

Palo Alto, 13 – 16 March 2023



Overview of Radiation Effects on LGADs

Evangelos – Leonidas Gkougkousis

~~(CERN)~~

SLAC – March 13th, 2023

•Introduction I

LGAD Technology

BROOKHAVEN
NATIONAL LABORATORY

NDL



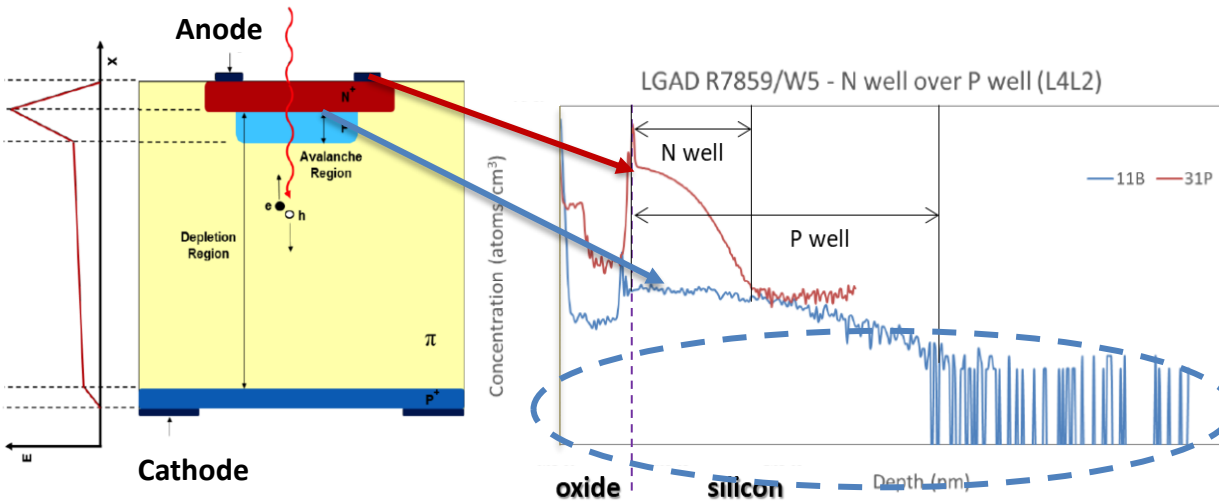
CNM
Centre Nacional de Microelectrónica



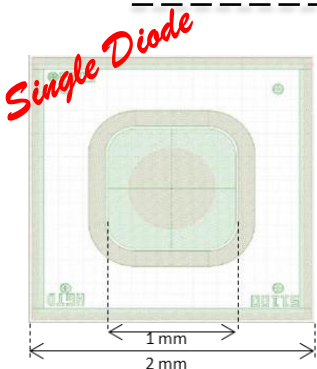
Teledyne e2v

HAMAMATSU
PHOTON IS OUR BUSINESS

- ✓ Invented at CNM, initially considered for tracking by IFAE, proposed for timing by UCSC
- ✓ HPK, CNM, FBK, MiCRON, BNL (USA), NDL (China), CiS, Teledyne (UK)



- ✓ Requires precise diffusion control for layer thickness:
 - ✓ Thin highly doped n-well layer (~ 1 – 1.5 μm)
 - ✓ Gain layer ~ 2 μm
 - ✓ p-stop ~ 3 -3.5 μm
- ✓ Different gain layer species possible:
 - ✓ Boron (standard)
 - ✓ Gallium
 - ✓ Boron + Carbon



- 4" Si-on-Si wafers (High Resistivity ~2 kΩ•cm)
- 50 μm thickness on 250 μm support wafer
- Different implantation species
- Single diodes of active area 0.7 x 0.7 mm
- 5 Neutron and proton fluences tested up to $6 \times 10^{15} \text{ n}_{\text{eq}}/\text{cm}^2$

Standard Boron
Boron + Carbon Spray
Gallium

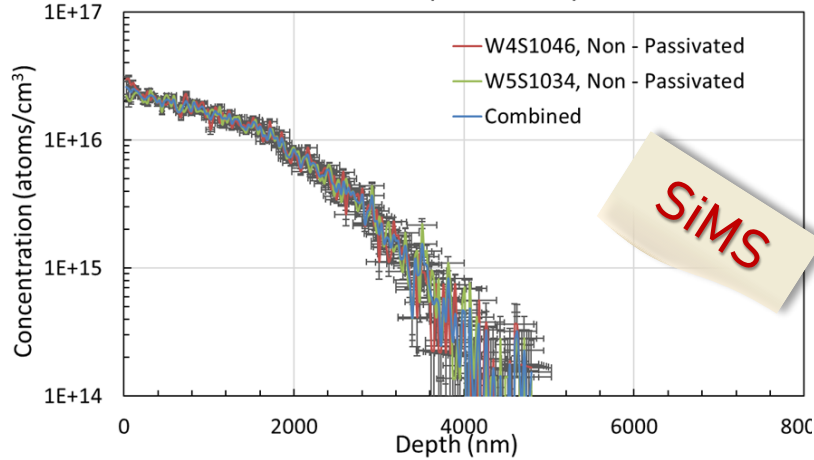
•Introduction III

E. - L. Gkougkousis: - [17th Trento workshop \(2022\)](#)
“Detailed process characterization of carbonated LGADs through
Secondary Ion Mass Spectroscopy”

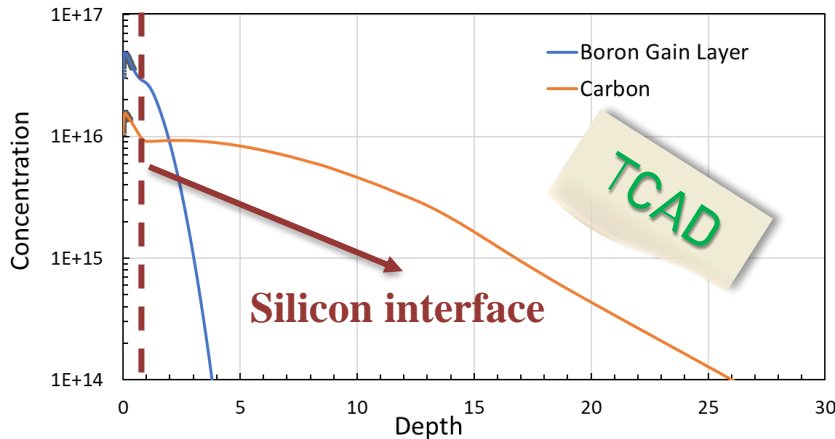
Doping Profiles - Carbon



CNM Run10478 - Gain Layer - Linearly reduced Profiles



CNM R10478 - TCAD Process simulation

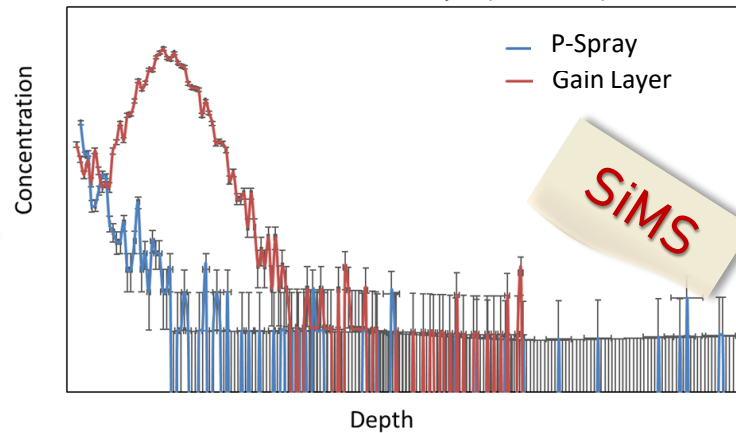


Best ever recorded
resolution



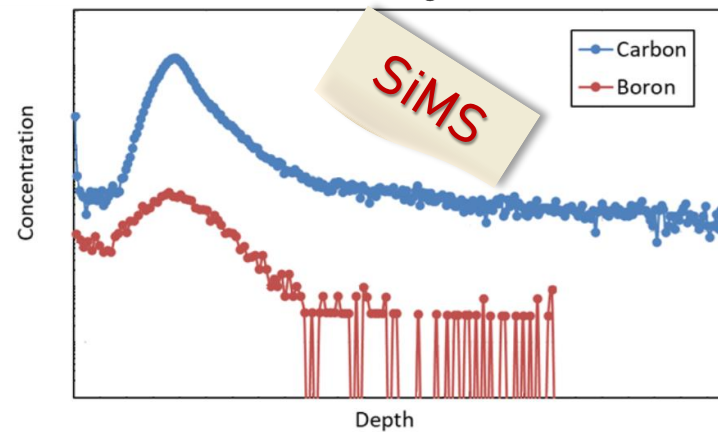
Element	Sensitivity (atoms/cm ³)
Carbon	$(2.22 \pm 0.08) \times 10^{16}$
Boron	$(1.35 \pm 0.58) \times 10^{14}$

FBK UFSD 2 - Gain Layer (In-Silicon)



Boron

FBK UFSD2, High Carbon

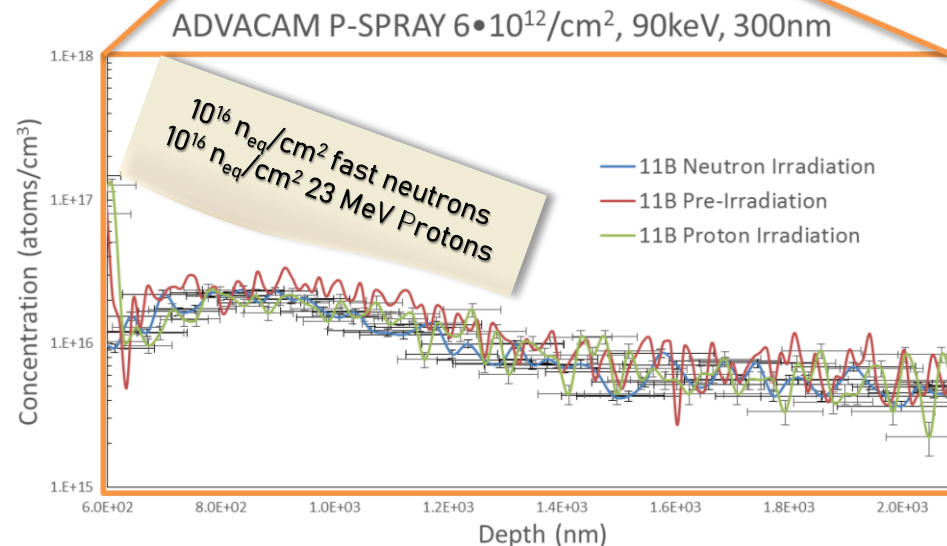
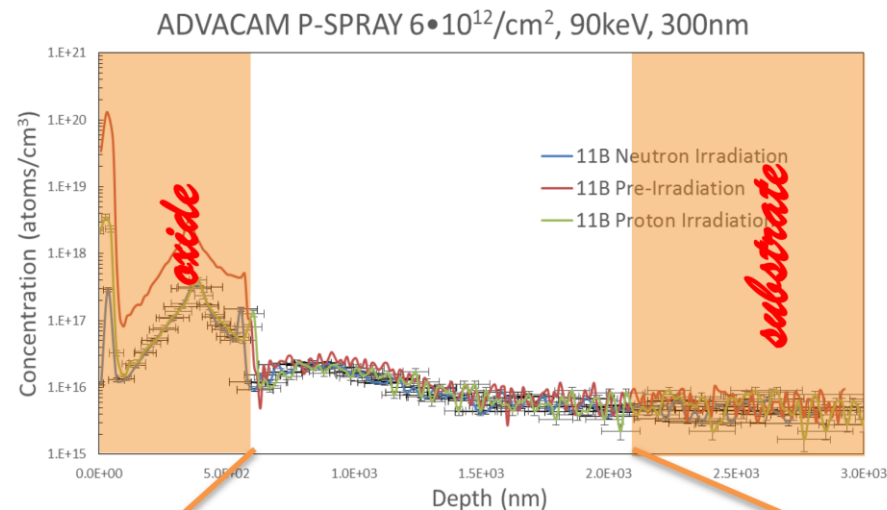


Carbon

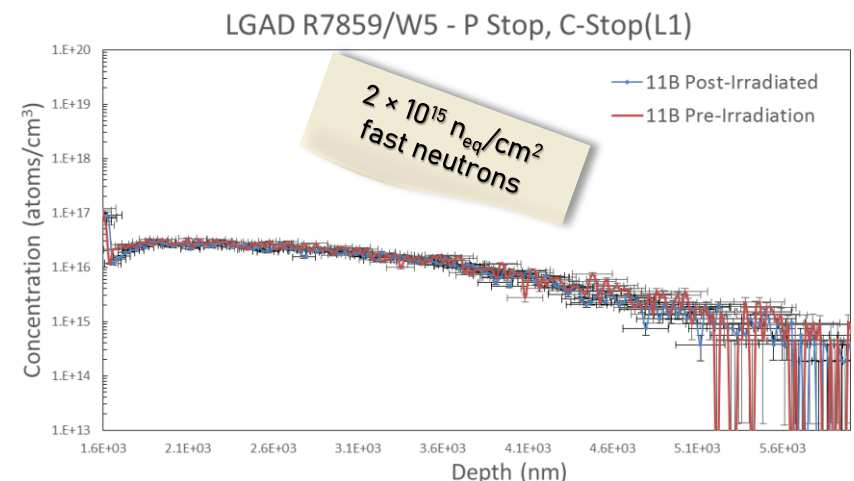
•Introduction IV

E. - L. Gkougkousis: [28th RD50 Workshop, Torino](#)
“Neutron Irradiated doping profile evaluation”
June 2016

Post-Irradiation Doping Profiles



- ✓ SiMS on irradiated boron implanted structures with high sensitivity
- ✓ High ($> 2 \text{ k}\Omega \cdot \text{cm}$) and low resistivity samples ($< 2 \text{ }\Omega \cdot \text{cm}$) p-type substrates tested under both proton and neutron irradiation (high resistivity is non-oxygenated, low is oxygenated)
- ✓ Up to fluences of $1 \times 10^{16} n_{\text{eq}}/\text{cm}^2$ no dopant redistribution was observed
- ✓ Boron DOES NOT diffuse (even as interstitial) under standard operation in sensors, neither does phosphorus



•Radiation Effects I

The Hamburg N_{eff} Model

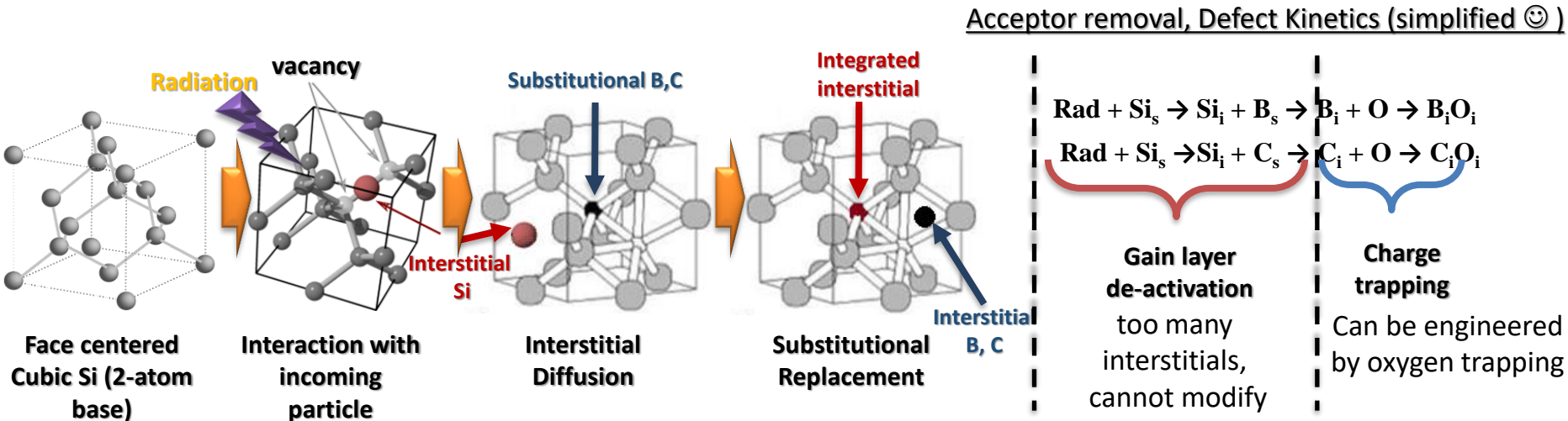
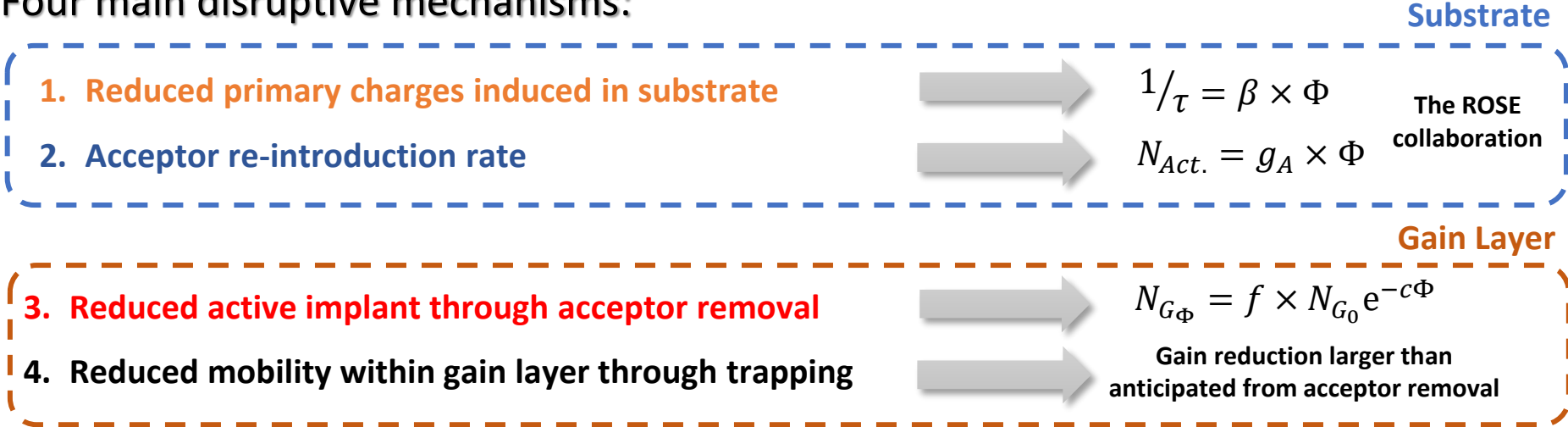
G. Lindstrom et al., NIM A 466(2001) 308-326
["Radiation damage in silicon detectors"](#)

Radiation damage modeling		
Constant Damage Terms	Acceptor Introduction	$\frac{dN_{acc.}^{con.}(t)}{dt} = g_{C_A} \times \Phi_{eq}(t)$
	Donor Introduction	$\frac{dN_{don.}^{con.}(t)}{dt} = g_{C_D} \times \Phi_{eq}(t)$
	Acceptor Removal	$\frac{dN_{acc.}^{rem.}(t)}{dt} = -c_{C_A} \times \Phi_{eq}(t) \times N_{acc.}^{rem.}(t)$
	Donor Removal	$\frac{dN_{don.}^{rem.}(t)}{dt} = -c_{C_D} \times \Phi_{eq}(t) \times N_{acc.}^{rem.}(t)$
Short term annealing	Acceptor Reduction	$\frac{dN_{acc.}^{short.}(t)}{dt} = g_A \times \Phi_{eq}(t) - k_A(T) \times N_{acc.}^{short.}(t)$
Long term annealing	Max Introducible Acceptors	$\frac{dN_{acc.}^{Max. long.}(t)}{dt} = g_Y \times \Phi_{eq}(t) - k_Y(T) \times N_{acc.}^{Max. long.}(t)$
	Acceptor Introduction	$\frac{dN_{acc.}^{long.}(t)}{dt} = k_Y(T) \times N_{acc.}^{Max. long.}(t)$

• Radiation Effects II

E. - L. Gkougkousis: [TIPP2021, May 2021](#)
“Comprehensive technology study of radiation hard LGADs”

Four main disruptive mechanisms:

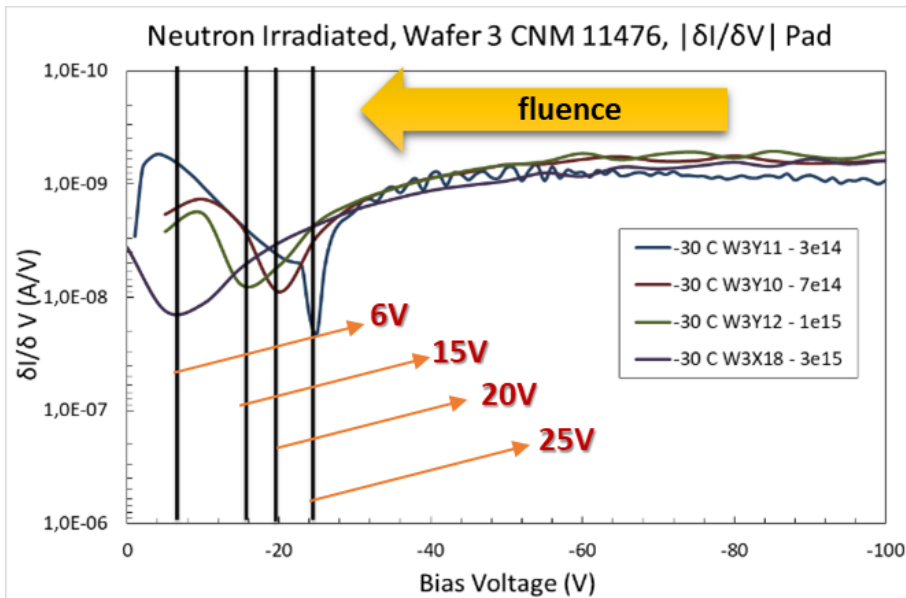
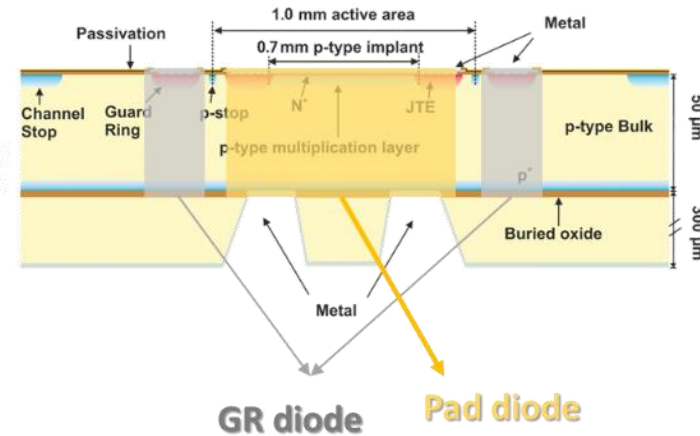


•The Leakage Current Transition Method (LCT) - I

$$f(V) = \left| \frac{\partial I_{pad}}{\partial V} \right|$$

E. - L. Gkougkousis et al.: "Comprehensive technology study of radiation hard LGADs"
J. Phys.: Conf. Ser. 2374 012175

- ✓ Probe active implant via depletion voltage
- ✓ Additional p-implantation gain layer creates secondary depletion region
- ✓ Mott-Schottky equation → **leakage current variation at gain layer depletion**
- ✓ Form of $|\partial I / \partial V|$ at depletion point corresponds to dopant transition function convoluted with instrument resolution (Gaussian X Gaussian)
- ✓ Depletion voltage determined Gaussian fit at depletion voltage for -10°C, -20°C & -30°C



