

Quantum State Tomography for Attosecond Ionization of Entangled Many-Electron Systems

Dr. Marco Ruberti

*Quantum Optics and Laser Science, Blackett Laboratory,
Imperial College London*

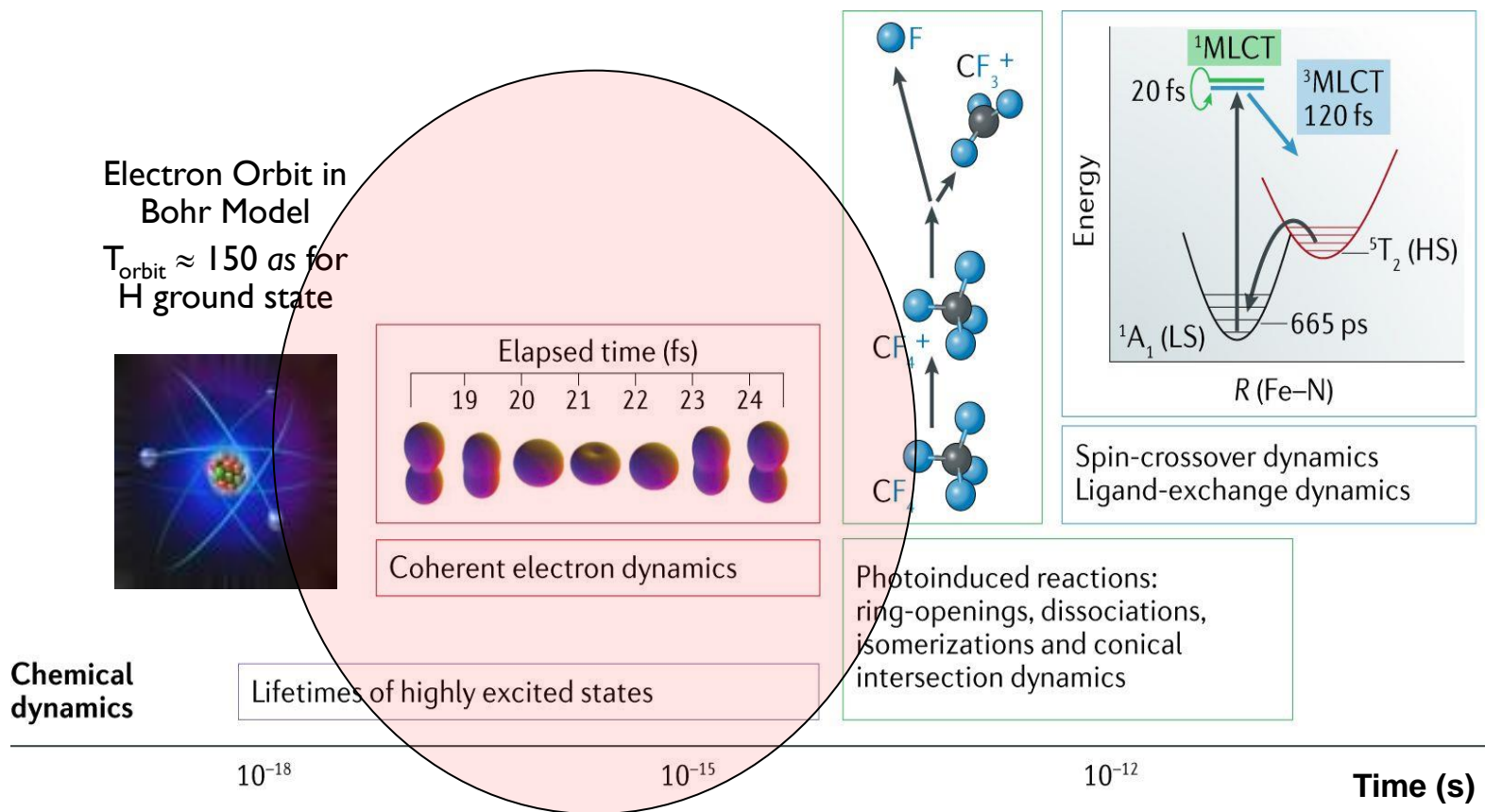
*Panofsky Fellowship Seminar,
SLAC, Stanford.*

