+e^- \rightarrow\text{ZH})$, the Higgs self coupling sensitivity for XCC will probably be comparable to ILC at $E_{\text{cm}}=550$ GeV
- There are strong synergies between XCC and the XFEL program at SLAC. Solutions to high energy/pulse XFEL production and focusing issues at XCC will lead to new opportunities in XFEL photon science.