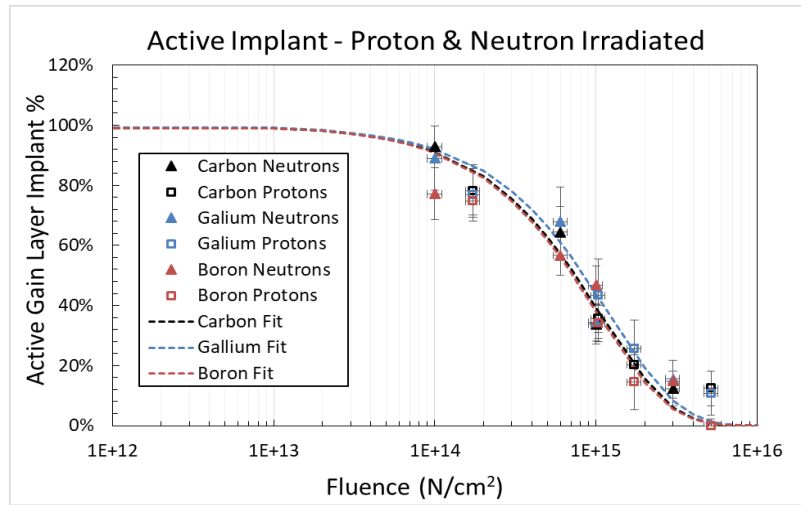
- Independent Gaussian fit for temperature
- Uncertainties estimated from propagation on fit sigma
- Fluences up to $3 \cdot 10^{15}$ n_{eq}/cm^2 in p^+ and n^0

$$V_d = \frac{\sum_{T=-10}^{-30} V_{d,T_i}}{n_T}$$

$$\delta V_d = \sqrt{V_{d,sys} + V_{d,stat}}$$

Average of fit sigma Standard deviation of V_d

•The Leakage Current Transition Method (LCT) - I



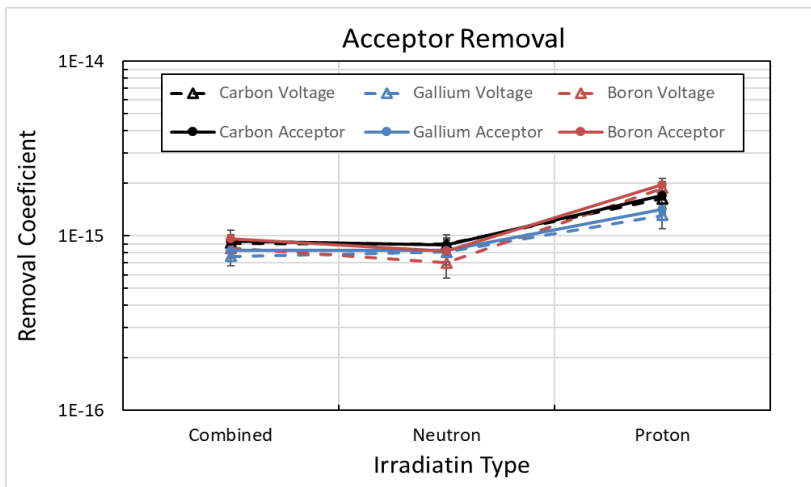
- Linear dependence assumption between V_{GL} and active implant
- Normalized exponential reduction fit model on gain and V_{GL}

$$G(\%) = e^{-C_G \Phi}$$

- Linearity hypothesis tested with independent C_v and C_G fits – full compatibility
- Constraints imposed on initial values to reflect charge measurements

Results

- **Compatible acceptor removal coefficients between all implants**
- Slight Ga advantage in p⁺ irradiation (23 GeV/c PS), higher mass reduces displacement probability in coulomb-only (far-field) interactions
- Quasi-identical performance for neutron irradiated (fast ~10MeV neutrons)
- **Identical gain layer de-activation for all dopants with fluence**

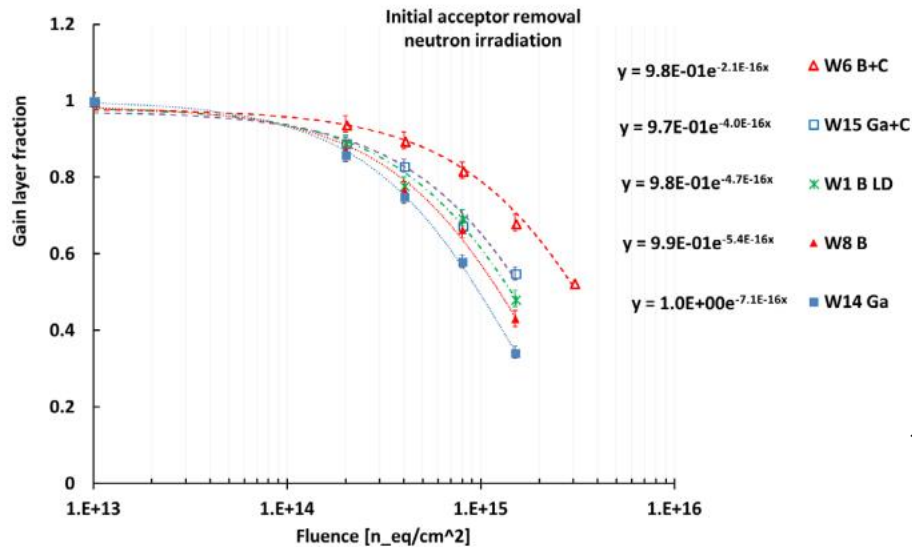


Acceptor Removal Coefficient	
Gallium	$(8.25 \pm 0.80) \times 10^{-16}$
Boron + Carbon	$(9.33 \pm 0.78) \times 10^{-16}$
Boron	$(9.69 \pm 1.04) \times 10^{-16}$

•The other side of the coin – FBK Carbonated

M. Ferrero et al.: [“Radiation resistant LGAD design”](#)

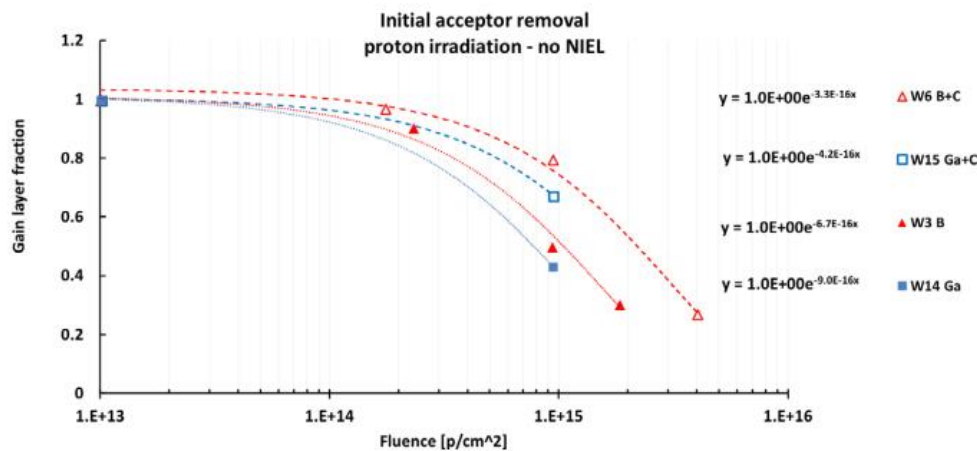
NIMA, Volume 919, 1 March 2019, Pages 16-26



- ✓ Carbon is directly implanted at the gain layer only with comparable concentrations
- ✓ An improvement is seen on the acceptor removal coefficient by a factor of 2

What does this mean?

- ✓ Irradiated Implants do not diffuse (to the nm level)
- ✓ Carbon only helps in acceptor removal when close to boron



**Acceptor removal
is a local process**

- ✓ The fact that proton and neutron irradiations fit in the same curve means that this is a point defect sensitive effect

• Radiation Hardness

Gain Layer Engineering

- ✓ First approximation: gain equivalent to charge in parallel plane capacitor:

$$G \sim e^{a \cdot d}$$

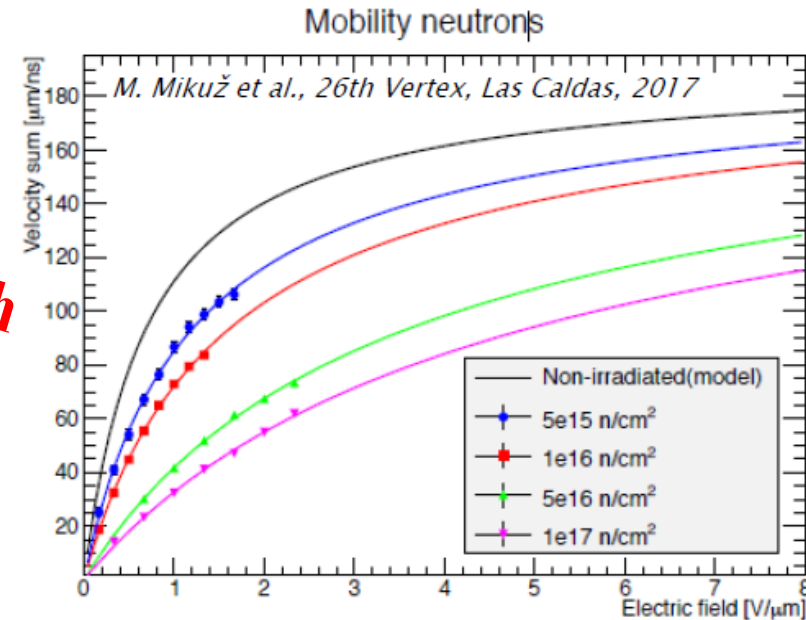
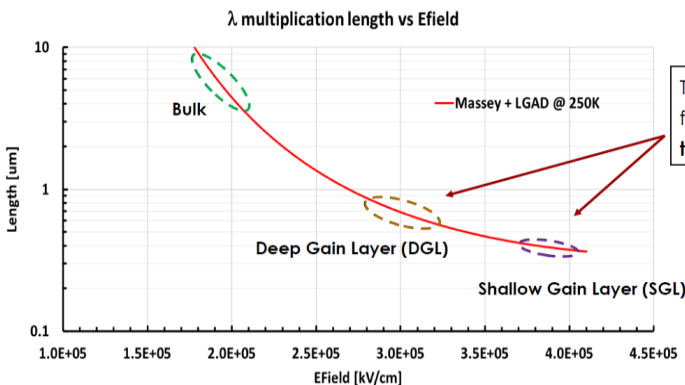
d: distance (gain layer thickness)
a function equivalent to inverse of mean free path ($1/\lambda$)
- ✓ In irradiated silicon, λ depends on fluence, temperature and field
- ✓ Higher fields mean shorter distances to acquire same kinetic energy
- ✓ Presence of scattering centres has to be compensated with higher fields

$$G(E, T, \phi, d) \propto e^{\alpha(E, T, \phi) \cdot d}$$

$\alpha(E, T, \phi)$ impact ionization coefficient

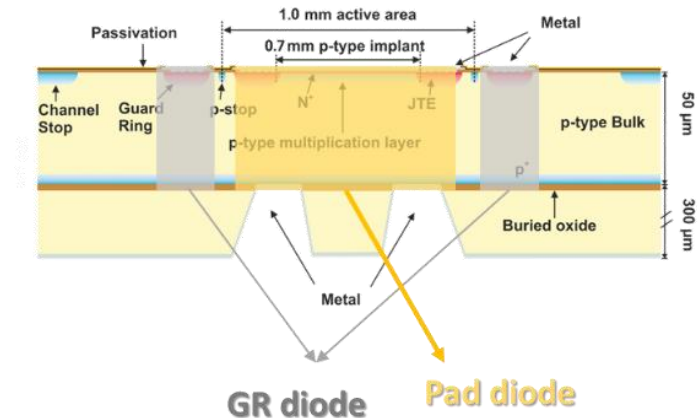
d = length of high E field

Gain depletion not enough



•Comparative Current to Gain Method (CG2C) - I

- ✓ Acceptor removal only gives information about active dopant, not gain
- ✓ Gain also depends on **trapping levels & doping profiles**
- ✓ Effects after irradiation for different defect concentrations
- ✓ For same amount of acceptor removal, different gain reduction expected



	Before Irradiation	After Irradiation
Pad Leakage Current $I_{pad}^{\Phi=0} = S \times I_s \times \left(e^{\frac{eV}{nkT}} - 1 \right) \times G(e^V, T, 0)$	$I_{pad}(\Phi) = S \times (I_{GR}^{\Phi=0} + \alpha\Phi) \times G(e^V, T, \Phi)$	$= f(V, \Phi)$
Guard Ring Leakage Current $I_{GR}^{\Phi=0} = I_s \times \left(e^{\frac{eV}{nkT}} - 1 \right)$	$I_{GR}(\Phi) = I_{GR}^{\Phi=0} + \alpha\Phi$	$\text{Expected substrate current increase}$

If we divide the two then:

$$f(V, \Phi) = \left| \frac{I_{pad}}{I_{GR}} \right|$$

Normalize with unradicated

➔

$$\frac{f(\Phi)}{f(\Phi = 0)} \sim G(e^V, T, \Phi)$$

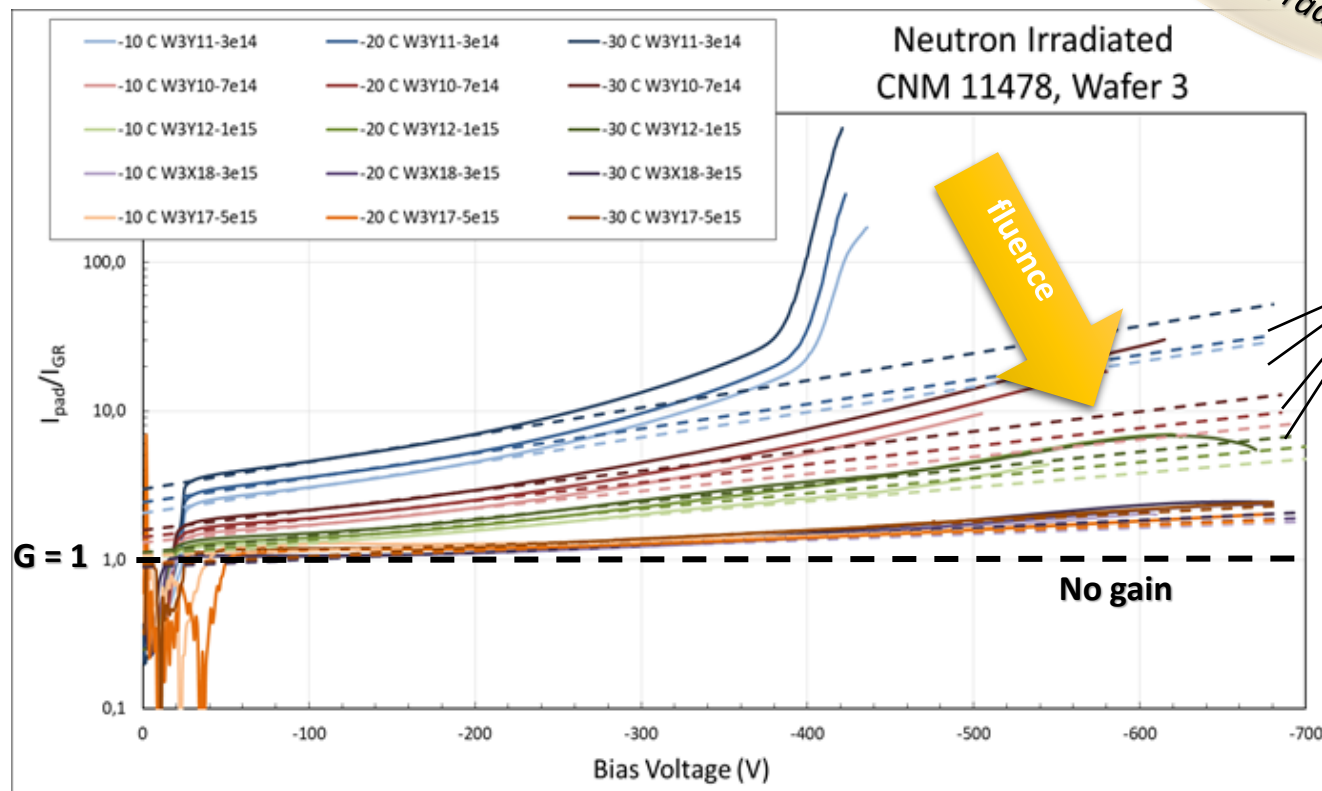
1. GR and pad on same substrate, all non-gain related irradiation effects can be normalized
2. Assumption that differences between GR n-type implant and pad n-type implant have minimal effects

•Comparative Current to Gain Method (CG2C) - II

- ✓ I_{GR}/I_{PAD} linear at the semi-log plane
- ✓ Gain Coefficient probed by slope of linear fit
- ✓ Different fits per temperature, reputed at -10 °C, -20 °C and -30 °C

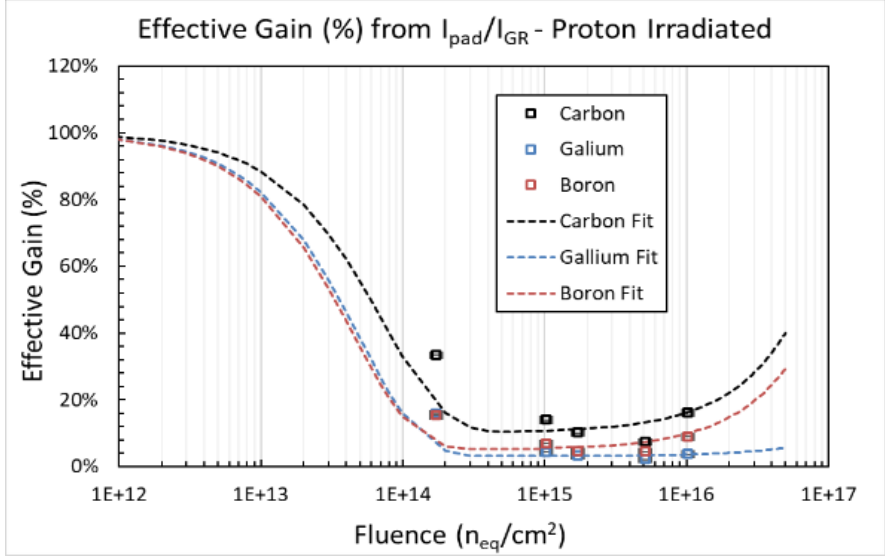
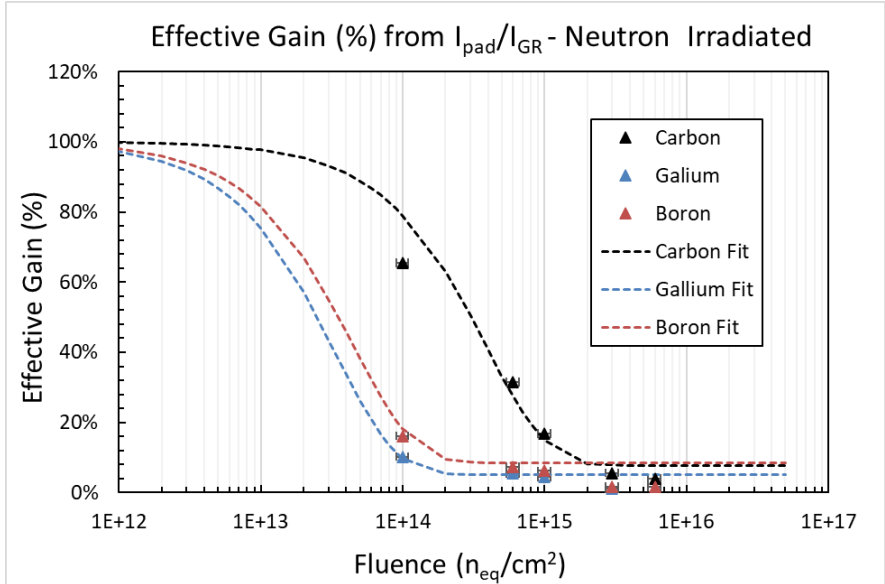
E. - L. Gkougkousis: [“Review of neutron irradiated 6” Sol LGAD sensors CNM 11486”](#)

35th RD50 Workshop, November 2019



Repeat for all implants
and radiation types

•Comparative Current to Gain Method (CG2C) - III



Acceptor level introduction rate

$$N_{eff}(\Phi) = N_{eff0} - N_c(1 - e^{-c\Phi}) + g_c\Phi$$

Effective dopant concentration Initial dopant concentration Removable dopant Gain extraction constant

Gain Reduction Coefficient	
Irrad. Type	C ± δC
Gallium	
n ⁰	(3.01 ± 0.9) × 10 ⁻¹⁴
p ⁺	(2.02 ± 0.11) × 10 ⁻¹⁴
Boron + Carbon	
n ⁰	(2.57 ± 1.1) × 10 ⁻¹⁵
p ⁺	(1.37 ± 0.24) × 10 ⁻¹⁴
Standard Boron	
n ⁰	(2.25 ± 0.39) × 10 ⁻¹⁴
p ⁺	(2.25 ± 0.28) × 10 ⁻¹⁴

Results

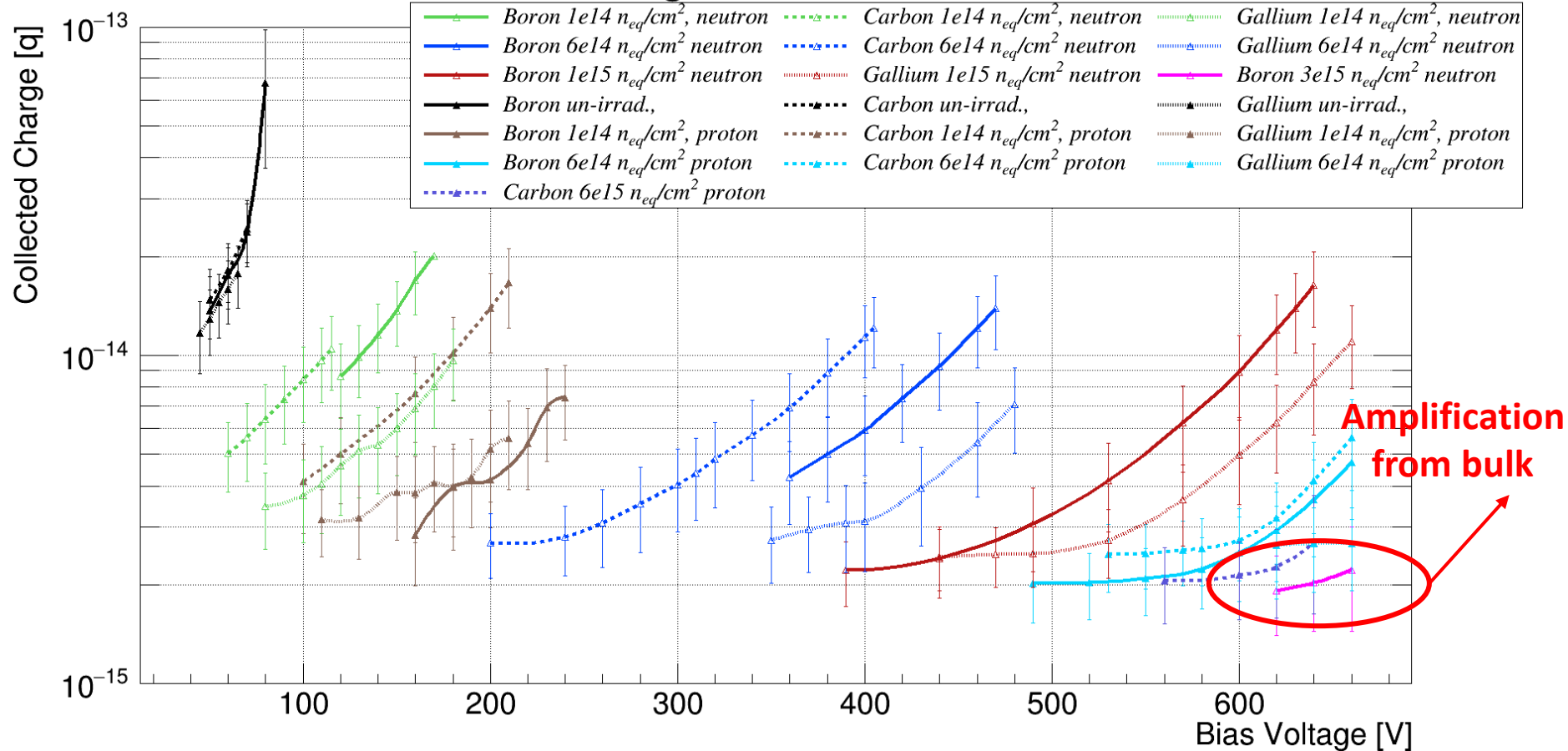
- Gallium and Boron perform similarly
- **Carbon + Boron is up to 2 times better in proton and up to 7-8 times better in neutron irradiation**
- Significant variation with implant type
- **Gain reduction coefficients are up to 10 x the previously estimated acceptor removal**