10 March 2021

Outline

- Introduction: Quantum **coherence** and **entanglement** upon attosecond ionization of atomic and molecular systems.
- Overview of my theoretical work: the **Time dependent B-spline restricted correlation space (RCS) ADC** *ab initio* method for ionization dynamics of many-electron systems.
Applications: **mixed state** of the molecular cation prepared by strong-field and attosecond ionization; **fully *ab initio* simulation of attosecond pump–probe experiment in polyatomic molecules**; mapping of the X-ray observables to the pump-prepared quantum electronic coherences in molecular cation.
- Research vision and program for Panofsky Fellowship: **quantum-information-theory** based perspective on attosecond physics.

Attosecond many-electron quantum dynamics in matter



Many-electron motion on timescales of 0.01fs- 10 fs

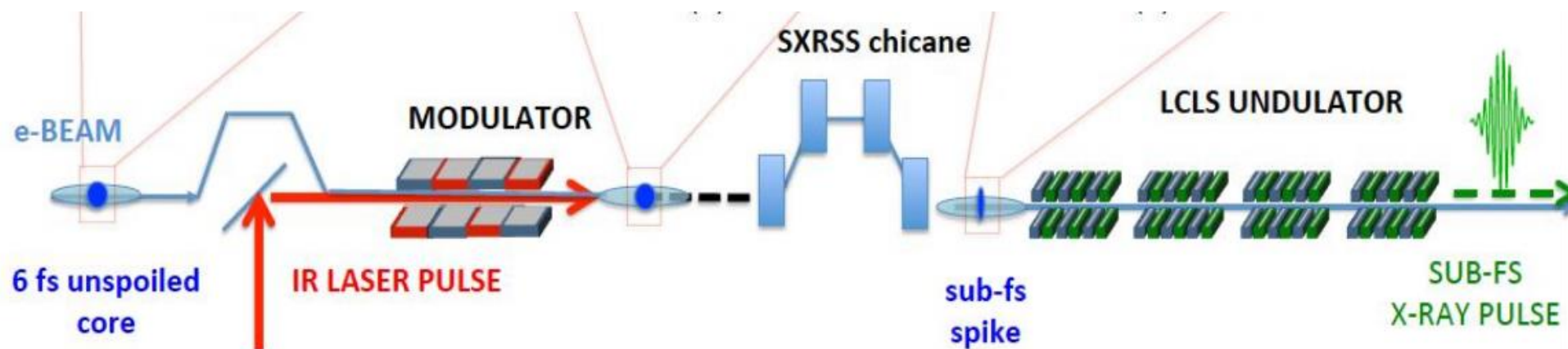
In a photophysical or photochemical process electronic and nuclear quantum dynamics will occur in a highly correlated fashion – we need tools that track this by observing the full dynamics of system i.e. both photoelectron and residual ion/neutral sub-states

Scientific Importance of Understanding Quantum Coherence and Electron-Nuclear Coupling in Photo-Excitation/Ionisation

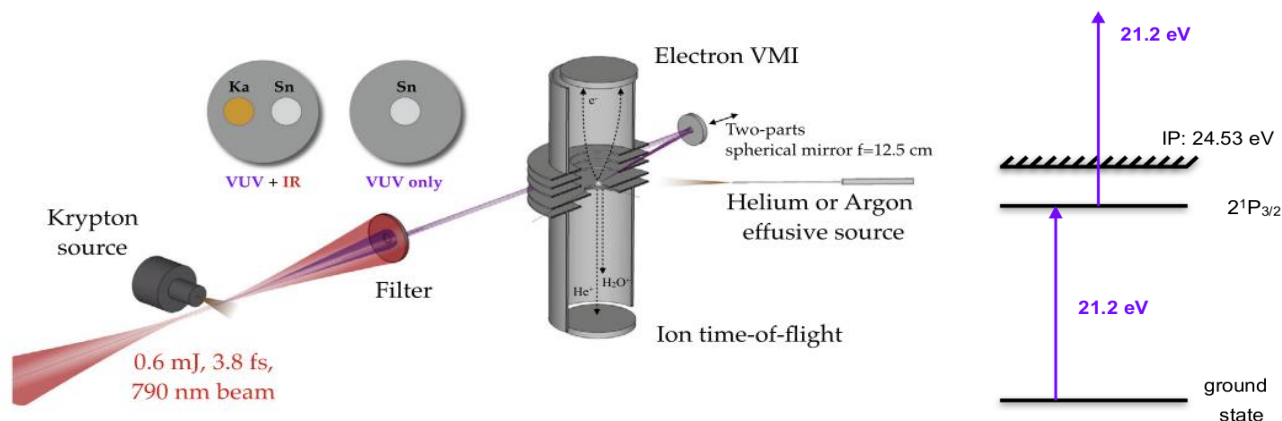
- Photo-excitation/ionisation is primary event in many processes from radiation damage to solar energy devices to photo-catalysed reactions
- Routes to control photo-physical/-chemical processes: *Imagine we can harness these fundamental, inherently quantum-mechanical concepts to develop control strategies of chemical reactivity of excited and ionized molecules by acting on the electronic degrees of freedom on attosecond time scale.*
- An experimental test-bed for quantum information in open multi-partite quantum systems: *entanglement* and *decoherence*
- Test our theories of *charge migration*, *electron-nuclear coupling* and *charge transfer* in quantum chemistry

X-ray FEL source development opens the door for attosecond pump – attosecond probe spectroscopy

From synchrotron radiation towards **attosecond X-ray pulses at LCLS-II**



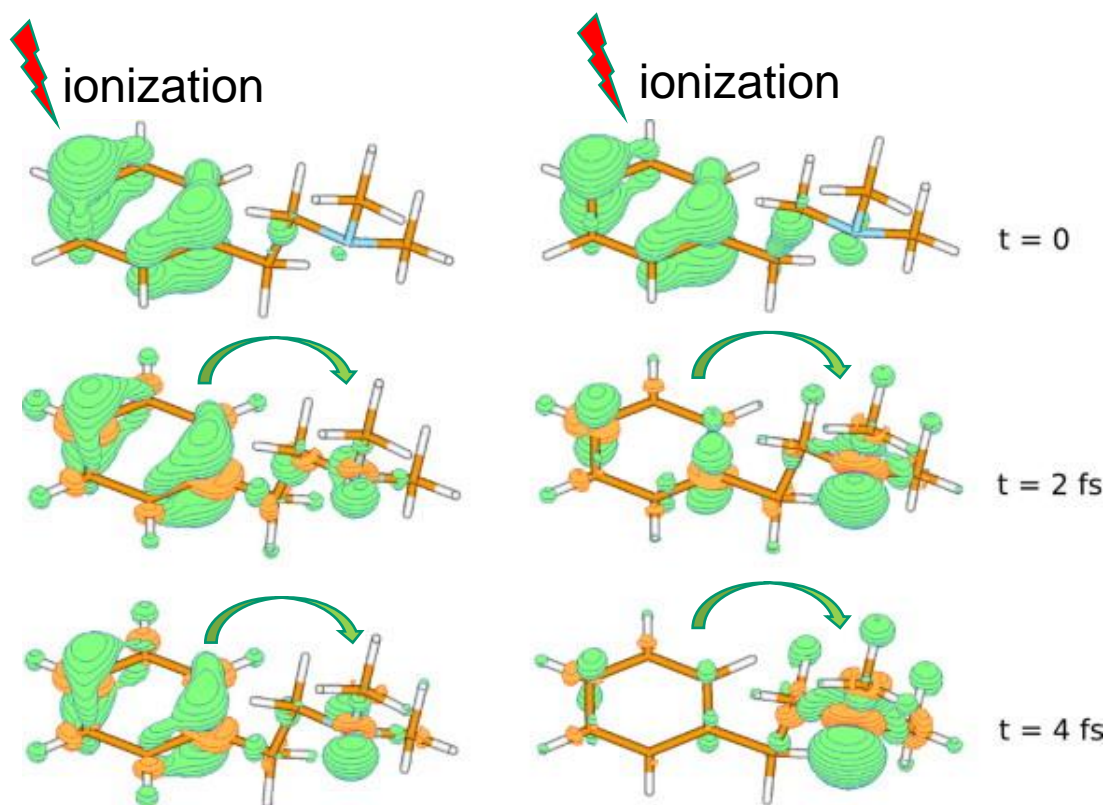
X-LEAP: 0.5 fs duration pulses, [Pellegrini 2017, Duris 2020]