•Collected Charge I

~940 series of measurements

- ✓ Each point corresponds to MPV of Landau x Gauss fit on 5k recorded events
- ✓ Measurements repeated for -10°C, -20°C & -30°C (see the backup)
- ✓ Gallium yields always 20% less charge for same voltage, carbonated 20% more

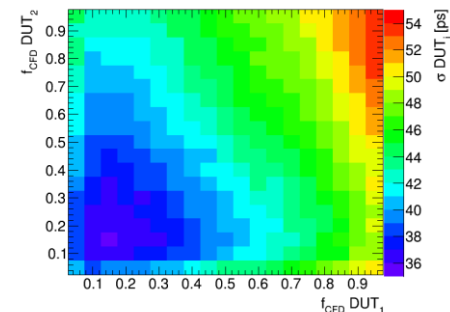
Collected Charge - B, Ga & B+C at -30°C



Time Resolution

E. - L. Gkougkousis: “[Acceptor removal and gain Reduction in proton and neutron irradiated LGADs](#)”

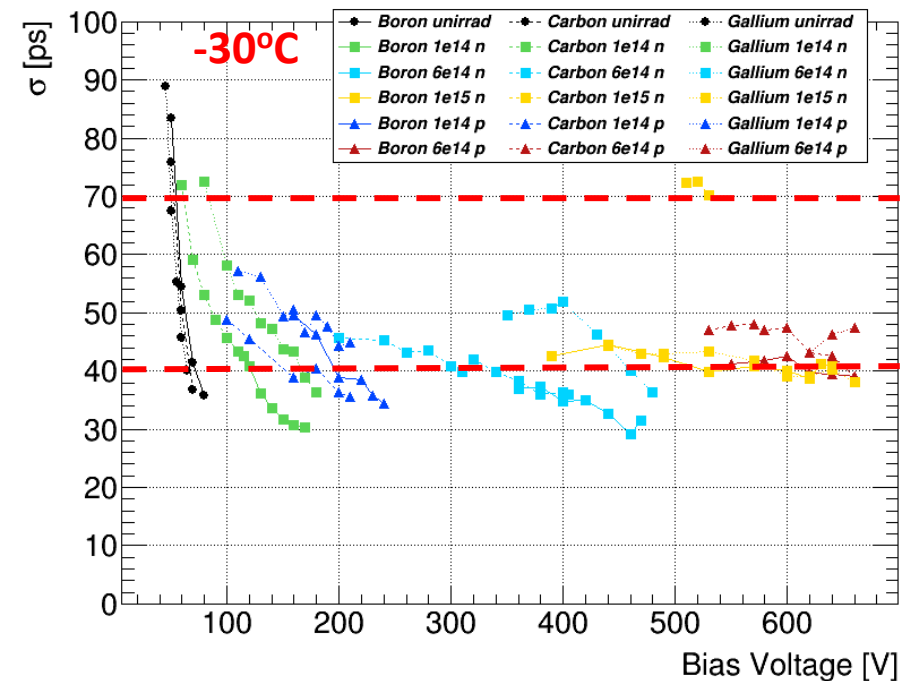
36th RD50 Workshop, June 2020



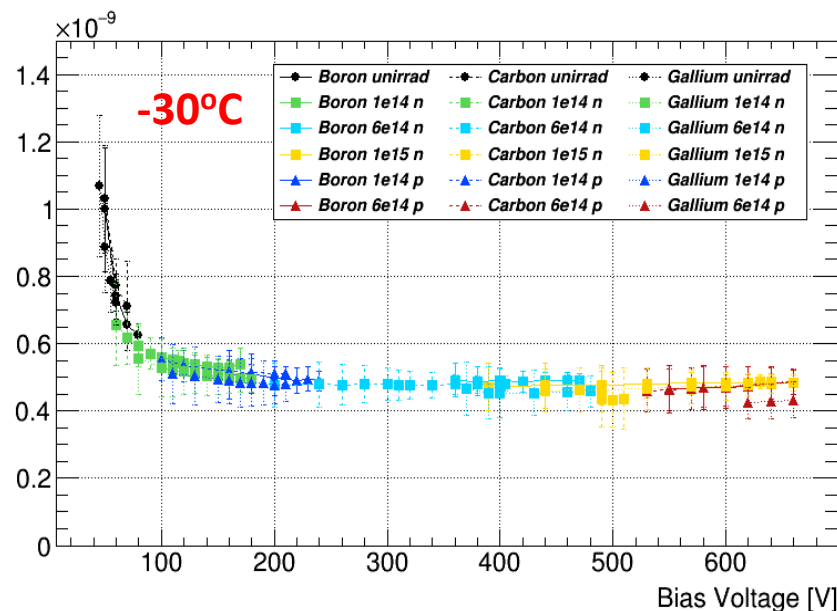
CFD Level optimization

$$(\sigma_{Dut})_{CFD_{ij}} = \sqrt{(\sigma_{Tot}^2)_{CFD_{ij}} - (\sigma_{Ref}^2)_{CFD_i}}$$

2D optimization plot – 0.5% binning



HL-LHC Range
Hise time (10%-90%) [s]

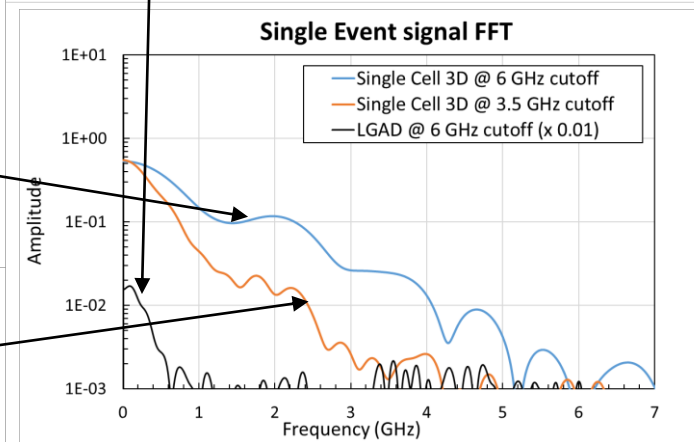
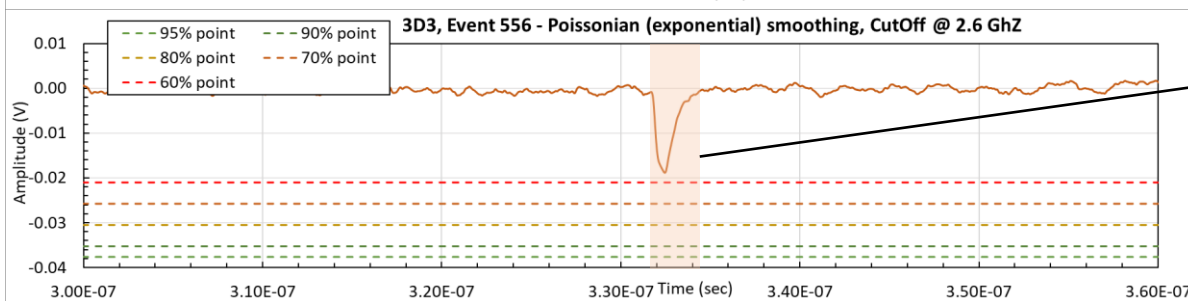
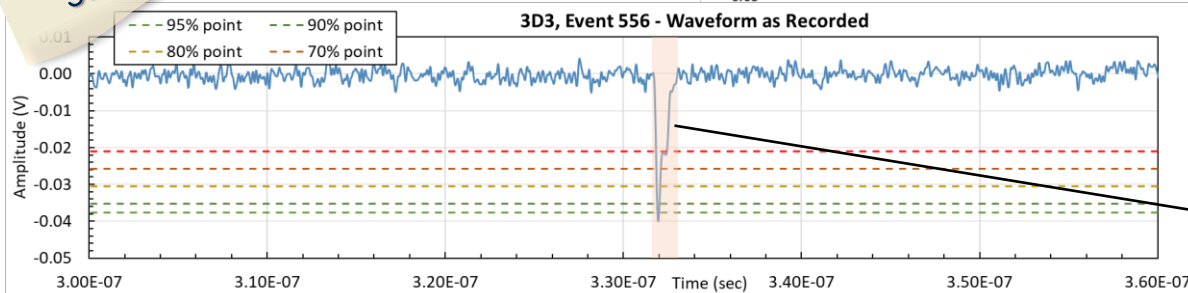
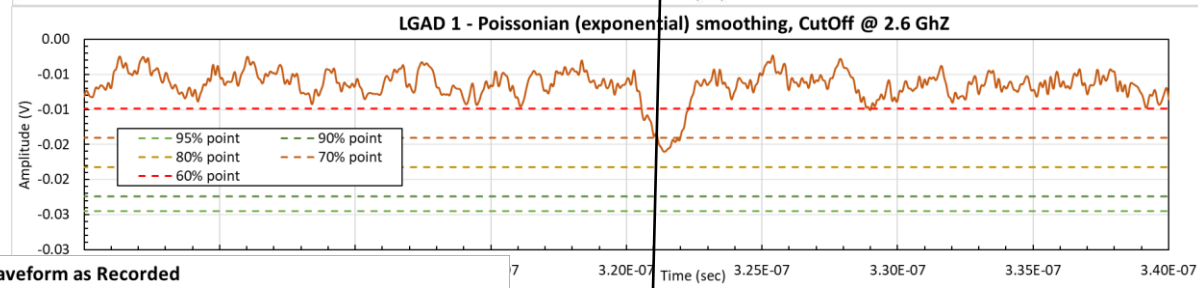
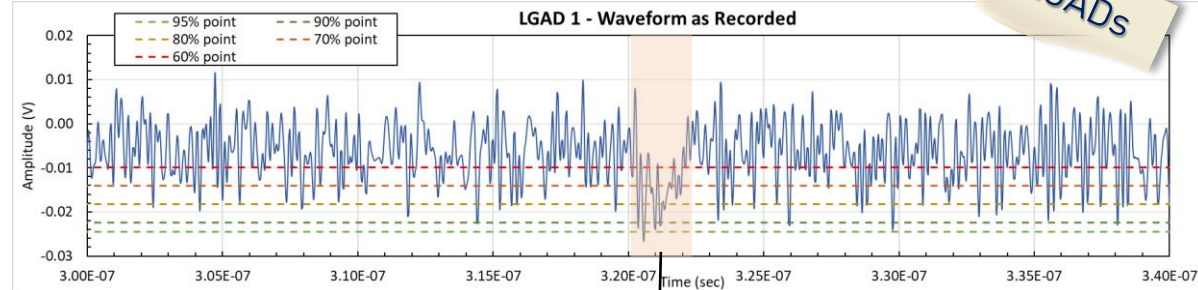
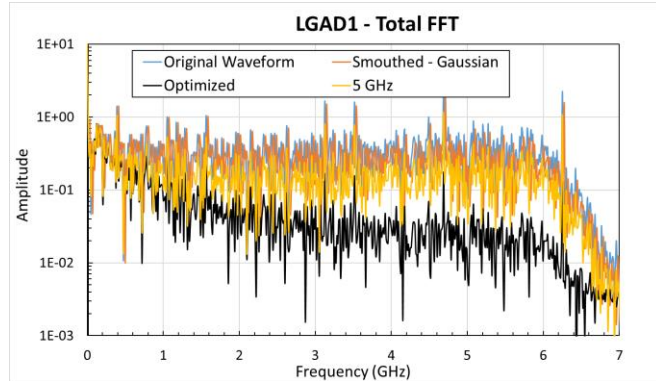


Time Resolution: $\sigma_{tot}^2 = \underbrace{\sigma_{timewalk}^2}_{\sigma_{Dist.}^2 + \sigma_{Landau}^2} + \underbrace{\sigma_{jitter}^2}_{\left(\frac{t_{rise}}{S/N}\right)^2} + \underbrace{\sigma_{conversion}^2}_{\left(\frac{TDC_{bin}}{\sqrt{12}}\right)^2} + \underbrace{\sigma_{Clock}^2}_{\text{Fixed Term } \sim 5-7 \text{ psec}}$

1. Similar behavior in terms of signal shape on all implants
2. Time resolution follow charge trend
3. Charge vs σt identical for all gain layer variations

Signal Analysis

The importance of Bandwidth



**60 % less amplitude
20 % more charge**

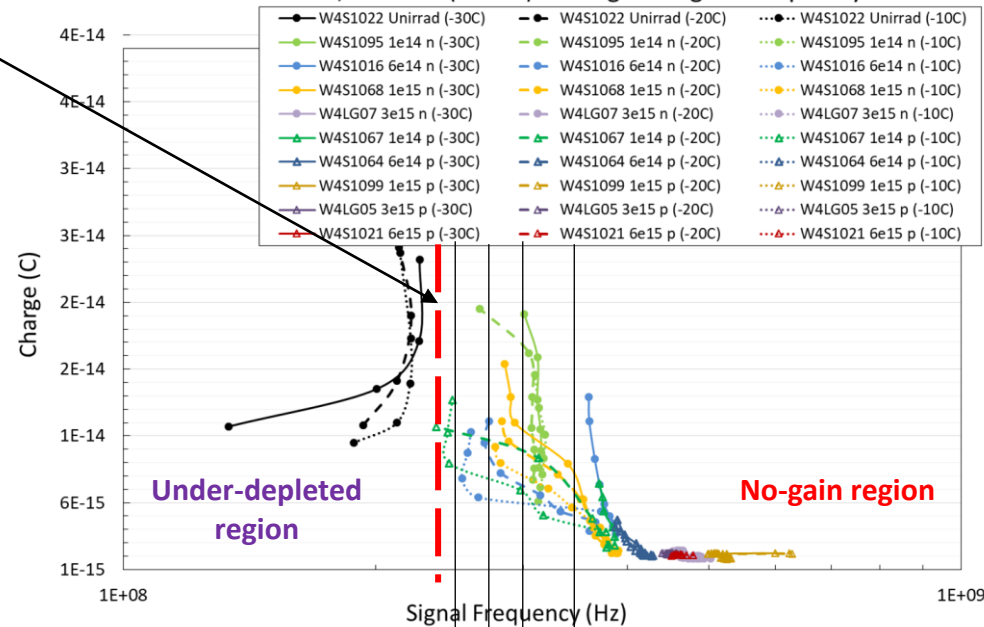
• Signal Analysis

FFT

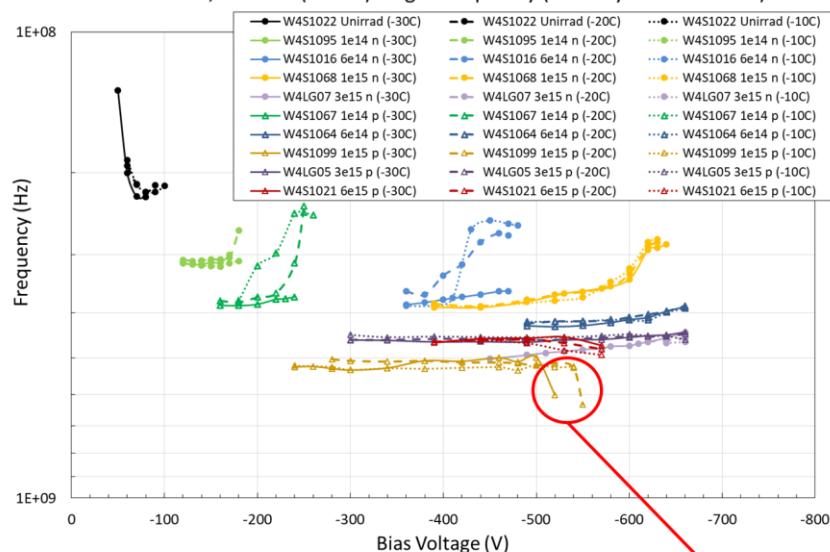
- ✓ FFT vs Voltage presents an asymptotic behavior towards a frequency
- ✓ Asymptotic frequency depends on fluence and remaining gain
- ✓ Signal frequency increases with voltage and decreases on the onset of multiplication

Asymptotic point ~ 250 MHz

CNM 10478, Wafer 4 (Boron) - Charge vs Signal Frequency



CNM 10478, Wafer 4 (Boron) - Signal Frequency (Primary FFT Harmonic)



Fluence

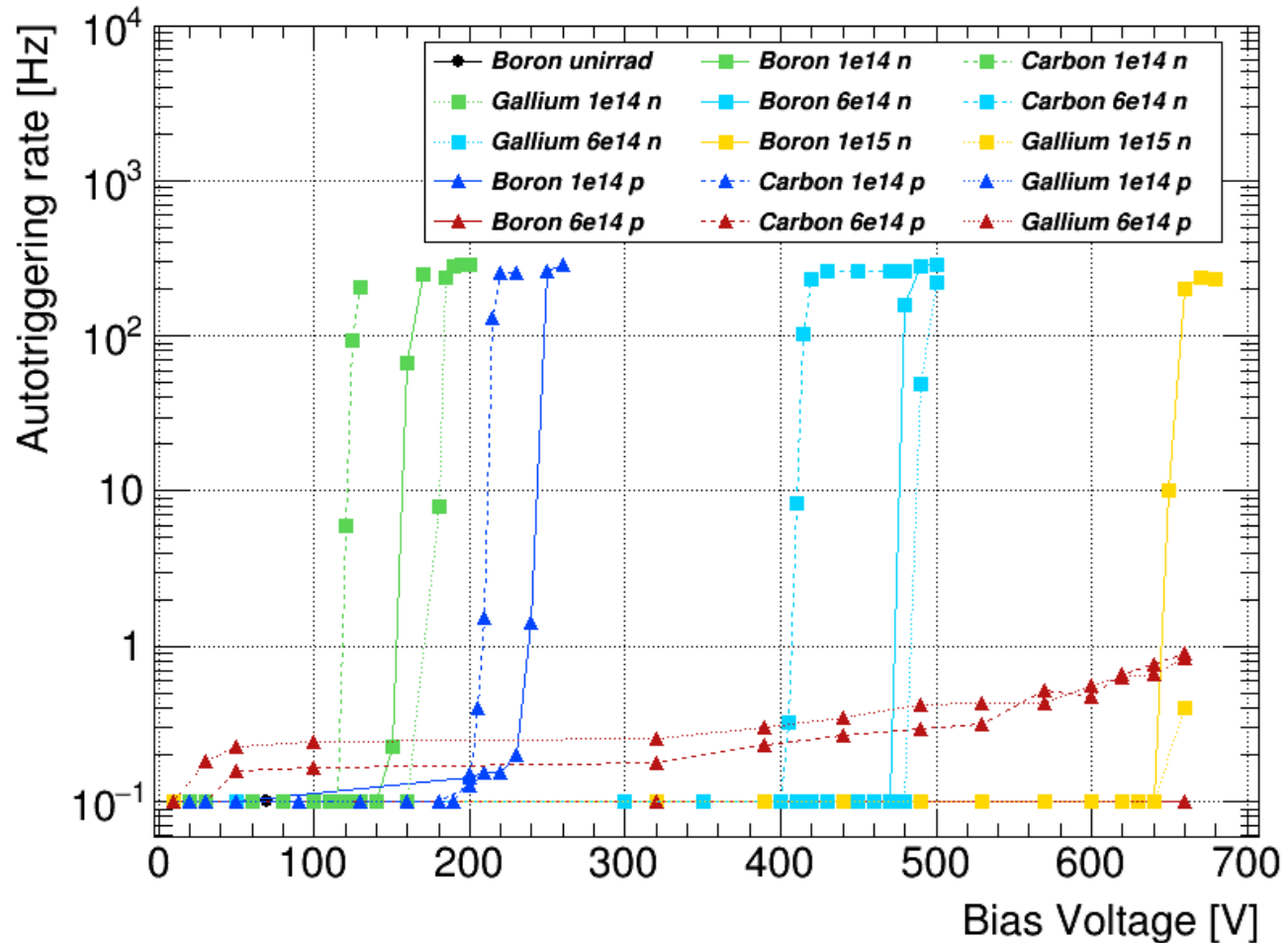
Gain

Asymptotics move to higher and higher frequencies as gain decreases and fluence increases

• Sensor Stability

Dark Rate

- ✓ All sensors with gain present dark rate at high field values
- ✓ Dark rate events result out of thermally induced electron-hole pairs drifting picked up by the field
- ✓ Random in nature, follow a Poisson distribution
- ✓ An inverted s-curve study for each sensor defines the stable region or the acceptable level of shot noise.



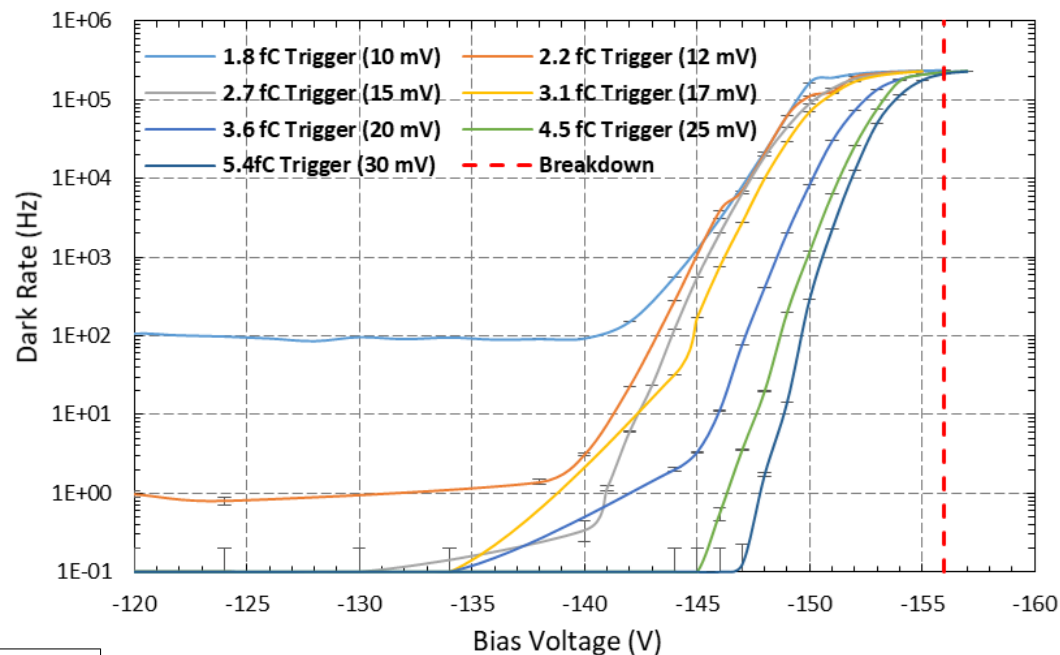
• Stability vs Threshold

Dark Rate

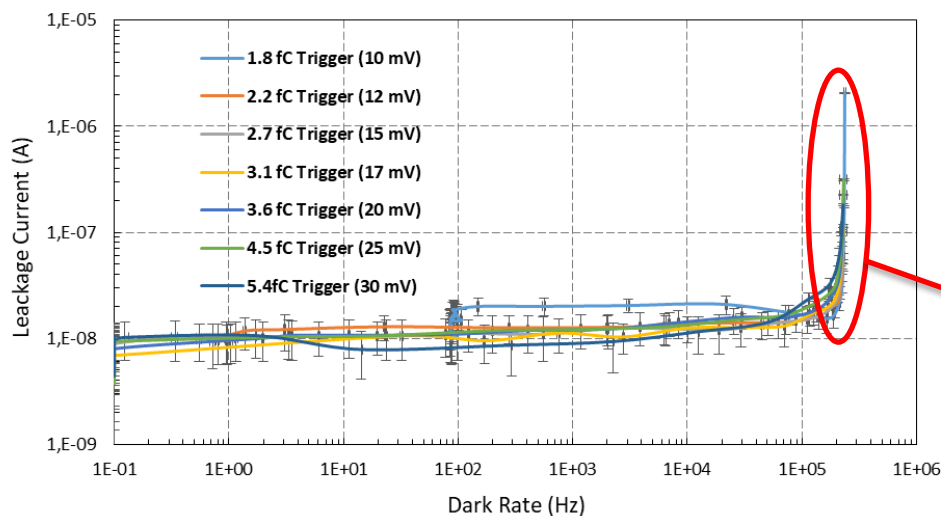
Threshold effect

- Un-irradiated HPK P2
- Breakdown ~ 156 V
- Measured at room temp
- Different Constant threshold triggers (1.8 – 5.4 fC) applied
- Bayesian Uncertainties
- Max saturation rate 230 kHz

HPK - P2W25 L17P12, Room Temp.



HPK -P2W25 L17P12, Room Temp.