HHG-based sources: [Tisch and co., Chem. Phys. Lett. **683** (2017) 38]

Hole migration: ultrafast bound-state charge dynamics following molecular ionization

- ❖ Valence photo-ionization can create a non-stationary state of molecular ion
- ❖ This leads to oscillations of the electron hole across the molecule that we call **hole migration**
- ❖ Purely electronic process. Eventually damped by nuclear motion



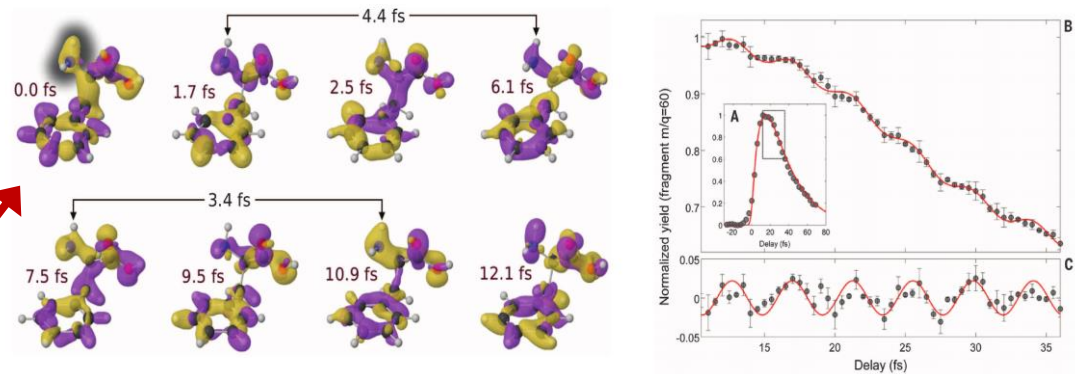
[Lünnemann, Kuleff & Cederbaum,
Chem. Phys. Lett. **450**, 232 (2008)]

Hole migration: ultrafast bound-state charge dynamics following molecular ionization

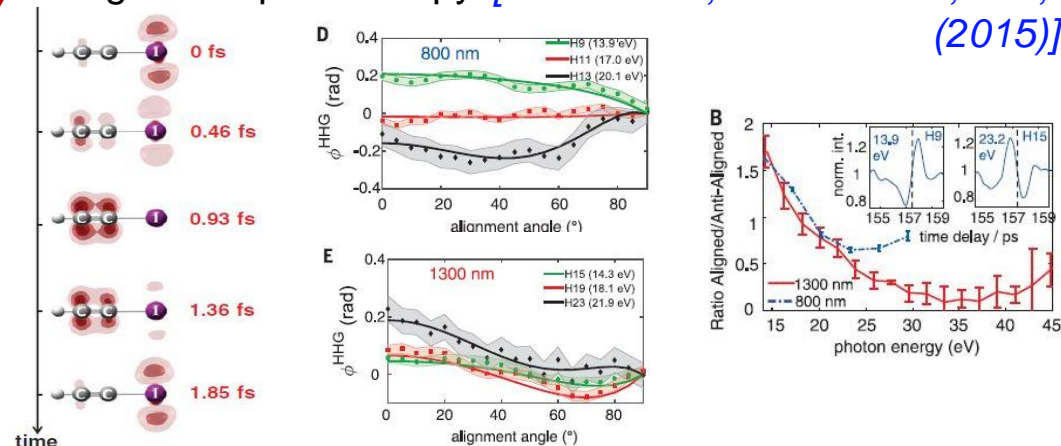
- ❖ Valence photo-ionization can create a non-stationary state of molecular ion
- ❖ This leads to oscillations of the electron hole across the molecule that we call **hole migration**
- ❖ Purely electronic process. Eventually damped by nuclear motion