Leakage Current effect

- Sensor far from breakdown
- Leakage current not demonstrate significant variation
- Stationary leakage current at exponential rate increase, breakdown over $1e5$ Hz

• Efficiency

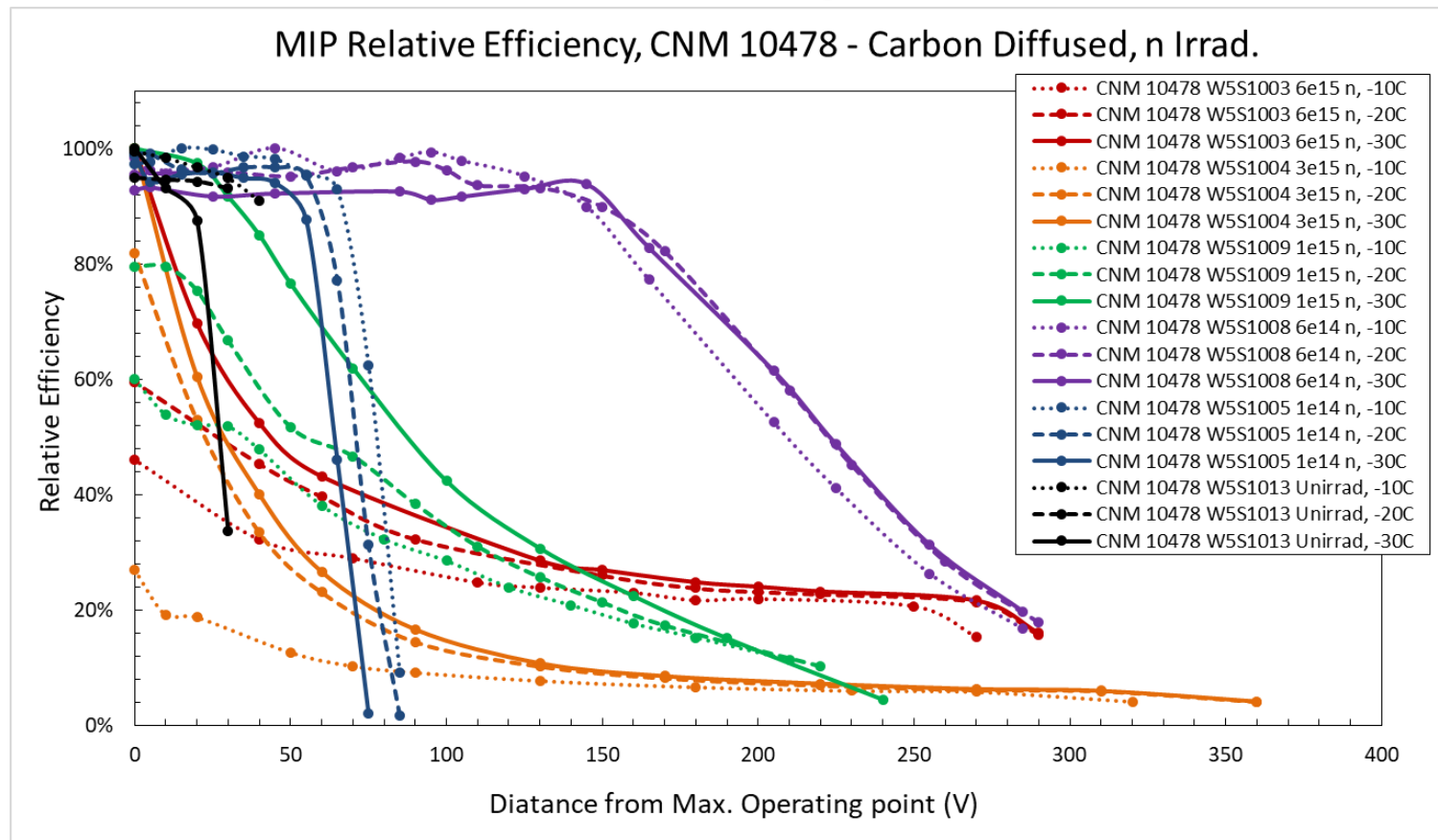
Head Room

- Measurements with radioactive ^{90}Sr source
- Define stable operation points satisfying the following conditions:

Sensor not in breakdown

Mean field inside sensor < 13.4 V/ μm

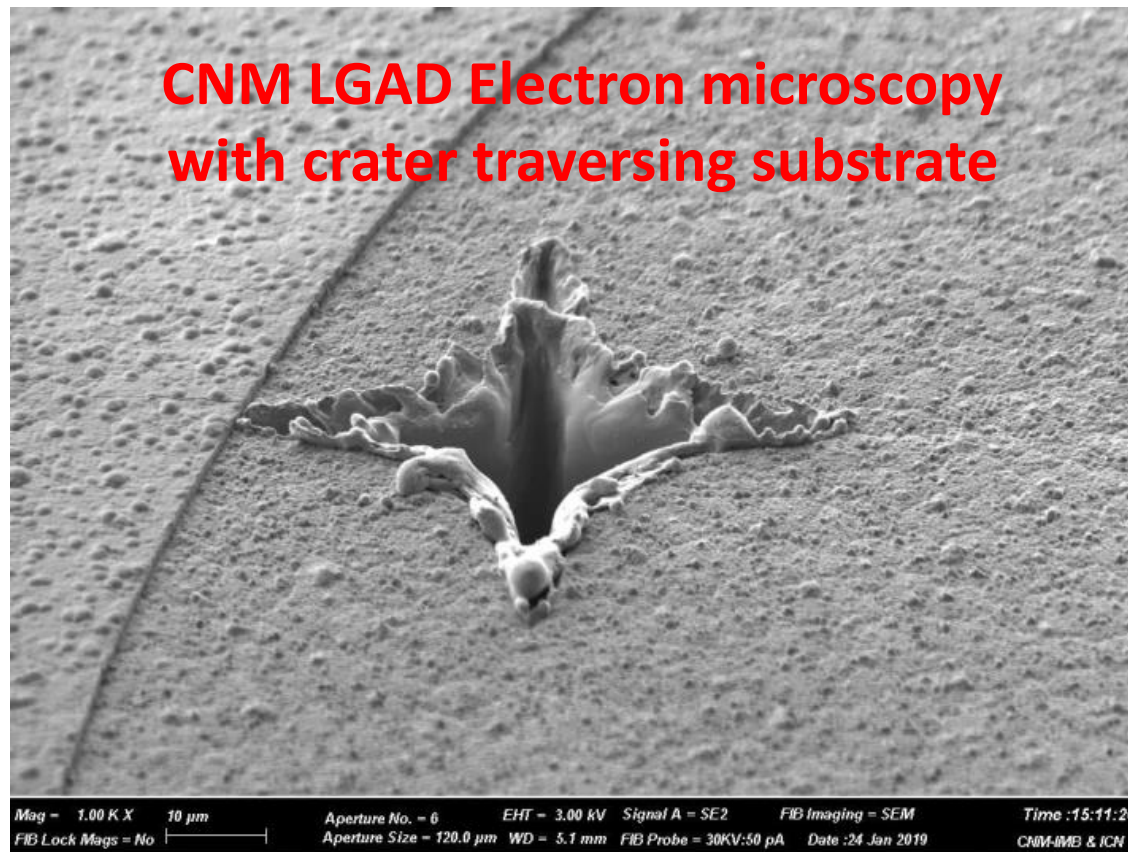
Auto-trigger rate < 1% of source trigger



•Single Event Burn-Out

Catastrophic failure

- Catastrophic breakdown events occur at mean bias voltages of $\sim 12 \text{ V}/\mu\text{m}$ for planar structures
- Effect observed on LGADs and planar pixels after irradiation
- High energy deposition close to a tap cluster creating a highly localized field variation which leads to high gain
- Observed in SPS test beams in 2017 and verified by lasers at ELI beamlines



Future developments

HAB = Half Activated Boron

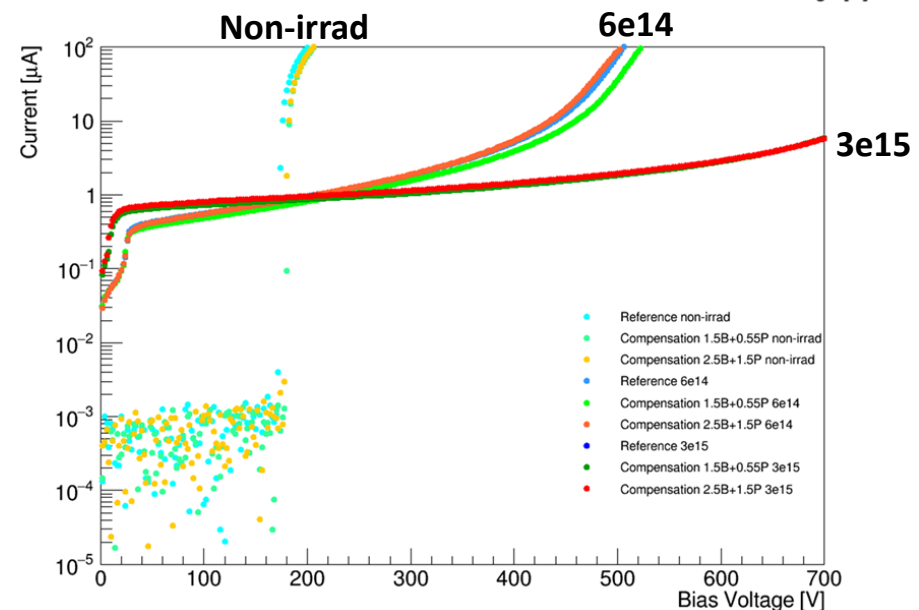
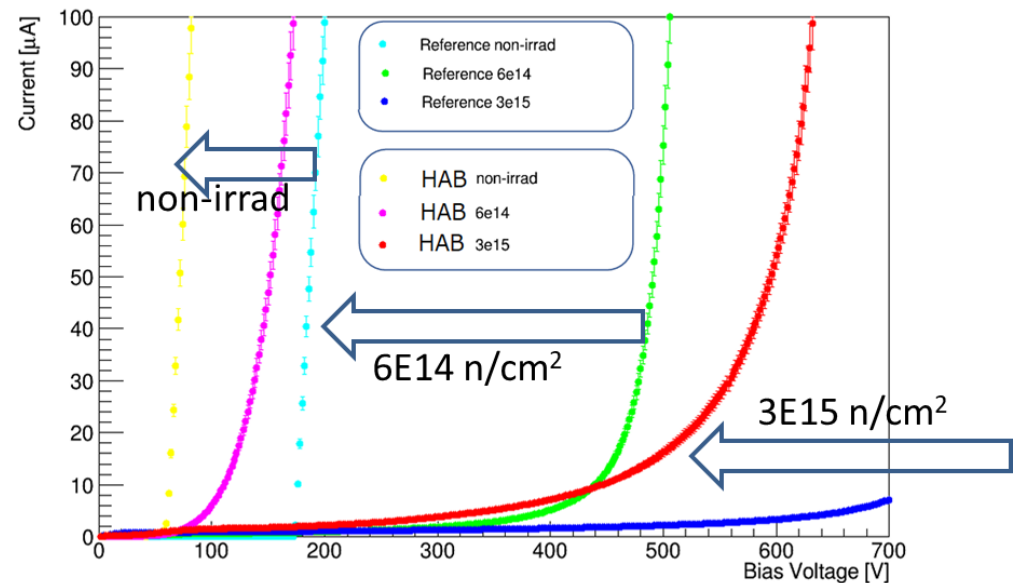
- Stop O₂ boron deactivation pathway by increasing amount of non-activated boron.
- Bi will capture the Oi before it encounters a substitutional boron
- Extremely promising first results

K. Hara et al.: [“Improvement of timing resolution and radiation tolerance for finely segmented AC-LGAD sensors”](#)

Trento 2023

Compensation

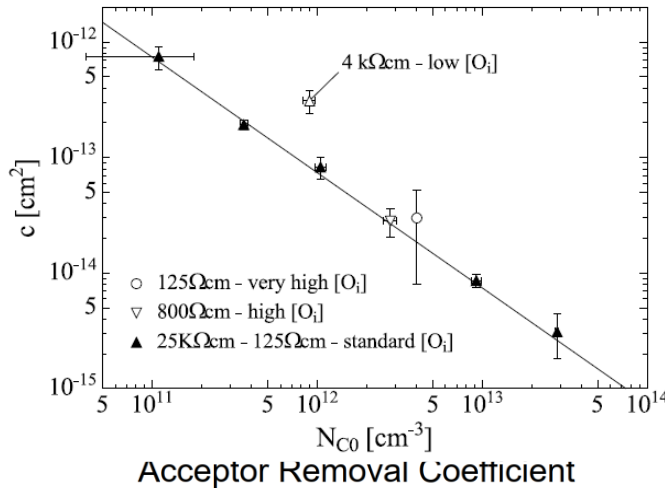
- Increase boron concentration but add some type of n-implant to maintain N_{eff} to acceptable levels
- If $C_D < C_A$ additional acceptors are disengaged to participate to the N_{eff} with irradiation
- First results not promising



•Future developments

Removal Coefficients

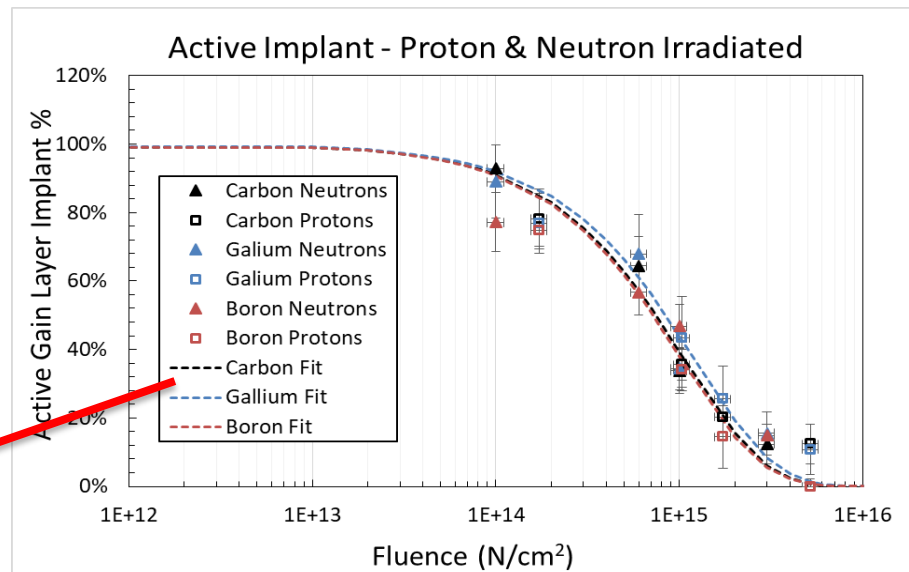
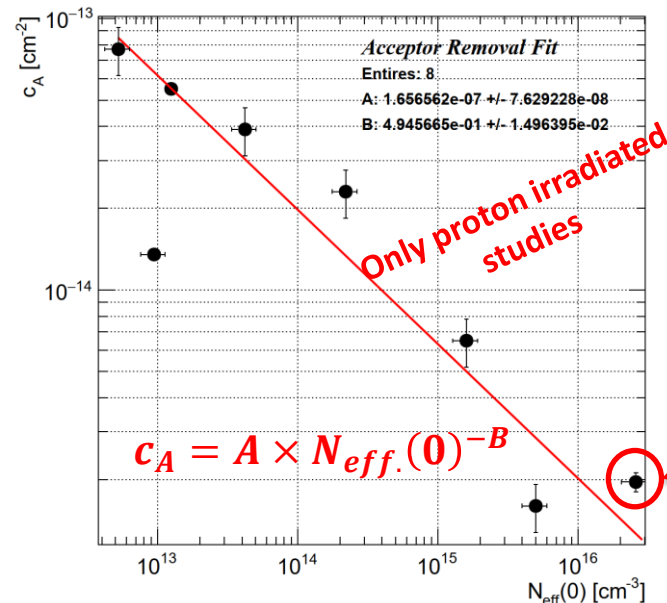
- ✓ Removal coefficient dependent on initial acceptor concentration
- ✓ Almost complete removal for high N_{eff} , ~ 40 % removal for high resistivity substrates



$$N_C(\Phi_{eq.}) = g_c \times \Phi_{eq.} - \underbrace{f_c}_{\text{Removable Fraction}} \times \underbrace{N_{eff.}(0)}_{\text{Original doping concentration}} \underbrace{(1 - e^{-c_A \Phi_{eq.}})}_{\text{Acceptor Removal Coefficient}}$$

Re-Introduction rate g_c

- ✓ > 0 primarily acceptor introduction
- ✓ < 0 primarily donor introduction



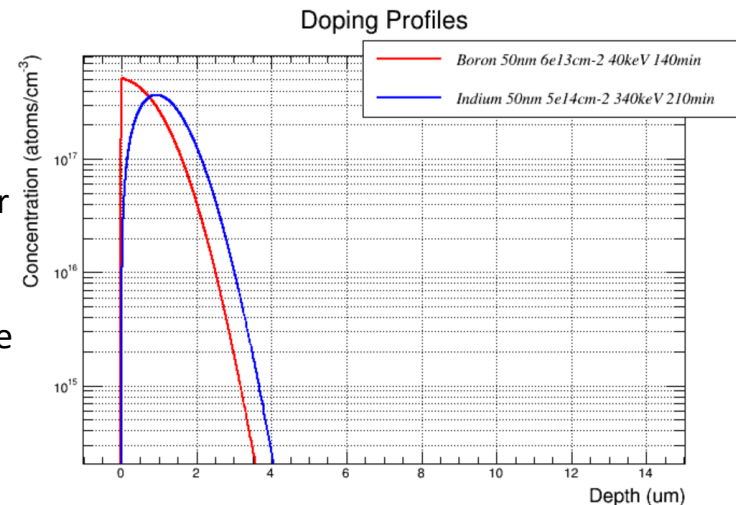
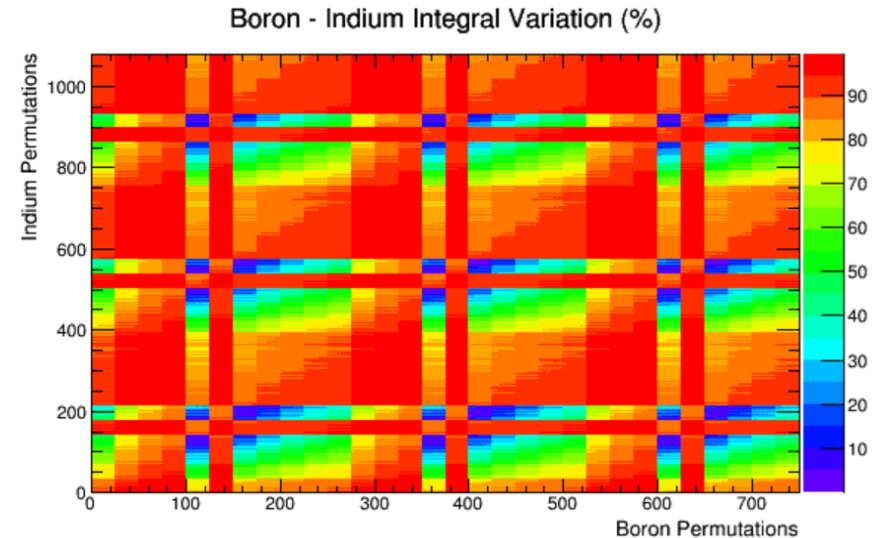
•Future developments

Alternative dopants

E. L. Gkougkousis: [“Parametric process optimization for Indium, Gallium and Boron dopants using TCAD simulation modelling”](#)

16th Trento Workshop

- ✓ Radiation damage lead to acceptor removal though defect kinematics
- ✓ Modify gain layer implants to generate beneficial defects for gain (**gain regulation**):
- **Lithium co-implantation:**
 - Boron with Lithium co-implantation demonstrates better neutron radiation hardness
- **Replace Boron with Indium**
 - Indium higher mass and lower reaction cross-section expected to generated less O_i defect clusters
- ✓ Implantation energy and doping profiles already optimized via TCAD simulations
- Lithium co-implantation ONLY on p-implant layers
 - Lithium is n-type but in low doses should not impact p layer
 - Proven to improve radiation hardness of solar cells after 1MeV neutron irradiation
 - Lowers annealing temperature when implanted in substrate
 - Defect engineering at low temperatures E. Oliviero et Al. ([link](#))
 - Original Solar cell study Weinberg et Al. ([link](#))




• Conclusions

Thoughts and discussion

- Three methods of radiation hardness:

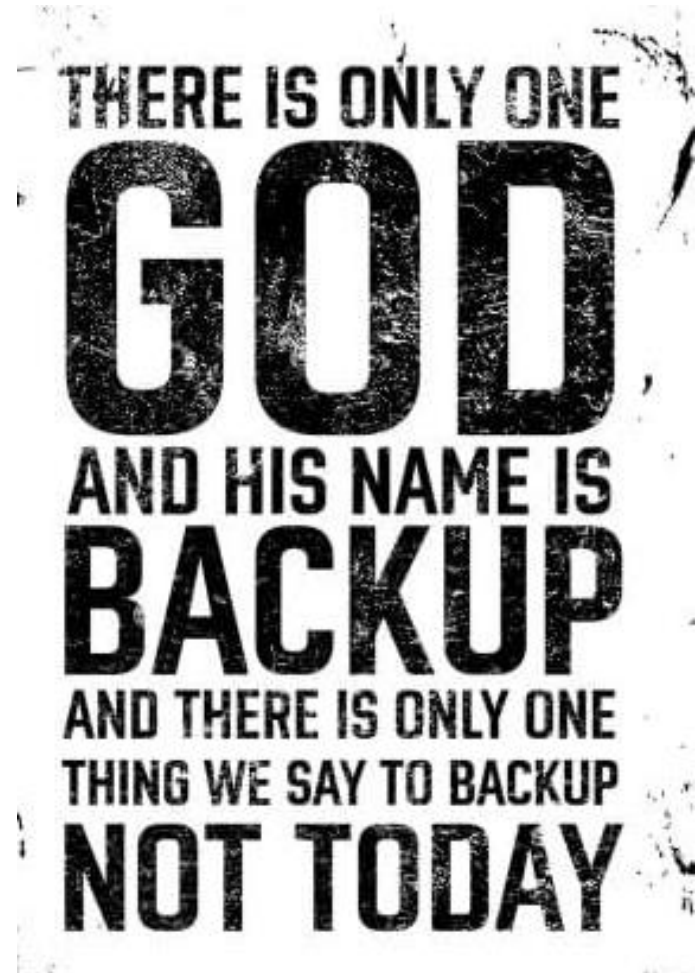
- | | |
|--------------------------------------|--|
| 1. Active Gain Implant: | No measureable improvement wrt different implants |
| 2. Effective Gain Estimation: | Gallium-Boron behave similarly
Carbon up to 2x better in neutrons / protons |
| 3. MIPs Charge collection: | 20 % improvement in required bias for Carbon
20 % degradation for Gallium |



Consistent with defect kinetics theory and an exponential field -gain dependence
Results consistent in all temperatures (-10°C, -20°C, 30°C)

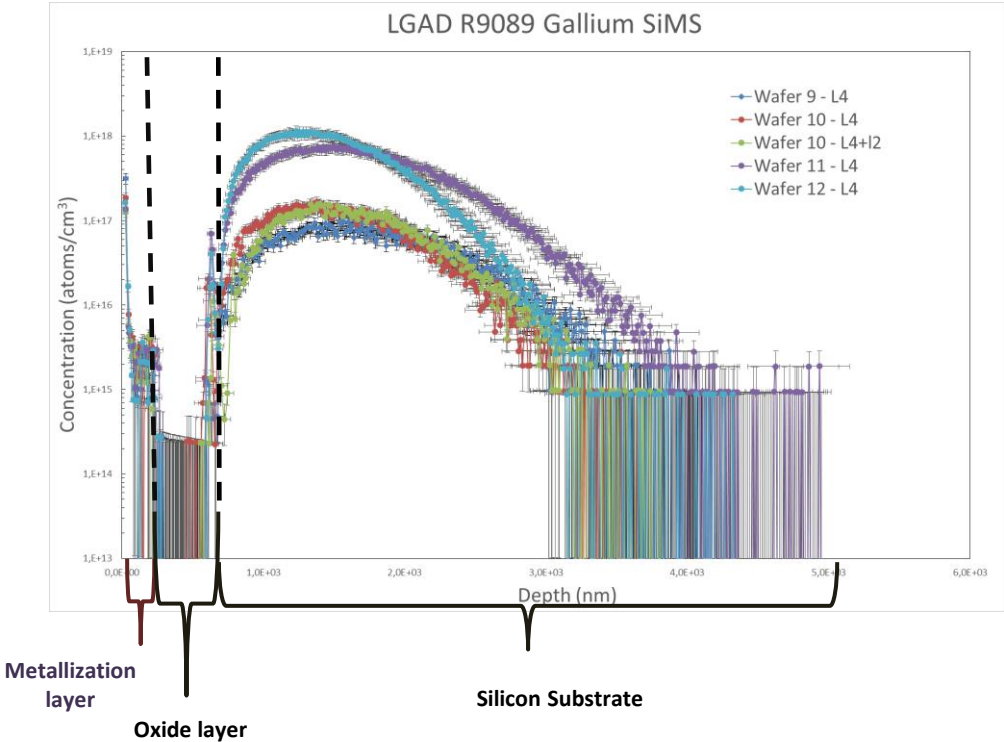
- No degradation in leakage current
- 15% degradation on available headroom in Carbon samples
- 15% degradation in stability of Carbon samples
- No effect on signal properties, efficiency, noise or timing
- In and Li co-implantation as next steps on defect engineering