First attempts of time-resolved observation

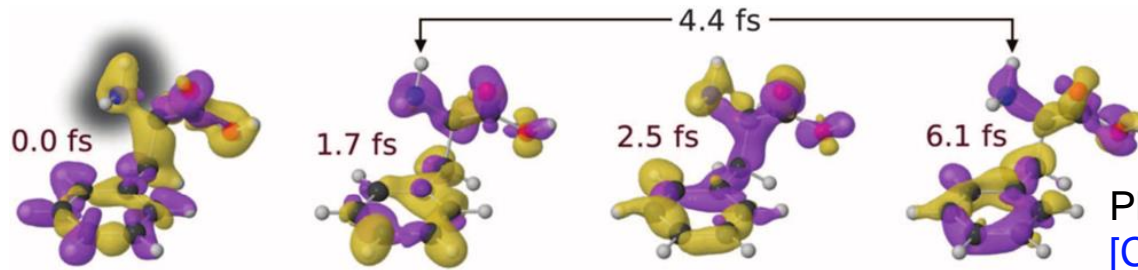
- 1) using attosecond pump – IR probe spectroscopy: [Belshaw et al., *J. Phys. Chem. Lett.* **3**, 3751 (2012); Calegari et al., *Science* **346**, 336 (2014)]



- 2) using HHG spectroscopy: [Kraus et al., *Science* **350**, 790, (2015)]

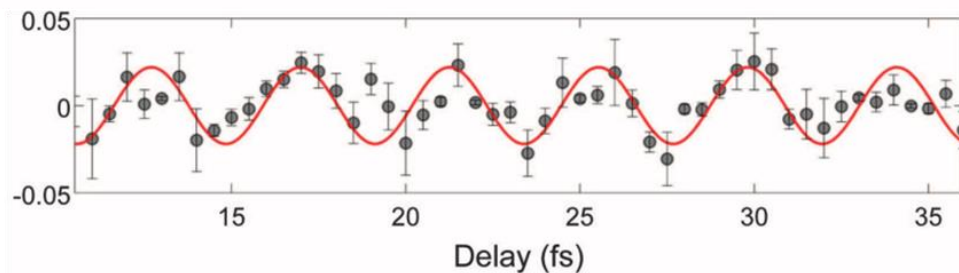


Hole migration: ultrafast bound-state charge dynamics following molecular ionization



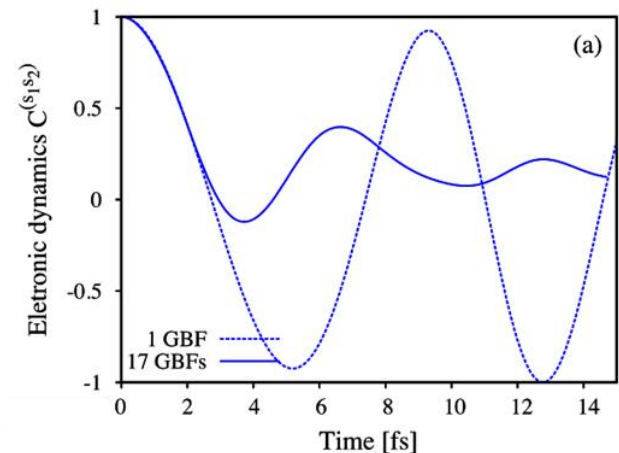
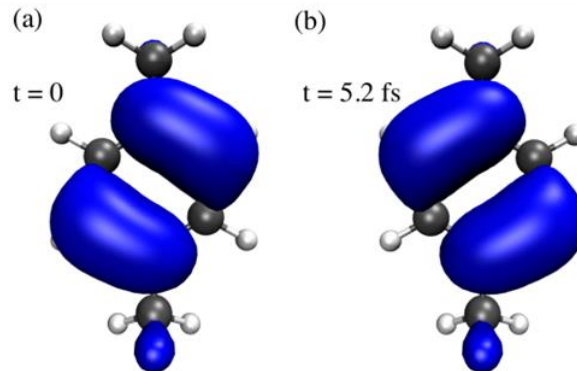
Phenylalanine
[Calegari *et al.*,
Science **346**, 336 (2014)]

**Example of
ultrafast
consequences
of coupling of
electronic and
nuclear states**



★ Paraxylene dipeptide -
Theory suggests
nuclear motion damps
the oscillations much
faster than the 2014
experiment implies...

[Vacher *et al.*,
PRL **118**, 083001 (2017)]



Hole migration: problem solved?

- ★ Hole migration is a simple effect of the coherent population of ionic eigenstates...
- ★ But the coherence of ionic state population is a highly non-trivial consequence of several photoionization mechanisms and is poorly understood.
- ★ We need a theoretical tool to model molecular photoionization and predict to what extent can the resulting ion be described by a wavefunction!

Needed:

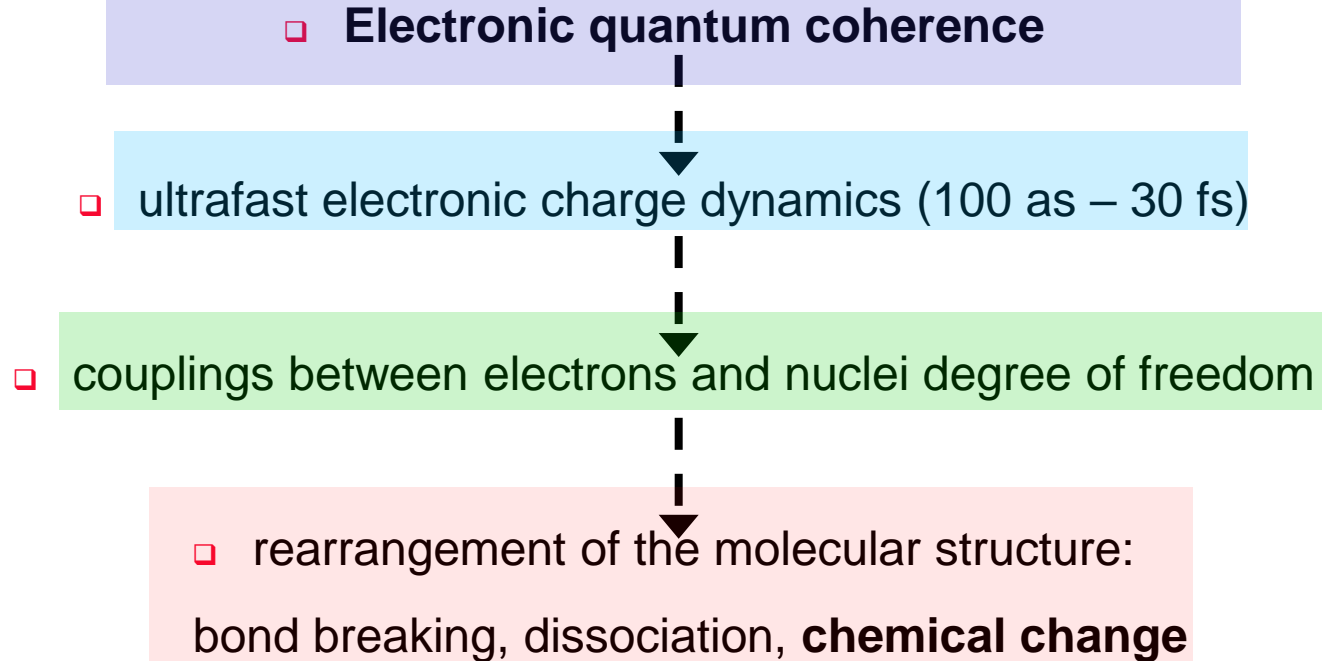
ab initio framework for

1- designing the
upcoming attosecond pump-probe studies through
reliable simulations