•Backup



•Standard Candle Process

Typical Profiles



- ✓ New process optimization require standard profile as a reference
- ✓ Use Secondary Ion Mass Spectroscopy (SIMS)profiles form LGAD gain layers
- ✓ Target Boron and Gallium process (well understood)
- ✓ Accuracy of 1e15/cm³

Gallium Nominal Parameters			
Nominal Dose [atoms/cm ²]	Annealing Temp [°C]	Annealing Time [min]	Implant. Energy [KeV]
1,00E+14	1100	180	195
1,00E+14	1100	100	
1,00E+14	1100	100	
1,00E+15	1100	180	
1,00E+15	1100	100	

Target optimization parameters

- ✓ Implantation energy
- ✓ Implantation dose
- ✓ Screen oxide layer thickness
- ✓ Diffusion Time
- ✓ Tilt Angle

•Carbon Calibration Profiles

Sensitivity Optimization

- ✓ The implant concentration is estimated in each case following:

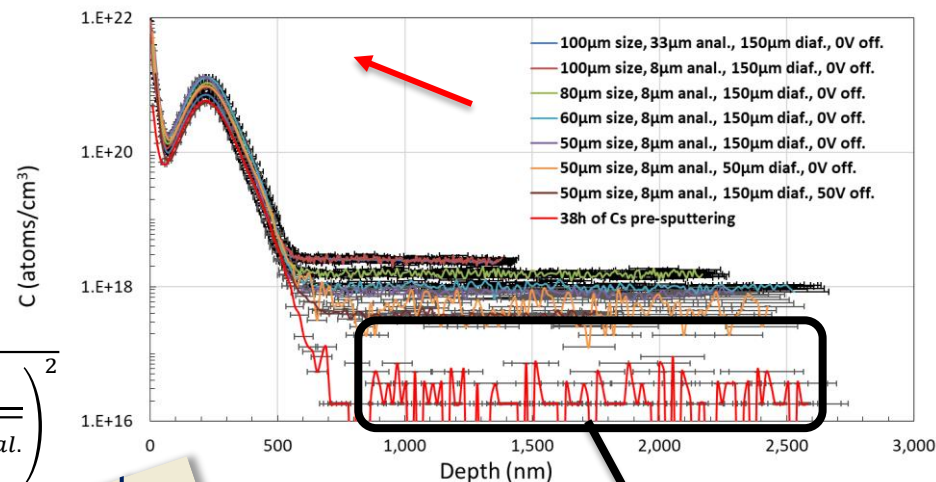
$$C = RSF \times \frac{i_i^{cal.}}{i_M^{cal.}}$$

$$\delta C = \sqrt{\left(\frac{i_i^{cal.}}{i_M^{cal.}} \times \delta RSF\right)^2 + \left(\frac{RSF}{i_M^{cal.}} \times \frac{1}{\sqrt{i_i^{cal.}}}\right)^2 + \left(RSF \times \frac{i_i^{cal.}}{(i_M^{cal.})^2} \times \frac{1}{\sqrt{i_M^{cal.}}}\right)^2}$$

- ✓ Without any additional optimization, a resolution of $(4.71 \pm 0.03) \times 10^{16}$ atoms/cm³ can be achieved
- ✓ The resolution increased for smaller raster sizes while maintain same beam intensity, resulting in higher observant signal intensity
- ✓ Downside of such an approach higher abrasion speed, lees points
- ✓ In essence this is the equivalent in measurement terms of statistical smoothing of profiles.
- ✓ Points recorded every 17 nm, limit of feature size one can probe for achieving such resolution

Best Recorded Resolution for SiMS on Carbon

Carbon Calibration Profile



Gaussian fit on point projection to estimate resolution form σ

Beam Parameters	Abrasion Speed v (nm/sec)	Scaling Factor RSF (atoms/cm ³)	Sensitivity S (atoms/cm ³)
100μm size, 33μm reg., 150μm dia., 0V off.	4.35 ± 0.20	$(2.77 \pm 0.06) \times 10^{22}$	$(4.85 \pm 0.11) \times 10^{16}$
100μm size, 8μm reg., 150μm dia., 0V off.	4.43 ± 0.21	$(3.61 \pm 0.08) \times 10^{22}$	$(1.80 \pm 0.01) \times 10^{17}$
80μm size, 8μm reg., 150μm dia., 0V off.	6.93 ± 0.34	$(2.62 \pm 0.06) \times 10^{22}$	$(1.48 \pm 0.005) \times 10^{17}$
60μm size, 8μm reg., 150μm dia., 0V off.	12.11 ± 0.65	$(1.82 \pm 0.04) \times 10^{22}$	$(7.89 \pm 0.02) \times 10^{16}$
50μm size, 8μm reg., 150μm dia., 0V off.	14.64 ± 0.84	$(1.45 \pm 0.03) \times 10^{22}$	$(7.44 \pm 0.01) \times 10^{16}$
50μm size, 8μm reg., 50μm dia., 0V off.	15.55 ± 0.91	$(1.05 \pm 0.02) \times 10^{22}$	$(1.78 \pm 0.002) \times 10^{17}$
50μm size, 8μm reg., 150μm dia., 50V off.	17.00 ± 1.04	$(5.56 \pm 0.14) \times 10^{21}$	$(4.71 \pm 0.03) \times 10^{16}$
38 h Cesium Pre - Sputtering	16.79 ± 1.02	$(6.04 \pm 0.14) \times 10^{21}$	$(2.22 \pm 0.07) \times 10^{16}$

•Carbon implant simulation

Dopant Defect
Cluster Models

Complex Cluster and BIC (boron interstitial) models

✓ Boron activation model:

- ✓ Boron activation is mainly interstitial driven
- ✓ BIC (Boron Interstitial Cluster) model simulates the process via clustering reactions:
$$\begin{aligned} \mathbf{B_i I_j} + \mathbf{V/I} &\rightarrow \mathbf{B_i I_{j-1}} / \mathbf{B_i I_{j+1}} \\ \mathbf{B_i I_j} + \mathbf{BI} &\rightarrow \mathbf{B_{i+1} I_{j+1}} \end{aligned}$$
- ✓ User defined cluster sizes to consider: **B, BI, BI₂, B₂I₁, B₃I₁, B₃I₂**
- ✓ Reaction rates can be set by user for each reaction (eg 0.3×10^{-10})

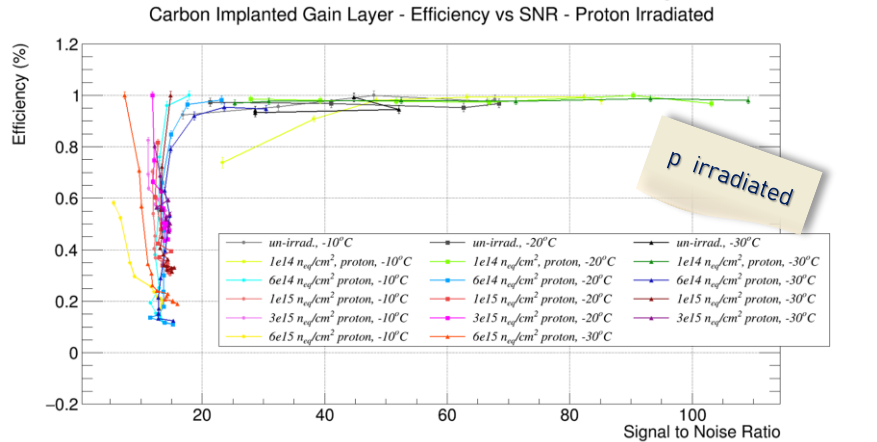
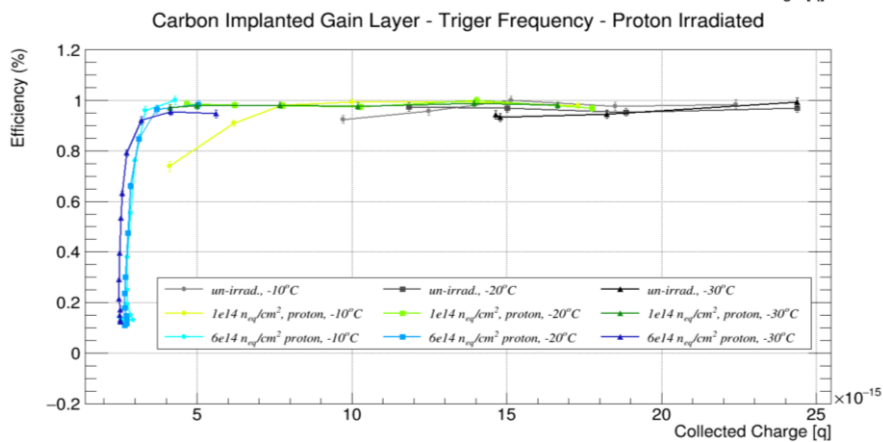
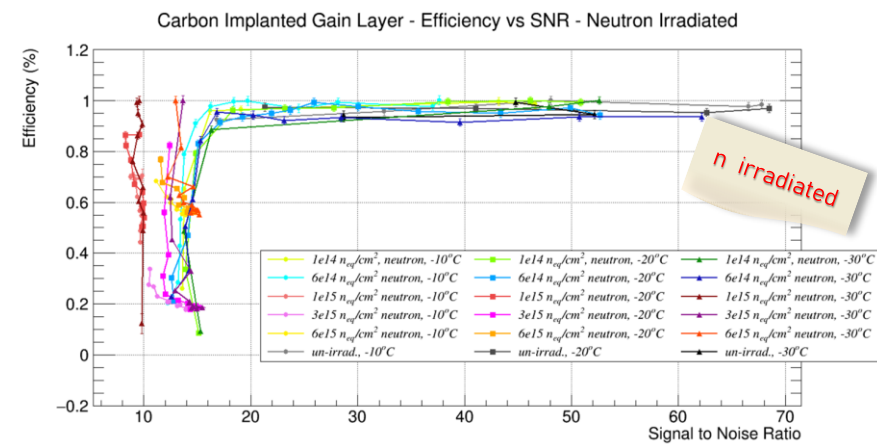
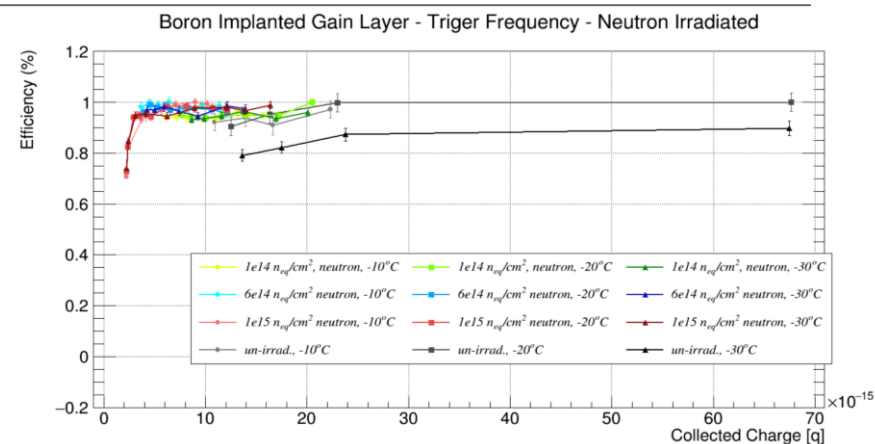
✓ Carbon activation model:

- ✓ The Carbon Cluster or Neutral Cluster Model sets initial cluster concentrations to 0 unless in amorphous regions
- ✓ No charged clusters are considered, solutions to $A_i I_j + I \leftrightarrow A_i I_{j+1}$ $A_i I_j + AI \leftrightarrow A_{i+1} I_{j+1}$ $A_i I_j + V \leftrightarrow A_i I_{j-1}$
- ✓ For Carbon, the following dedicated clusters are computed: C₃I₂, C₄I₂, C₄I₃, C₅I₃, C₅I₄

✓ Boron/Carbon activation/deactivation models:

- ✓ The ComplexCluster Model considers cluster formation between dopants and Vacancies / Interstitials in Si
- ✓ Such process can be described generally as: $\mathbf{n_1 \times Imp.A + n_2 \times Imp.B + n_3 \times V/I + n_4 \times e^- \rightarrow A_{n1} B_{n2} (V/I) n_5 + n_6 e^-}$
- ✓ In the carbon/boron case, the simplest reaction to consider is: $\mathbf{C + B + I \rightarrow BCI + e}$
- ✓ A final charge of 1.0 is expected in such a case
- ✓ For the moment using the Initial concentration as provided after MC implantation by Crystal Trim

•Comparative Studies - Efficiency



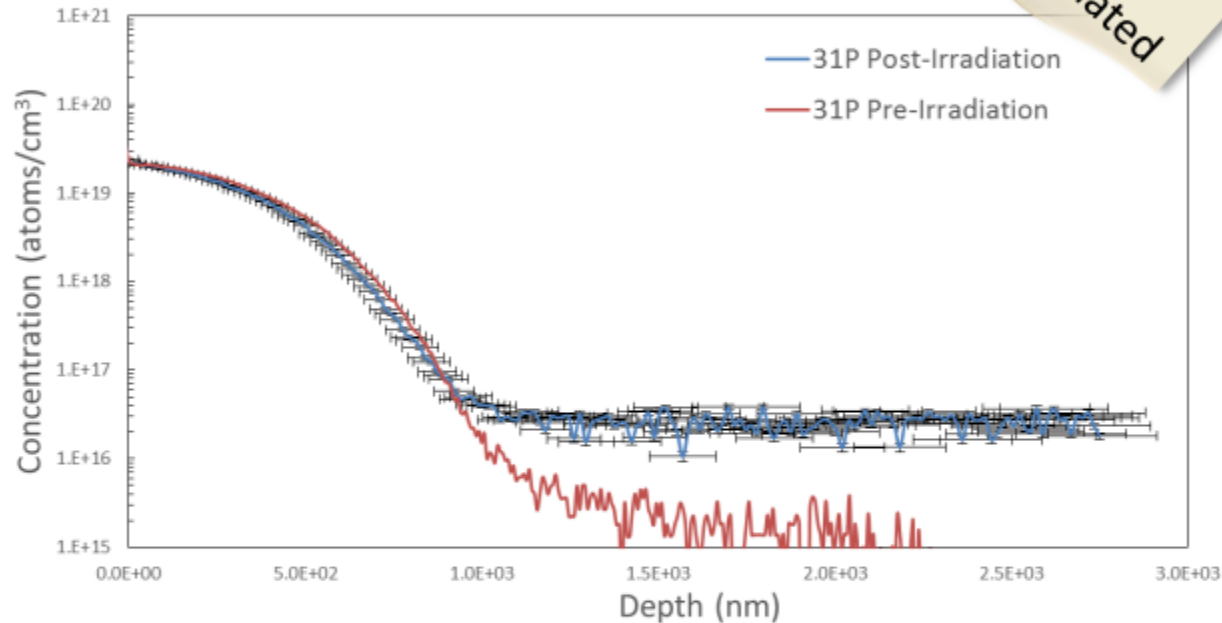
•Phosphorus Test structures

Cis n-in-n $10^{15}/\text{cm}^2$ @130keV

CiS n Implant $10^{15}/\text{cm}^2$, 130keV, 100nm

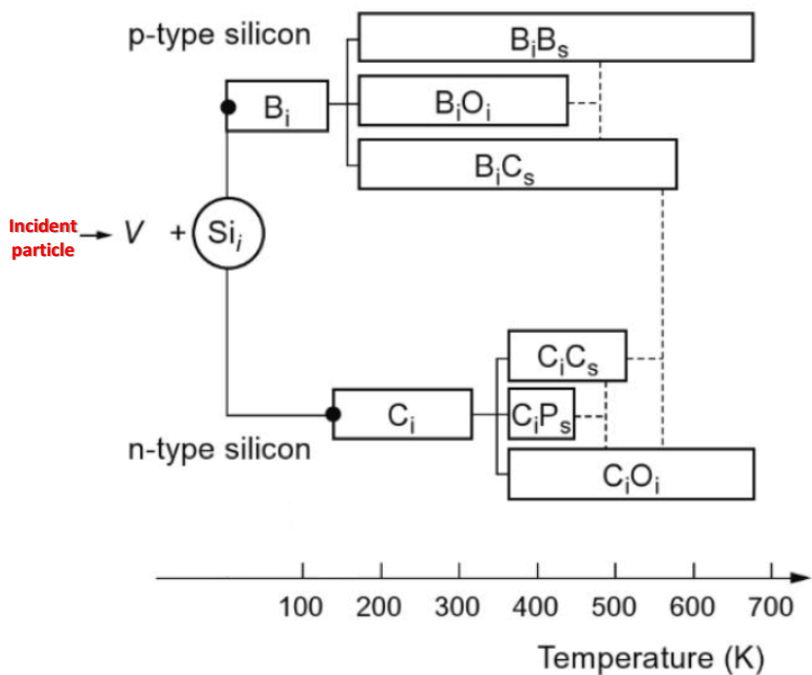
JIS
Irradiated

- ✓ Fluence of $10^{16} n_{eq}/\text{cm}^2$
- ✓ Thermal neutrons
- ✓ Cooled during storage and transport



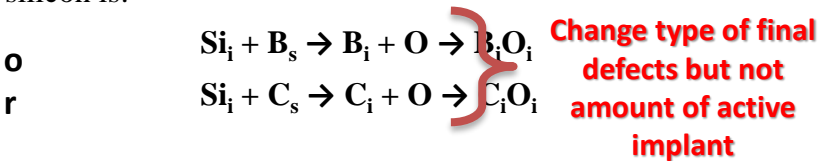
- The doping profile of the neutron irradiated samples seems unaffected
- Agreement within uncertainties
- Higher detection limit due to timing constraints induces deviations at lower part of the profiles

•Radiation Effects



Acceptor removal, Defect Kinetics (simplified ☺)

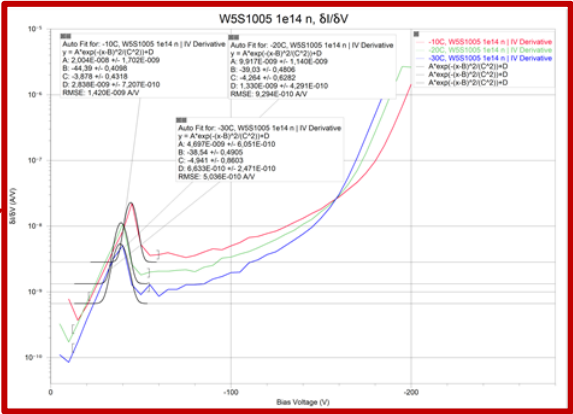
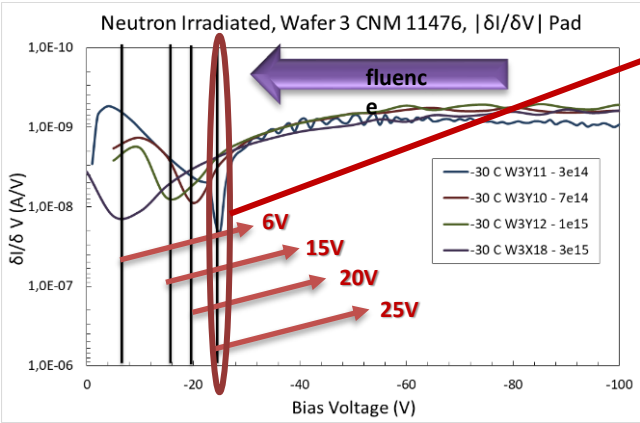
- Incident particle hits silicon atom and created Vacancy (V) and Interstitial Silicon (Si_i)
- Si_i Propagates and can transform substitutional Boron/Carbon to B_i/C_i (interstitial),
- B_i/C_i can form several defects, but the most prominent in high resistivity silicon is:



- Since B_i and C_i both compete for the same Si_i , if we introduce more Carbon we would expect to from less $B_i O_i$ defects and more $C_i O_i$
- If we exchange Boron with a less mobile (heavier) atom (Ga), then we should also enhance $C_i O_i$ defects instead of $Ga_i O_i$

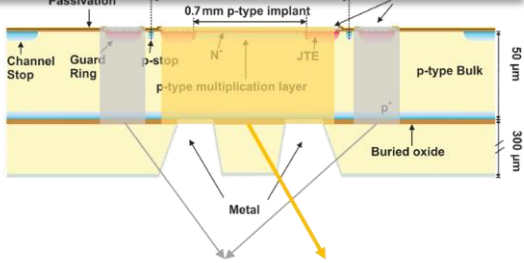
The Derive and Fit Method - I

- ✓ Probe active implant by depletion voltage
- ✓ Additional p-implantation gain layer creates secondary depletion region
- ✓ Mott–Schottky equation → **leakage current variation at gain layer depletion**
- ✓ Form of $|\partial I/\partial V|$ at depletion point corresponds to dopant transition function convoluted with instrument resolution (Gaussian X Gaussian)
- ✓ Depletion voltage determined Gaussian fit at depletion voltage for -10°C, -20°C & -30°C



- Independent Gaussian fits for each temperature
- Uncertainties estimated from propagation of fit sigma
- Fluences up to $3 \cdot 10^{15} \text{ n}_{eq}/\text{cm}^2$ in p^+ and n^0

Gkougkousis V., RD50 Workshop Talk, November 2019:
[link](#)

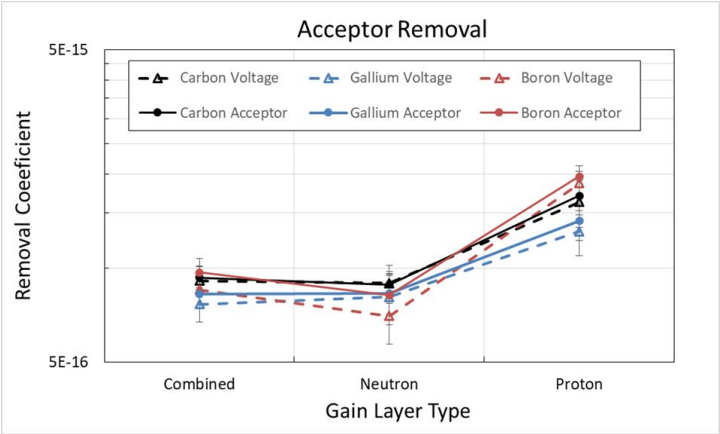
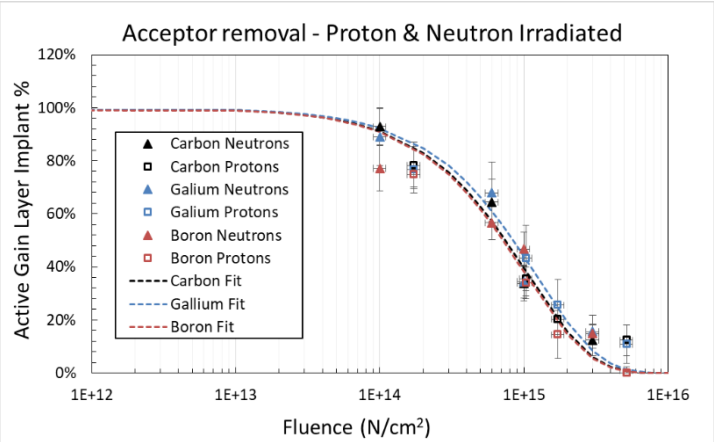


$$V_d = \frac{\sum_{T=-10}^{-30} V_{d,T_i}}{n_T}$$

$$\delta V_d = \sqrt{V_{d.sys} + V_{d.stat}}$$

Average of fit sigma Standard deviation of V_d

The Derive and Fit Method - II



- Linear dependence assumption between V_{GL} and active implant
- Normalized exponential reduction fit model on gain and V_{GL}

$$G(\%) = e^{-C_G \Phi}$$

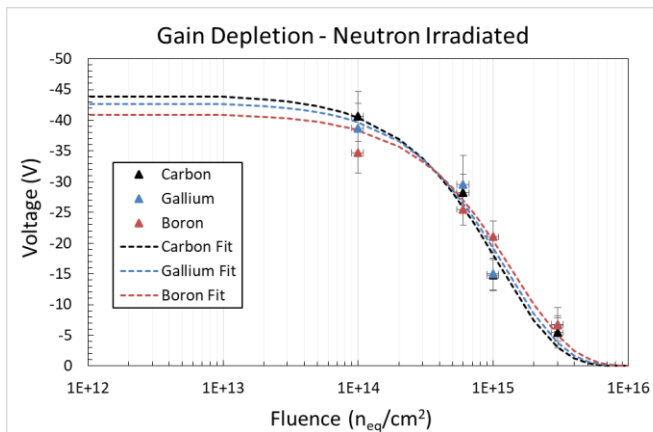
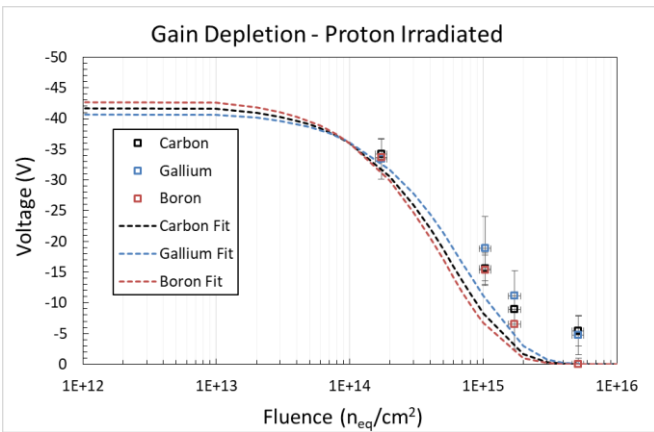
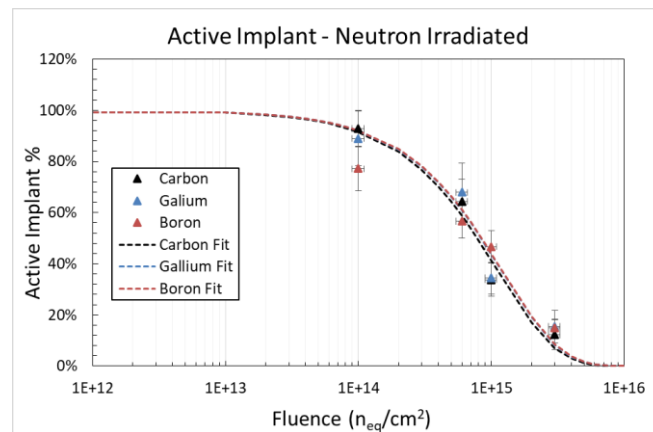
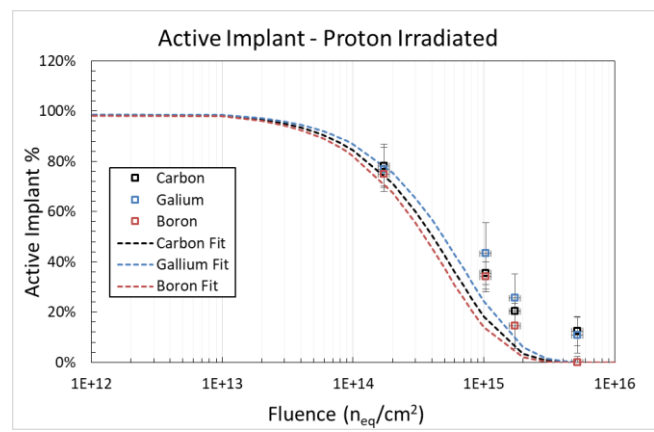
- Linearity hypothesis tested with independent C_v and C_G fits – full compatibility
- Constraints imposed on initial values to reflect charge measurements

Results

- Compatible acceptor removal coefficients between all implants
- Slight Ga advantage in p^+ irradiation (23 GeV/c PS), higher mass reduces displacement probability in coulomb-only (far-field) interactions
- Quasi-identical performance for neutron irradiated (fast ~ 10 MeV neutrons)
- Identical gain layer de-activation for all dopants with fluence**

Acceptor Removal Coefficient		
Irrad. Type	C	δC
Gallium		
Combined	8.25E-16	7.98E-17
n^0 irradiated	8.28E-16	1.16E-16
p^+ irradiated	1.41E-15	1.88E-16
Boron + Carbon		
Combined	9.33E-16	7.78E-17
n^0 irradiated	8.85E-16	8.76E-17
p^+ irradiated	1.70E-15	2.23E-16
Standard Boron		
Combined	9.69E-16	1.04E-16
n^0 irradiated	8.19E-16	1.35E-16
p^+ irradiated	1.96E-15	1.60E-16

•The Derive and Fit Method - II



•Comparative Studies II - Stability

Self-trigger time:

$$\Delta T_{trig}^i = \frac{\sum_{j=1}^{n-1} (T_{j+1}^{trig} - T_j^{trig})}{n}$$

Self-trigger Rate:

$$F_{trig}^i = \frac{1}{\Delta T_{trig}^i}$$

Median of several rate measurements

$$\widetilde{F_{trig}} = \frac{F_{trig} \lfloor (\#k+1) \div 2 \rfloor + F_{trig} \lceil (\#k+1) \div 2 \rceil}{2}$$

Uncertainty on trigger rate:

$$\delta \widetilde{F_{trig}}(\%) = \sqrt{\frac{(N_{over} + 1) \times (N_{over} + 2)}{(N + 2) \times (N + 3)} - \frac{(N_{over} + 1)^2}{(N + 2)^2}}$$

Efficiency is a binary magnitude, Bayesian approach implemented

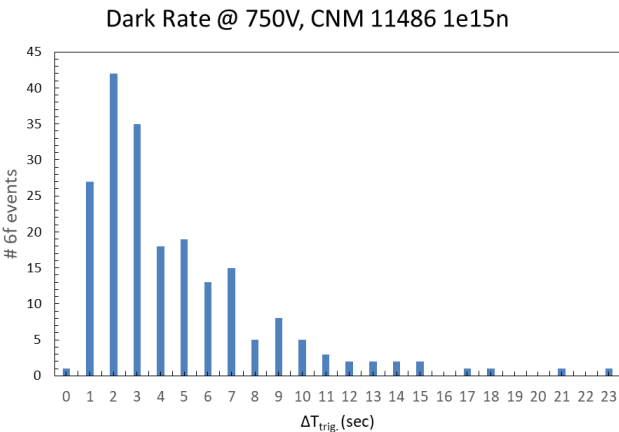
Sigmoid Dark rate Fit:

$$R_{Dark\ Rate} = \frac{R_{max}}{1 + e^{C \times (V_{50\%} - V)}} + R_{BaseLine}$$

Max, recordable rate
Inst. saturation point

50% of maximum voltage point

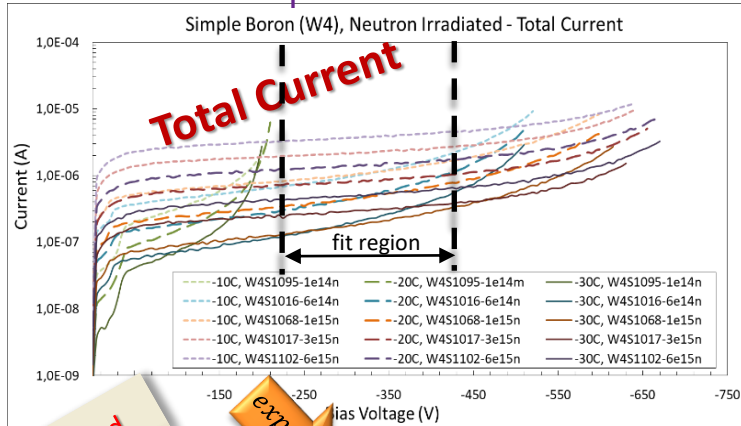
Baseline trigger rate
(noise, radioactivity)



- ✓ Sensors with intrinsic gain present dark rate at higher biases
- ✓ Brownian thermal electrons following Poisson distribution
- ✓ As gain increases, the amount of charge necessary for an event to cross trigger threshold decreases
- ✓ Shot thermal noise increases with voltage
- ✓ Evaluation performed at the 2 fC threshold
- ✓ Values estimated from Poissonian fit on event frequency distribution (1000 events)

• Breakdown Voltage

Current Multiplier



Method

exp. fit

Exponential Fit: $I = b \cdot m^V$

Acceptance Criteria: $R^2 \geq 99\%$

Expected current: $I_{norm} = b \cdot m^{V_i}$

Current Multiplier: $M(V) = \left| \frac{I_{pad} + I_{GR}}{I_{norm}} \right|$

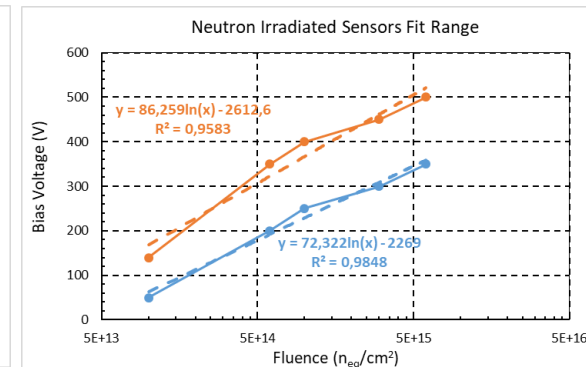
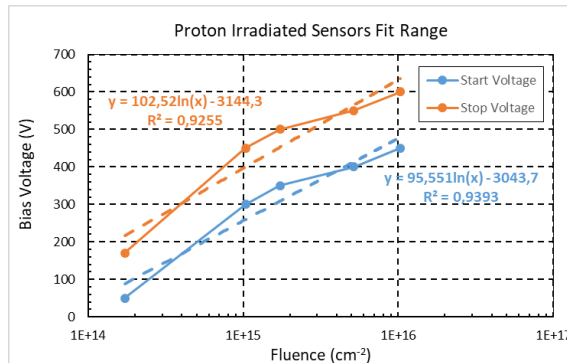
Breakdown: $V_{brw} \rightarrow M(V) > 2$

- ✓ Measure total leakage current (-10°C, -20°C, -30°C)
- ✓ Select a stable voltage range where behaviour follows exponential law
- ✓ Define common for all temperatures stable voltage range, after depletion and much before breakdown
- ✓ Perform exponential fit requesting $R^2 \geq 99\%$ (same range as in the gain reduction fits - same constraints)
- ✓ Calculate the multiplier with respect to the expected current
- ✓ **Define breakdown in multiplier value (Is it really exponential??)**

Un-irradiated: $I_{pad}^{\Phi=0} = I_s \times (e^{\frac{eV}{nkT}} - 1) \times G(e^V, T)$

Irradiated: $I_{pad}(\Phi) = (I_{pad}^{\Phi=0} + \alpha\Phi) \times G^*(e^V, T, \Phi)$

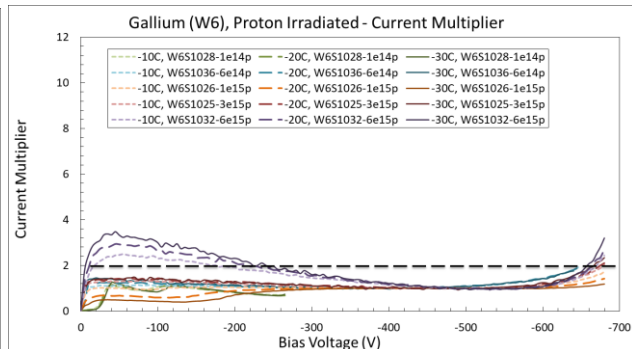
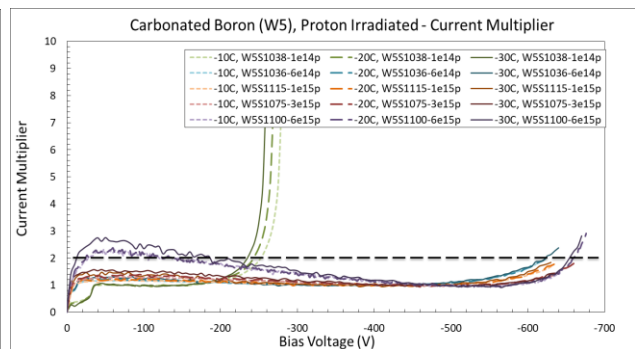
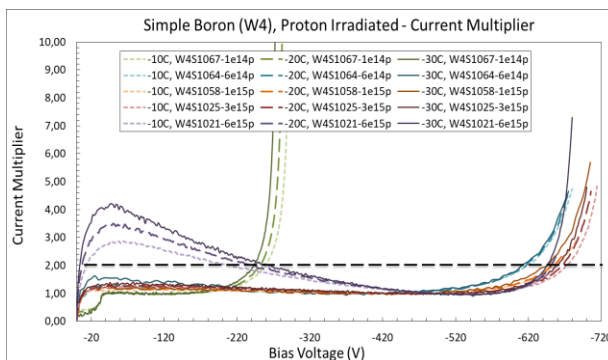
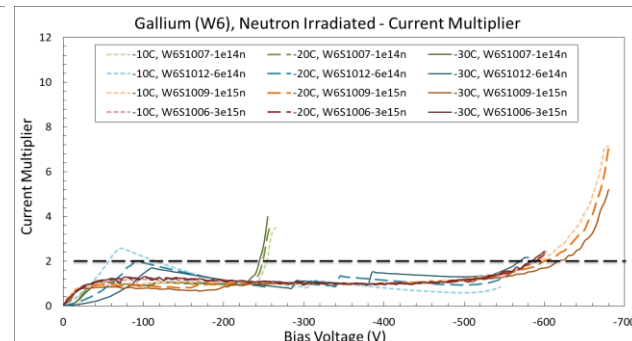
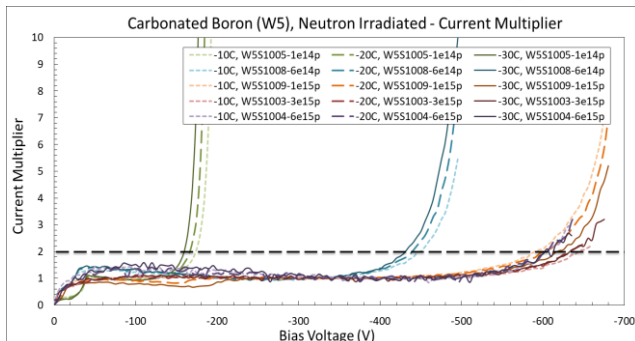
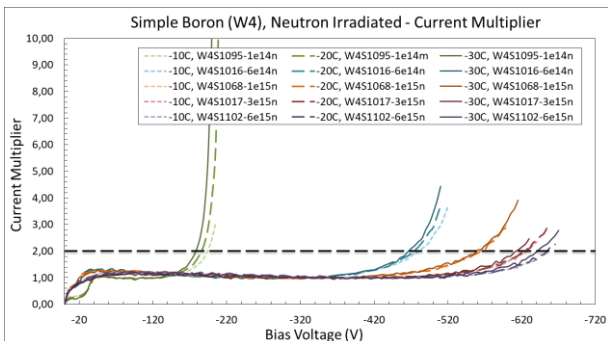
Function of acceptor removal, exponential to fluence and voltage plus a linear term



• Breakdown Voltage

- ✓ Independent fit for each temperature
- ✓ Identical fit regions across all temperatures
- ✓ Identical fit regions for same fluence across all three implants

Constraints

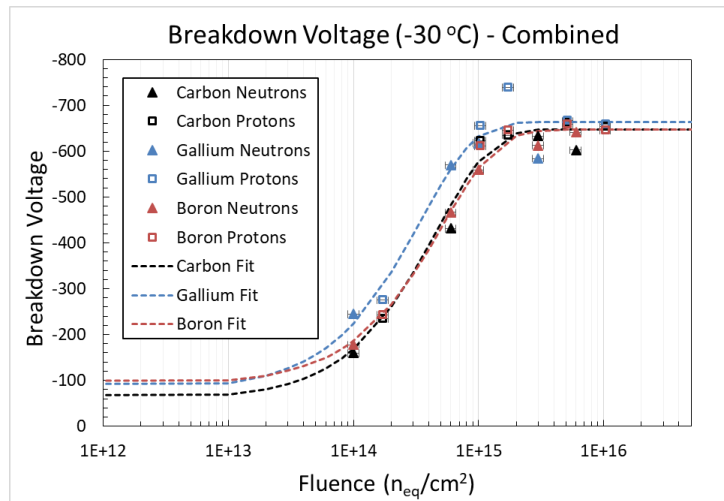


•Breakdown Voltage

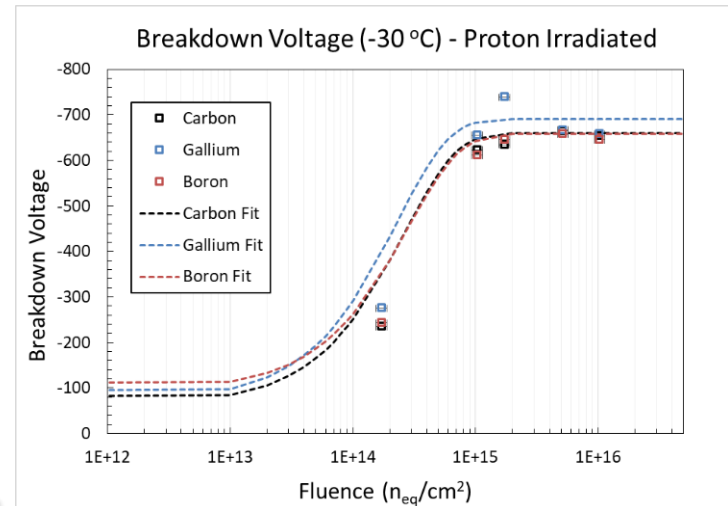
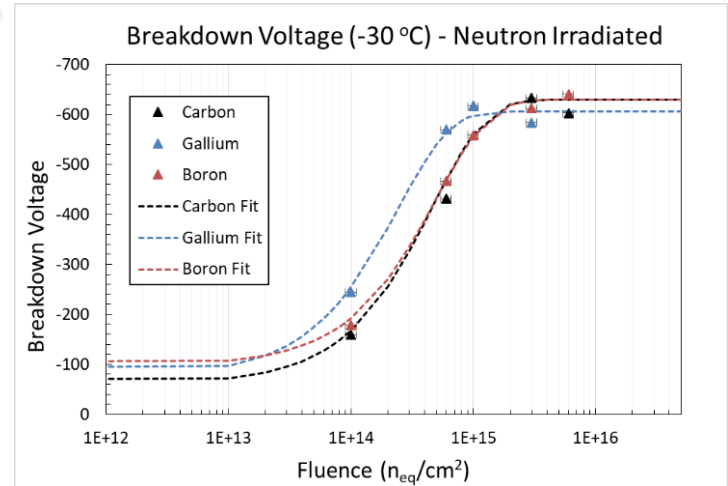
Model

$$V_b = (V_{max} - V_0)(1 - e^{-c\Phi}) + V_0$$

Breakdown of PIN Un-irradiated breakdown voltage



- ✓ Carbon and boron are compatible
- ✓ Gallium presents higher breakdown voltage (most possibly due to process variation)
- ✓ All implants compatible with sigmoid approach
- ✓ Highest breakdown voltage after irradiation independent of gain – exclusively process dependent



• Introduction ii

Timing Concepts

Time Resolution: $\sigma_{tot}^2 = \underbrace{\sigma_{timewalk}^2}_{\sigma_{Dist.}^2 + \sigma_{Landau}^2} + \underbrace{\sigma_{jitter}^2}_{\left(\frac{t_{rise}}{S/N}\right)^2} + \underbrace{\sigma_{conversion}^2}_{\left(\frac{TDC_{bin}}{\sqrt{12}}\right)^2} + \underbrace{\sigma_{Clock}^2}_{\text{Fixed Term } \sim 5-7 \text{ psec}}$

$$\frac{t_{rise}}{S/N} \approx \frac{N}{dV/dt}$$
$$\left[\frac{V_{th}}{S/t_{rise}}\right]_{RMS} \propto \left[\frac{N}{dV/dt}\right]_{RMS}$$

$\left(\left[\frac{V_{thr}}{S/t_{rise}}\right]_{RMS}\right)^2$

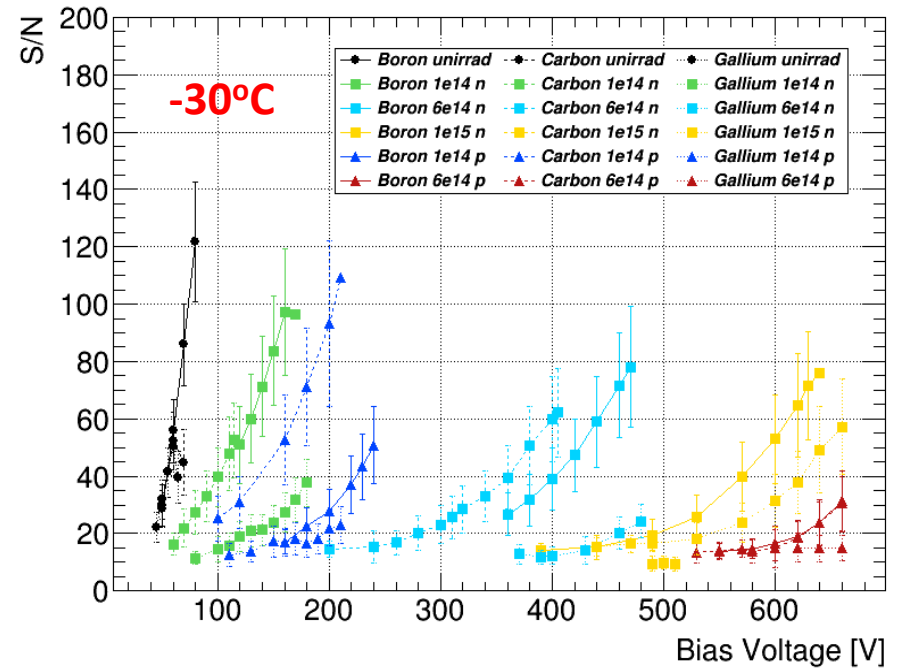
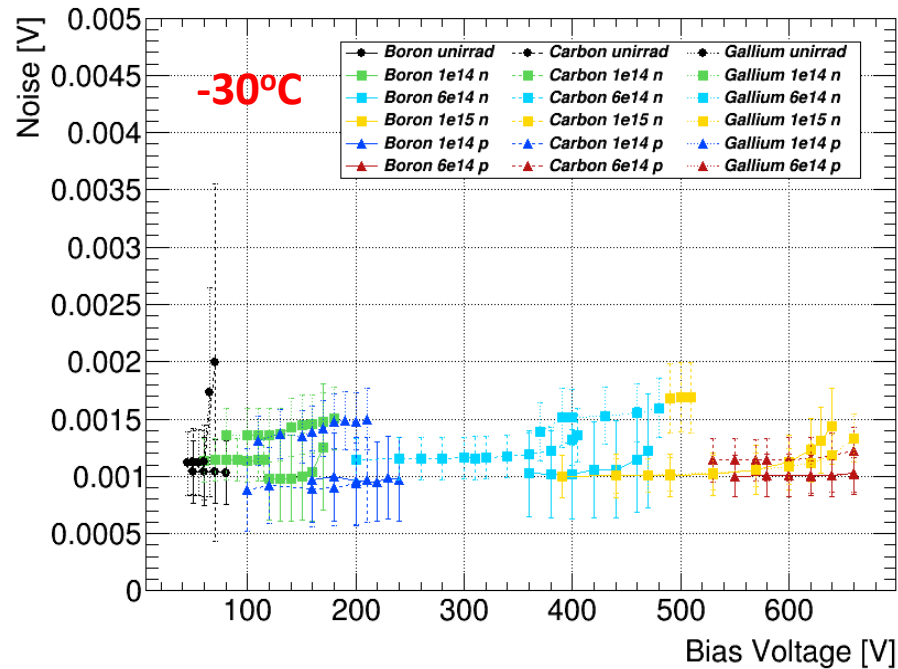
Where: S signal
 N noise
 V_{th} CFD threshold
 t_{rise} rise time

Fast time resolution:

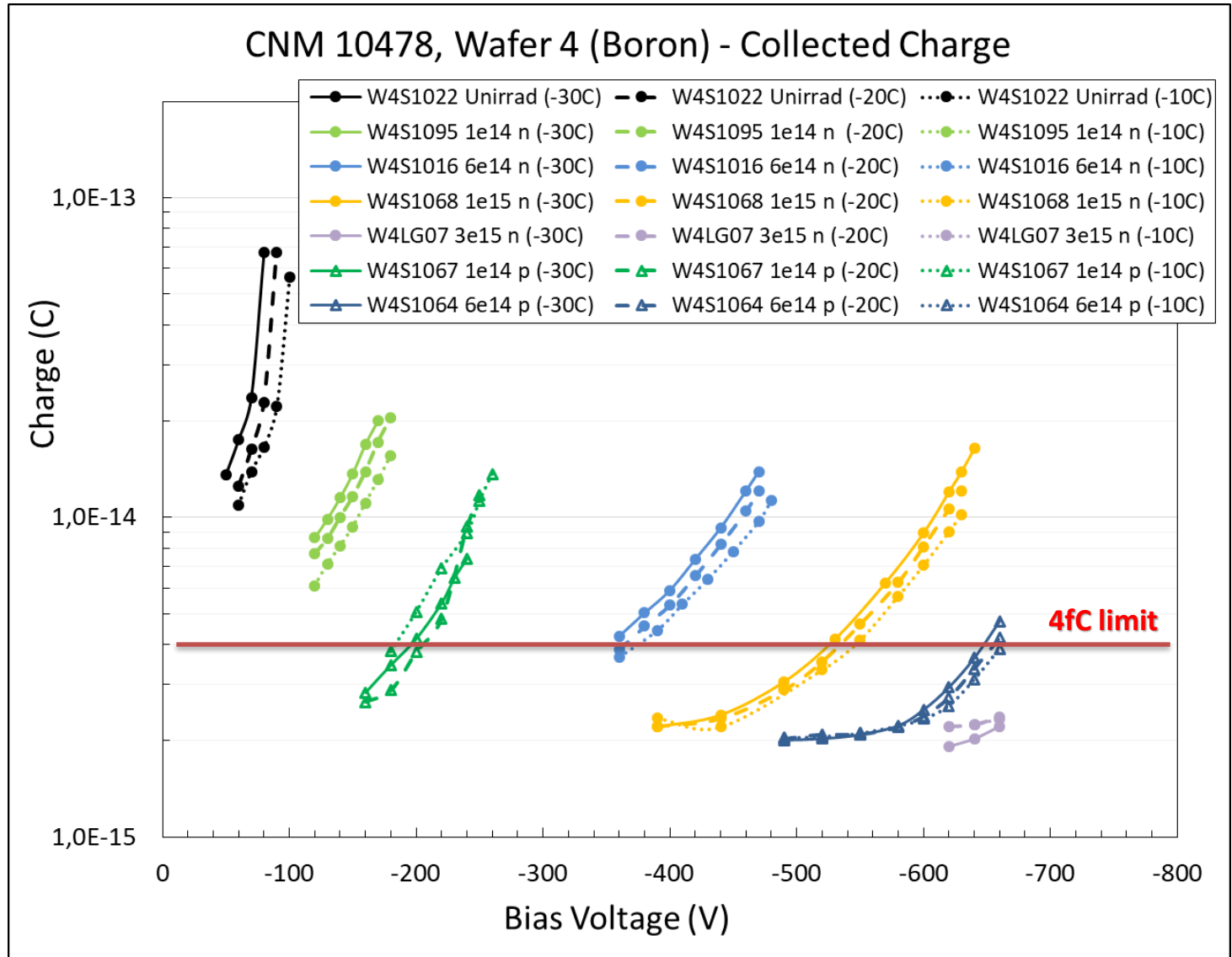
- ✓ Maximize slope (large fast signals)
- ✓ Minimize noise to minimize jitter
- ✓ Implement time walk correction (CFD, ToT, ToA, ect)
- ✓ Uniform field with to minimize distortion term

Thin silicon sensors with internal gain

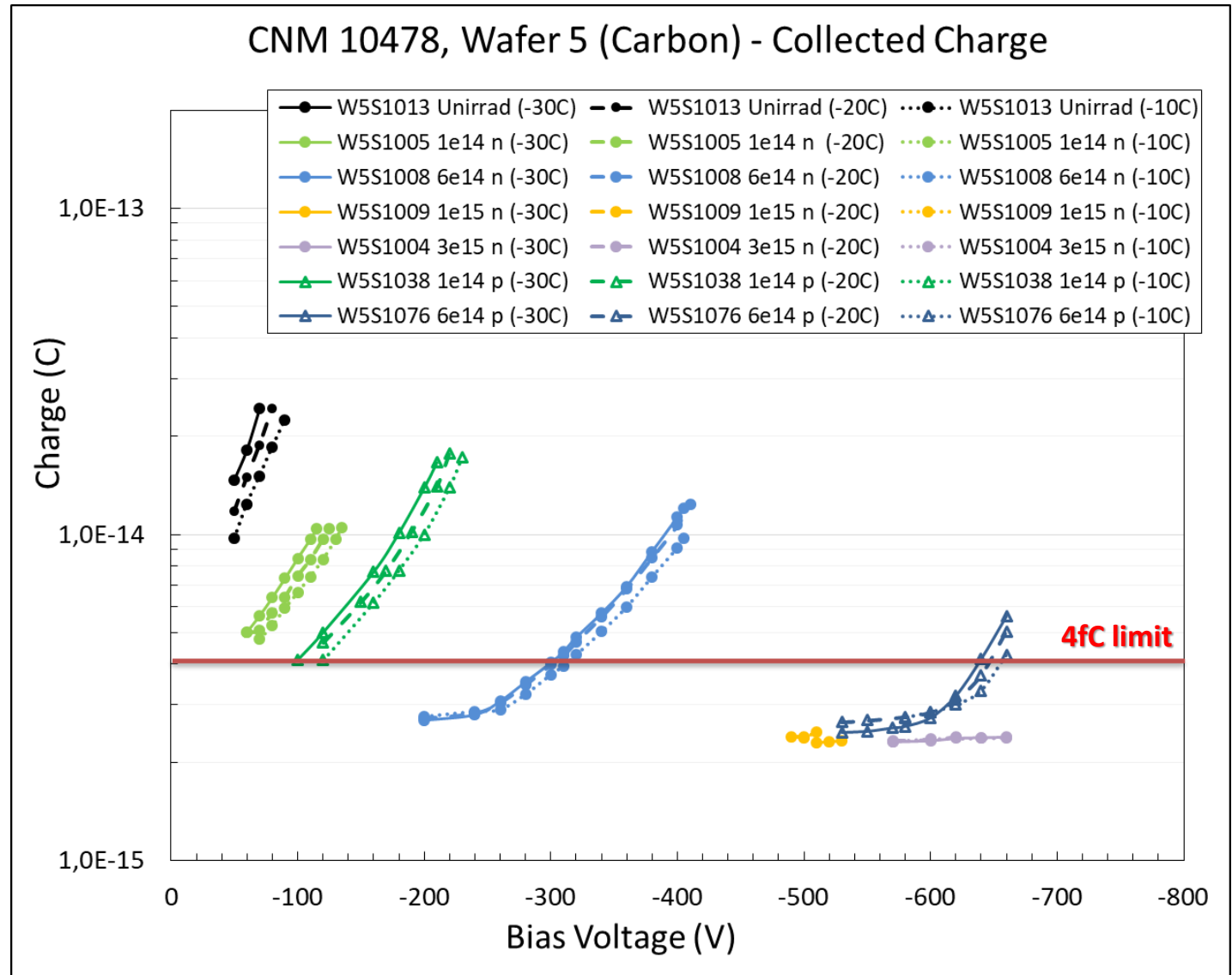
•Noise and S/N Ratio



•Collected Charge - Boron



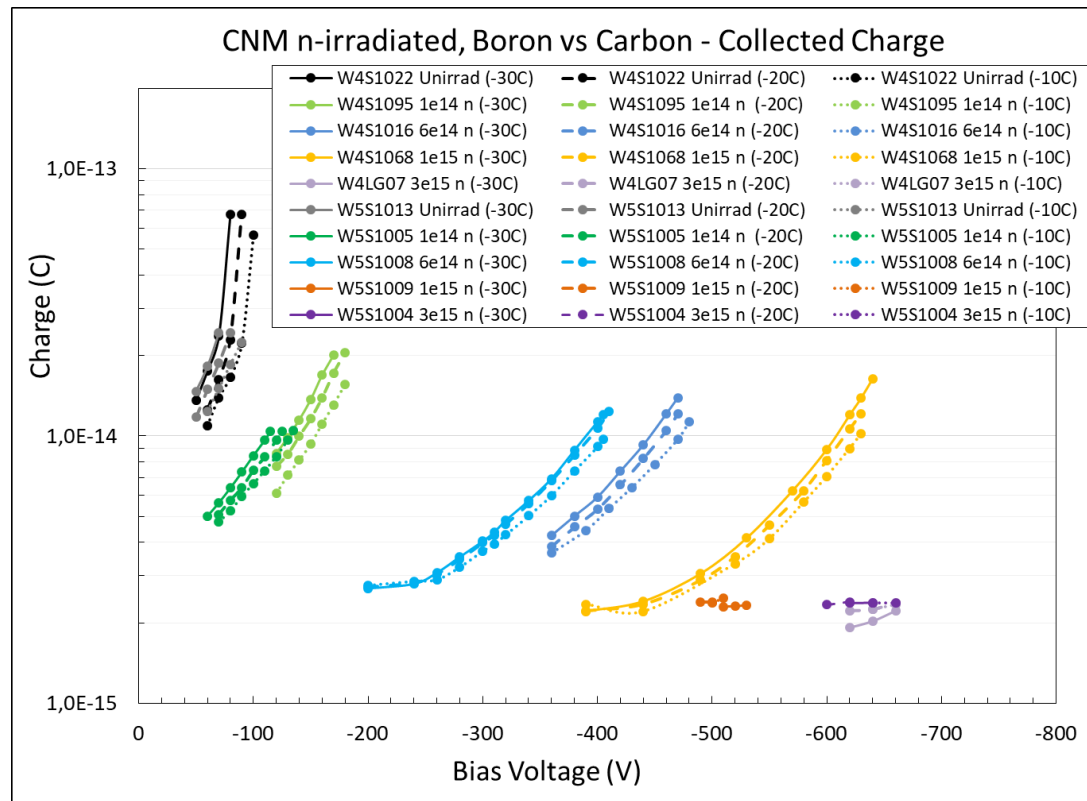
•Collected Charge – Boron + Carbon



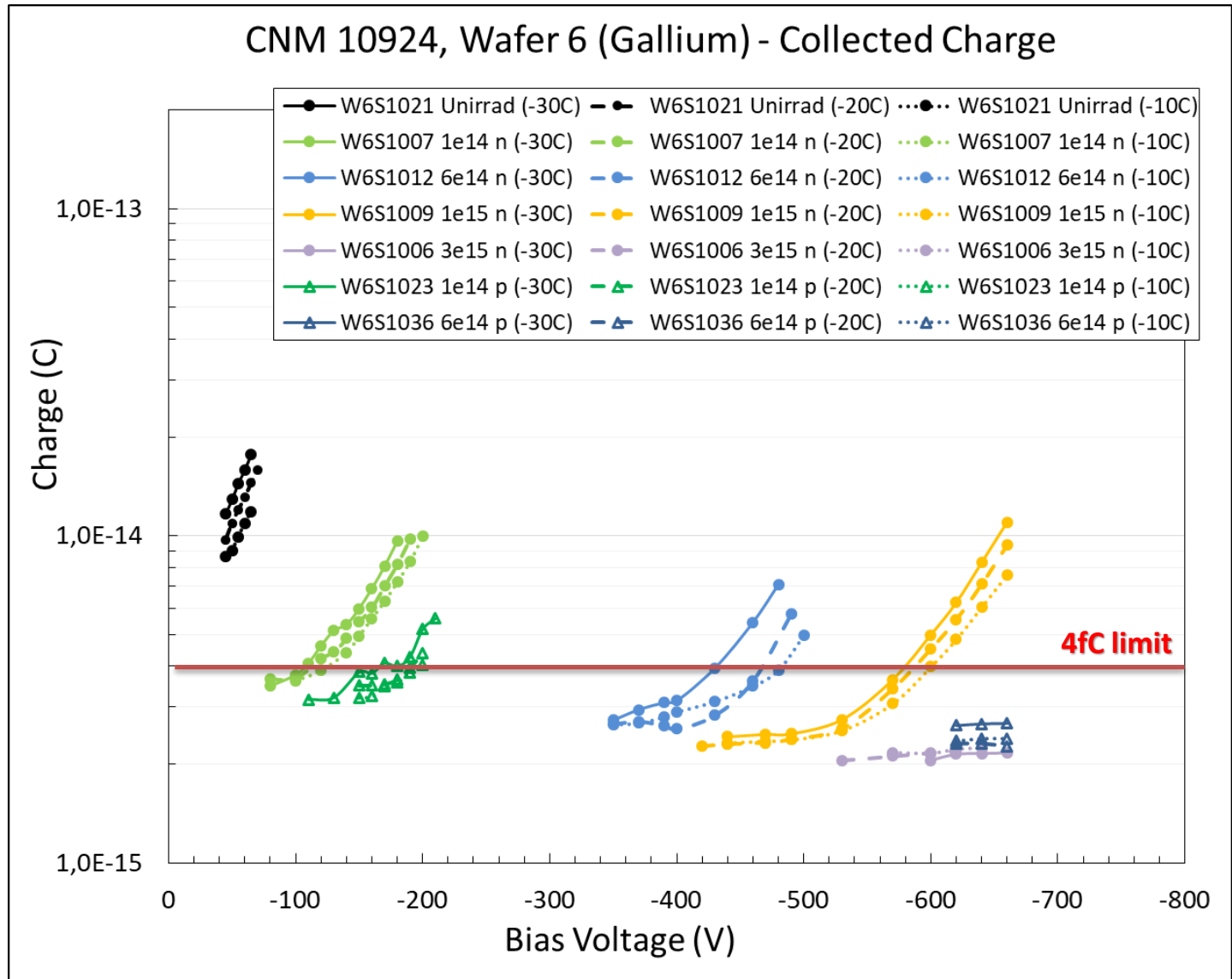
• Sensor R&D

LGADs – Charge collection

- ✓ Unirradiated gain tuned to be identical for boron/gallium/carbon implanted sensors for accurate comparison
- ✓ Irradiated Carbon infused sensors present higher charge at lower HV
- ✓ Gallium implanted sensors are 10% worse than standard process boron
- ✓ Carbon is 20% better across the spectrum with respect to boron



•Collected Charge – Gallium



•Dark Rate

Concepts & Methods

- ✓ Sensors with gain present dark rate at high enough voltages
- ✓ Dark rate events result of thermal movement and random in nature
- ✓ Follow the Poisson distribution

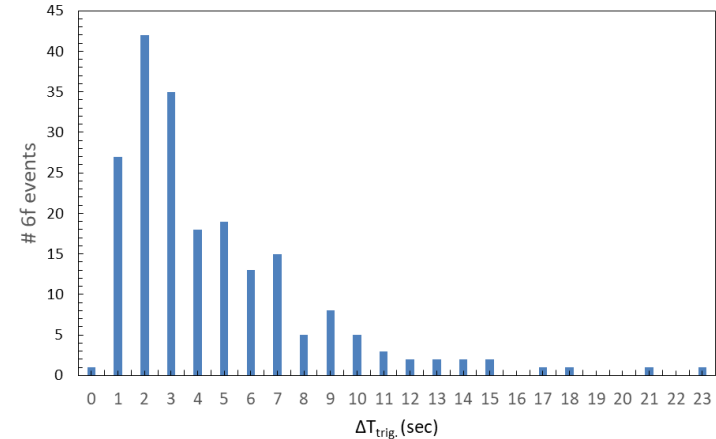
Quantification

- ✓ Study the time between consecutive self-triggering
- ✓ Use mean of 4 events (3 values) to reject cosmic background

Self-trigger time:
$$\Delta T_{trig}^i = \frac{\sum_{j=1}^{n-1} (T_{j+1}^{trig} - T_j^{trig})}{n}$$

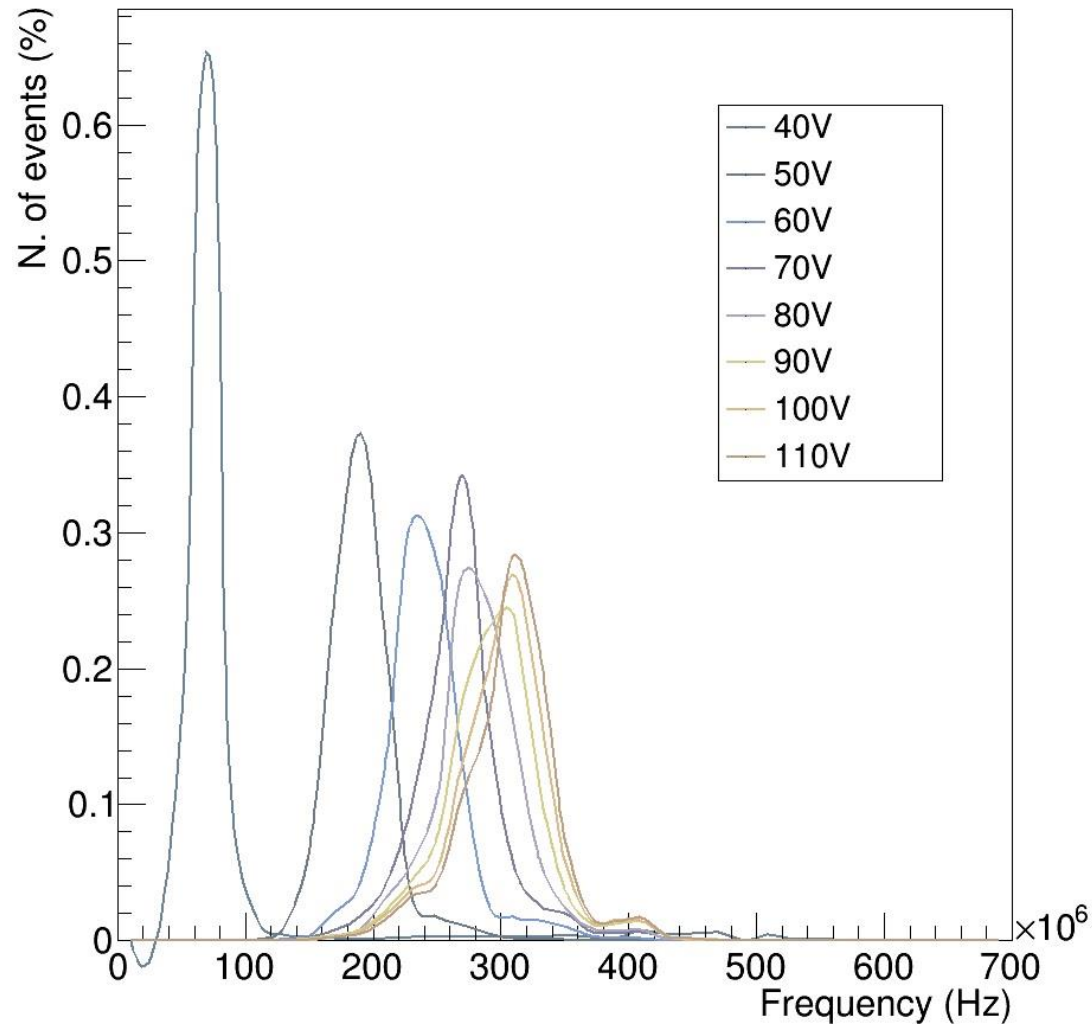
Self-trigger Rate:
$$R_{trig}^i = \frac{1}{\Delta T_{trig}^i}$$

Dark Rate @ 750V, CNM 11486 1e15n



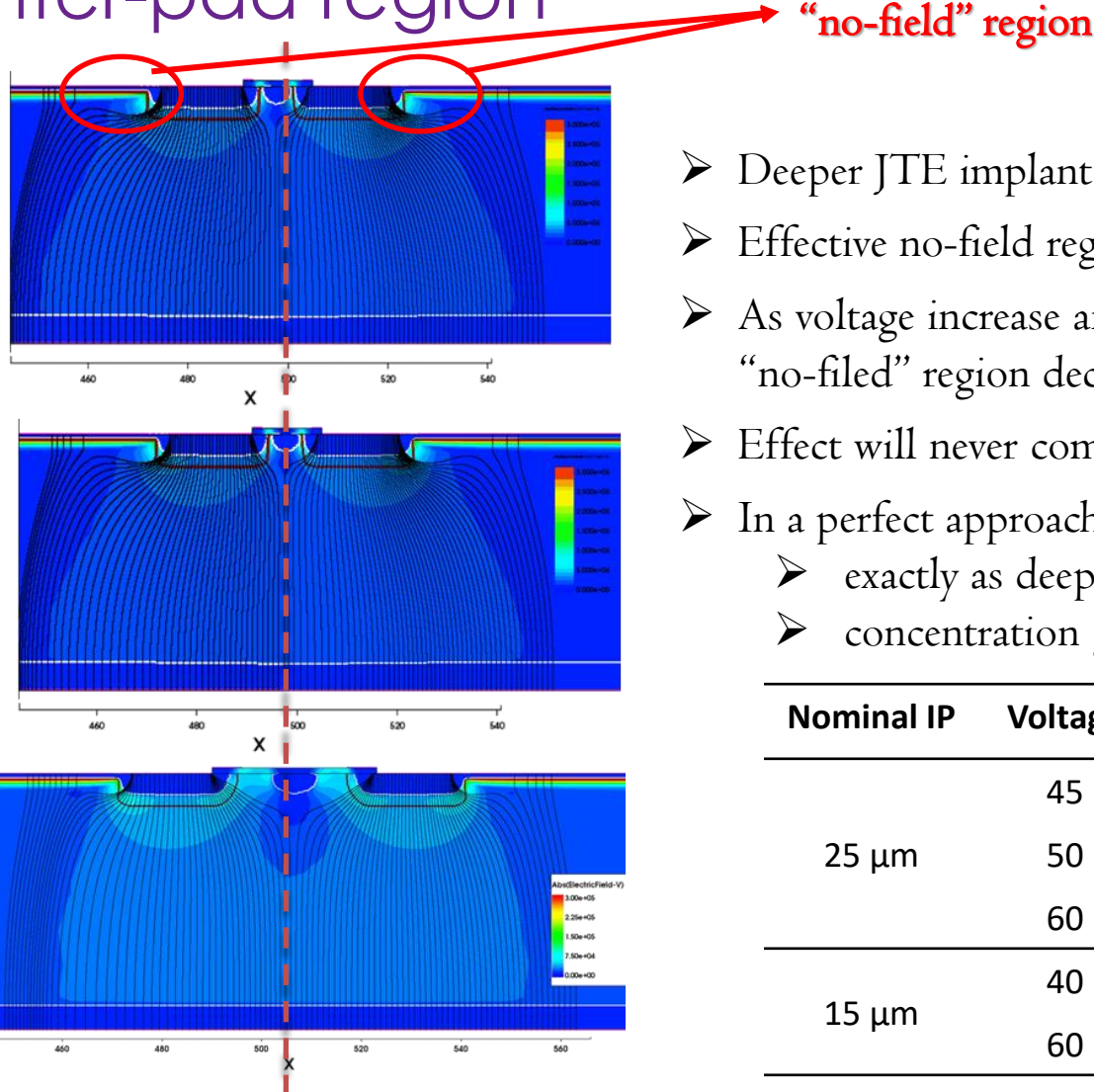
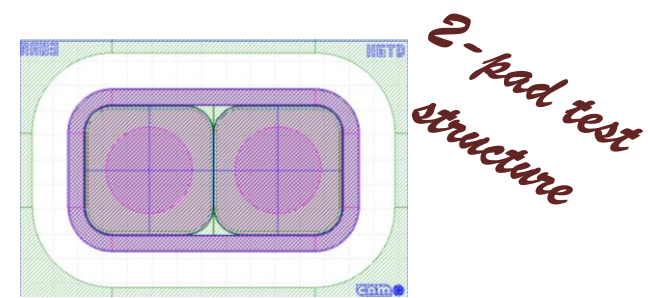
•Signal Evolution with bias in LGADs

Signal FFT - 1e14n, -30C



• Main Development Points

Inter-pad region



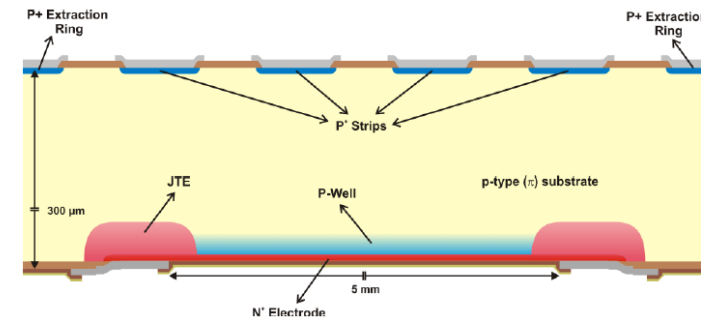
- Deeper JTE implant with respect to multiplication layer
- Effective no-field region created next to the boarder
- As voltage increase and gain field become more significant, “no-filed” region decreases
- Effect will never completely disappear
- In a perfect approach, JTE would be:
 - exactly as deep as gain layer
 - concentration gradient tuned to match gain layer

Nominal IP	Voltage	Measured IP	IP efficiency (%)
25 μm	45	103.0 ± 1.3	2,2
	50	91.8 ± 1.3	10,1
	60	69.2 ± 2.4	17,4
15 μm	40	111.9 ± 4.9	14.8
	60	70.2 ± 4.8	29.4
45 μm	120	68.4 ± 4.6	55.1

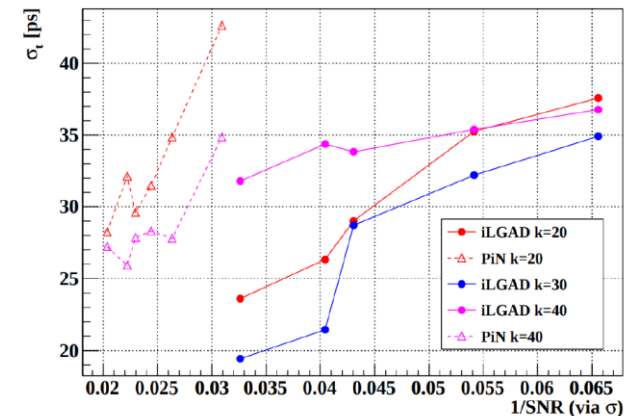
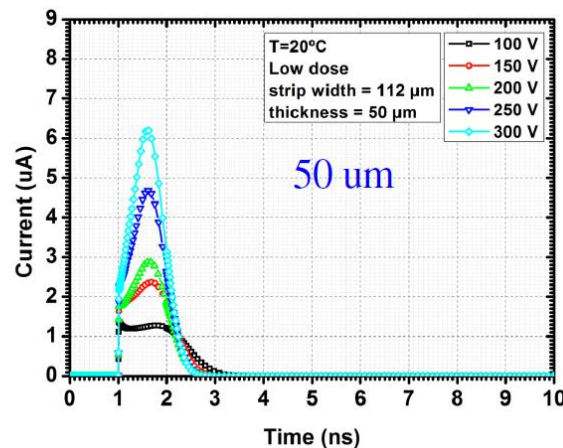
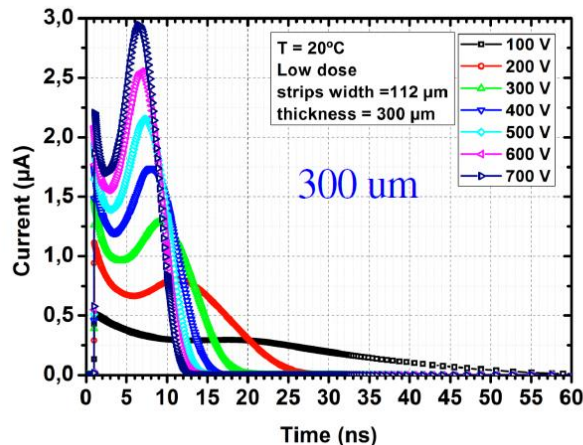
• Geometric efficiency

Inter-pad region, i-LGADs

- Approached based on non-segmentation of gain layer
- Double sided process with **NO** possibility for support structures (thin sensors extremely difficult)
- Multiplication layer multiplies carriers
- Two contribution signal, primary electron collection and subsequent multiplication of holes
- Distance between two signals depends on drift time, thinner sensors ($< 50 \mu\text{m}$) should be usable for time



- Typical pulse duration $\sim 10 - 20$ times of equivalent thickness LGAD
- Very good for timing, bad in SNR terms
- Need too go to $50 \mu\text{m}$ thick devices for any realistic applications



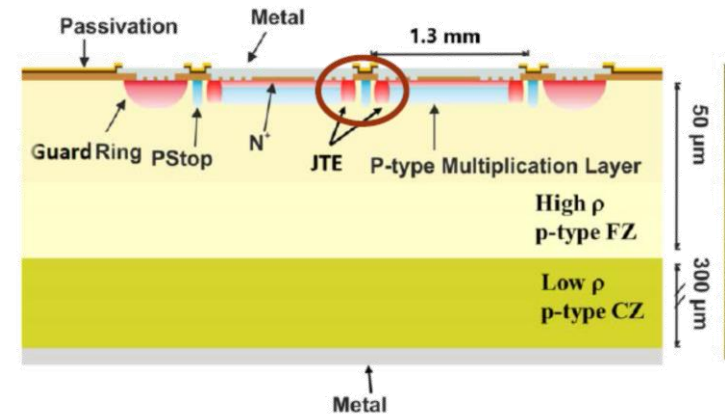
• Inter-pad area

Inactive regions

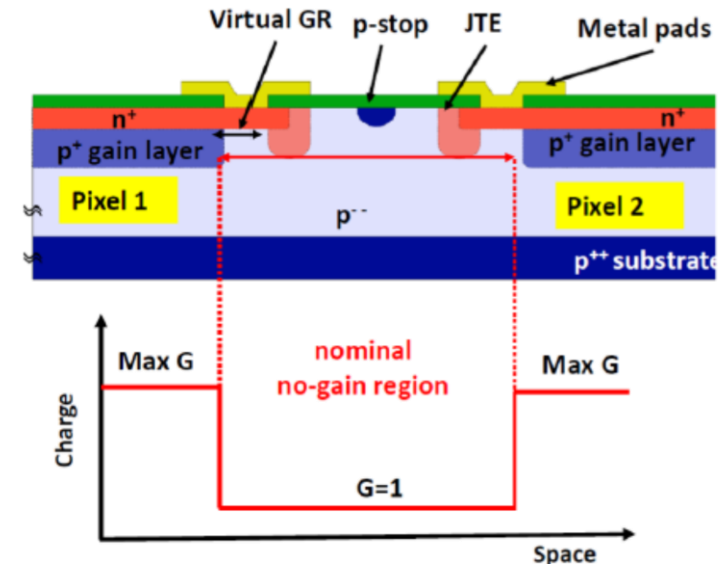
- ✓ High Field region in the gain layer
- ✓ Field needs to be controlled in pad edges where values increase due to geometry
- ✓ Introduction of electrical isolation implant JTE (Junction Termination Extension)
- ✓ Typical JTE geometries introduce 50 – 150 μm inactive area between adjacent gain layers
- ✓ Dead area varies with field (bias voltage) in a non-linear way (see next slide)
- ✓ Overall fill factor reduction at 1mm pads:

JTE size	Fill Factor reduction
50 μm	~ 10 %
150 μm	~ 30 %

CNM Technology



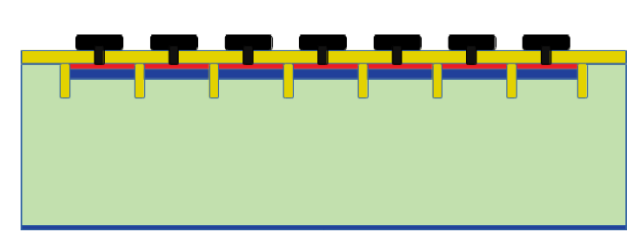
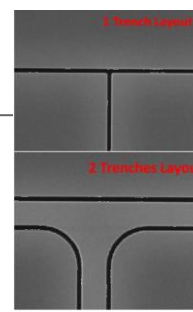
FBK Technology



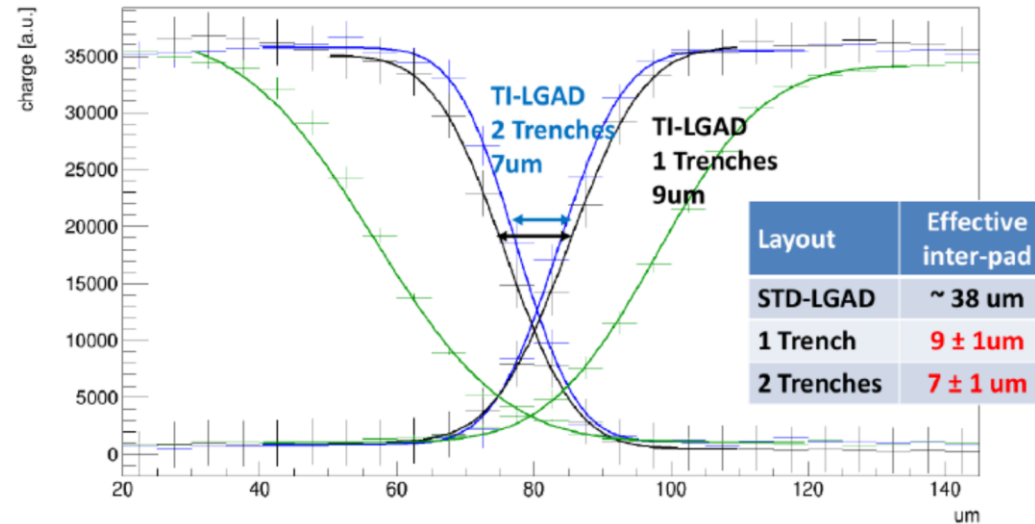
• Inter-pad area

Inter-pad region, Ti-LGADs

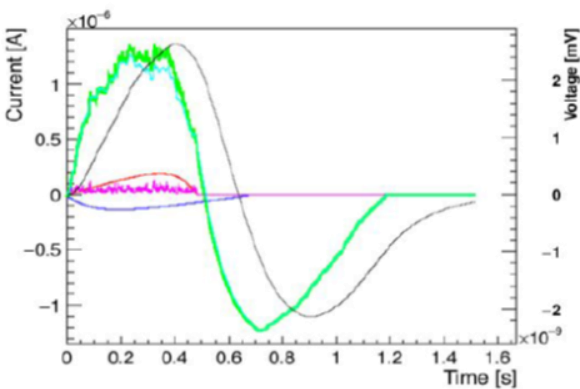
- Approach based on removing JTE completely
- Using trenches as electron diffusion barrier and field containment (JTE and p-stop)
- DRIE trenches comparable to SIPM processes $\sim 1\mu\text{m}$ thick, filled with oxide
- Relatively understood process
- Based on the distance from the edge an opposite polarity signal present on adjacent pad (probably dependent on trench depth)



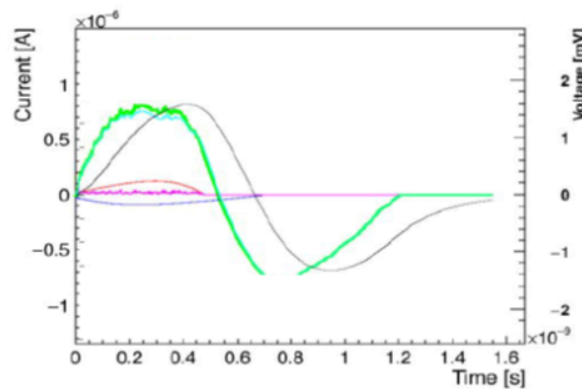
Laser characterization



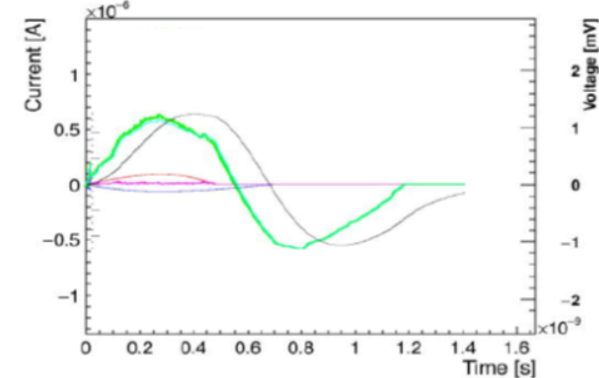
(a) $35\mu\text{m}$ from mid gap



(b) $65\mu\text{m}$ from mid gap



(c) $95\mu\text{m}$ from mid gap



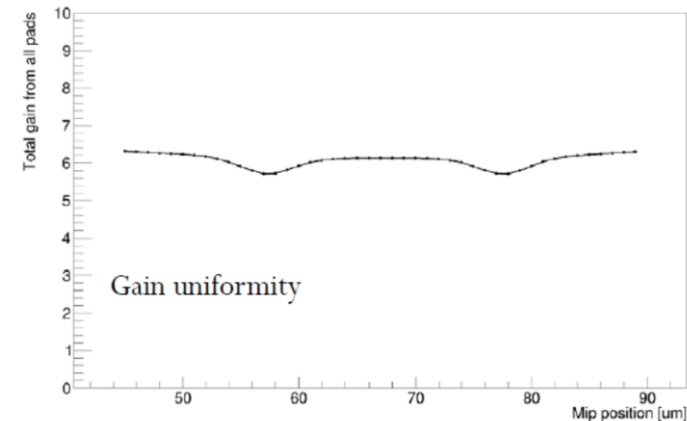
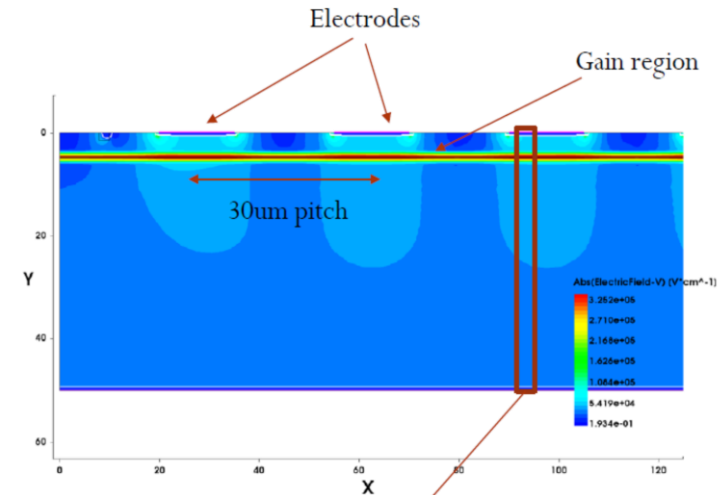
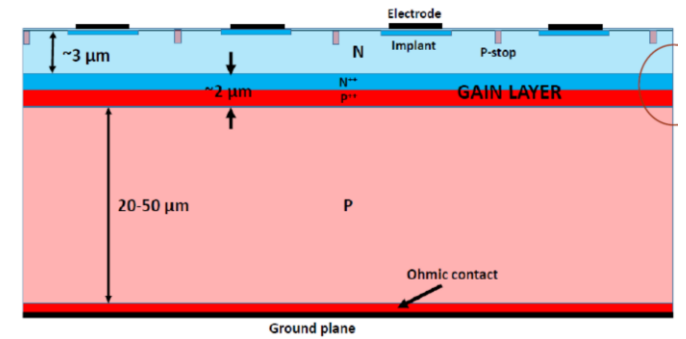
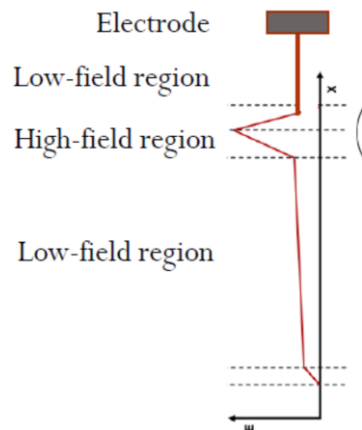
• Inter-pad area

Deep Junction LGADs

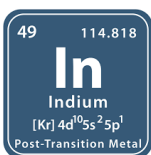
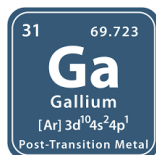
- Based on the non-segmentation of the gain region
- Move the gain layer away from the surface and deeper into the substrate
- Requires not one but two implants in a relatively deep region ($\sim 5 \mu\text{m}$) from surface
- Top implant has to be lower concentration than deeper gain implants, leading to depletion rather than impanation techniques

Status and issues

- Opposite sign signal on adjacent pad
- Only theoretical simulations from UCSC, BNL will work on a process
- Susceptible to crosstalk, but dependent probably on gain layer positioning
- Increase in process complications



• Implantation Parameters – Energy 1



Implantation energy

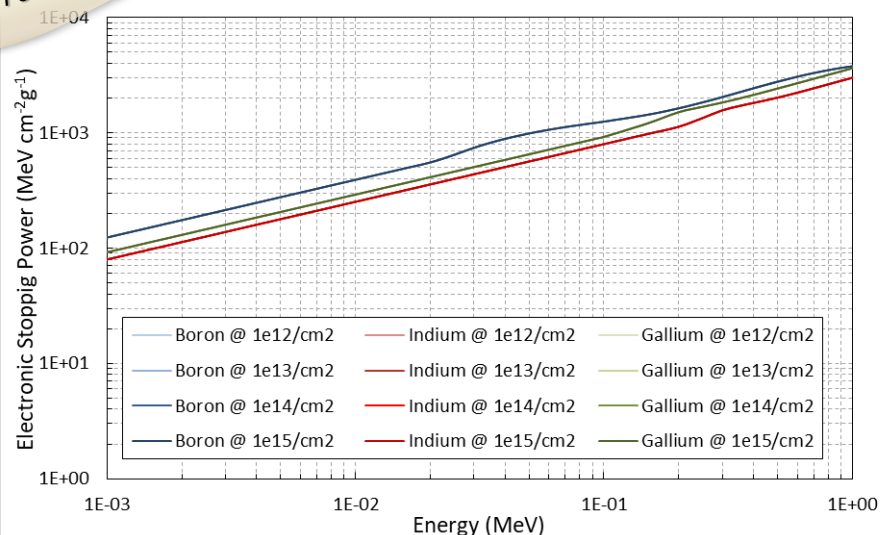
Stopping Power

High energy
electronic stopping power
(ionization)

Low energy regime
nuclear stopping power
(elastic scattering)

Electronic Stopping Power in Si

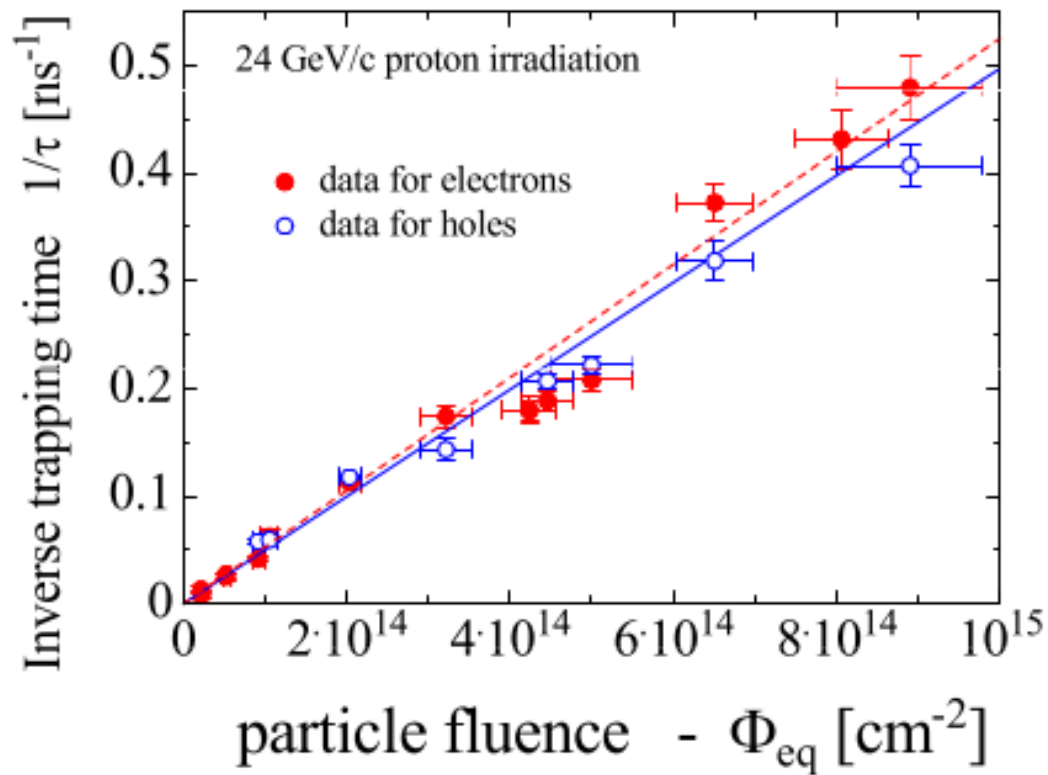
Electronic only Stopping Power in Silicon (SRIM)



Electronic stopping power (ionization) consistency

- ✓ Varies with incident particle dose
- ✓ Typical doses in range of $10^{12} \text{ cm}^{-2} - 10^{15} \text{ cm}^{-2}$
- ✓ Simulate with:
 - Real time implementation of SRIM (2013 version)
 - Implantation in pure ^{14}Si
 - No delta-ray assumption
 - Four different implantation doses tested (10^{12} , 10^{13} , 10^{14} , 10^{15})
 - Energy range 1 keV – 1MeV
- ✓ No variation for any of the three implants in the dose rate of interest

• Mobility & Trapping



• Mobility & Trapping

N_{eff} – Dynamic Model

Radiation damage modeling		
Constant Damage Terms	Acceptor Introduction	$N_{acc.}^{con.}(t) = g_{c_A} \times \int_0^t \Phi_{eq.}(\tau) \partial \tau$
	Donor Introduction	$N_{don.}^{con.}(t) = g_{c_D} \times \int_0^t \Phi_{eq.}(\tau) \partial \tau$
	Acceptor Removal	$N_{acc.}^{rem.}(t) = f_{c_A} \times N_{eff.}(0) \left(1 - e^{-c_{c_A} \int_0^t \Phi_{eq.}(\tau) \partial \tau} \right)$
	Donor Removal	$N_{don.}^{rem.}(t) = f_{c_D} \times N_{eff.}(0) \left(1 - e^{-c_{c_D} \int_0^t \Phi_{eq.}(\tau) \partial \tau} \right)$
Short term annealing	Acceptor Reduction	$N_{acc.}^{short.}(t_i) = g_A \times \int_{t_{i-1}}^{t_i} \Phi_{eq.}(\tau) \partial \tau / \delta t \times \frac{(1 - e^{-k_a(T_i) \times \delta t})}{k_a(T_i)} + N_{acc.}^{short.}(t_{i-1}) \times e^{-k_a(T_i) \times \delta t}$
Long term annealing	Max Introdactable Acceptors	$N_{acc.}^{Max. long.}(t_i) = g_Y \times \int_{t_{i-1}}^{t_i} \Phi_{eq.}(\tau) \partial \tau / \delta t \times \frac{(1 - e^{-k_Y(T_i) \times \delta t})}{k_Y(T_i)} + N_{acc.}^{Max. long.}(t_{i-1}) \times e^{-k_Y(T_i) \times \delta t}$
	Acceptor Introduction	$N_{acc.}^{long.}(t_i) = N_{acc.}^{long.}(t_{i-1}) +$ $g_Y(T) \times \frac{\int_{t_{i-1}}^{t_i} \Phi_{eq.}(\tau) \partial \tau}{k_Y(T) \delta t} \times (k_Y(T) \times t + e^{-k_Y(T)t} - 1) +$ $N_{acc.}^{Max. long.}(t_i) \times (1 - e^{-k_Y(T)t})$