2- interpretation of the
ongoing mixed attosecond/femtosecond pump-probe studies
through reliable simulations

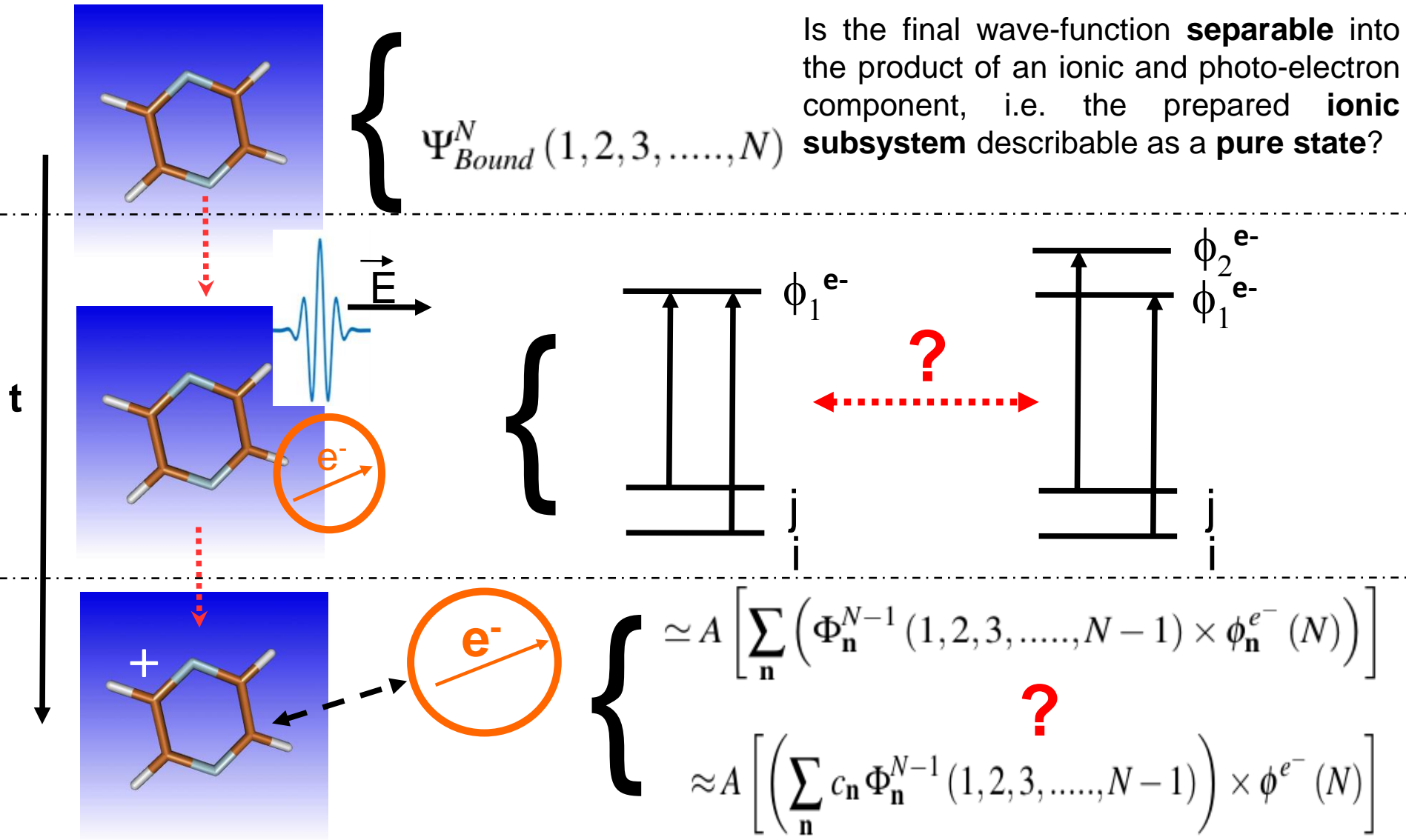
**A strong coupling between theoretical,
computational and experimental efforts is key
for the success of attosecond science**

Goals (...Attochemistry...?)

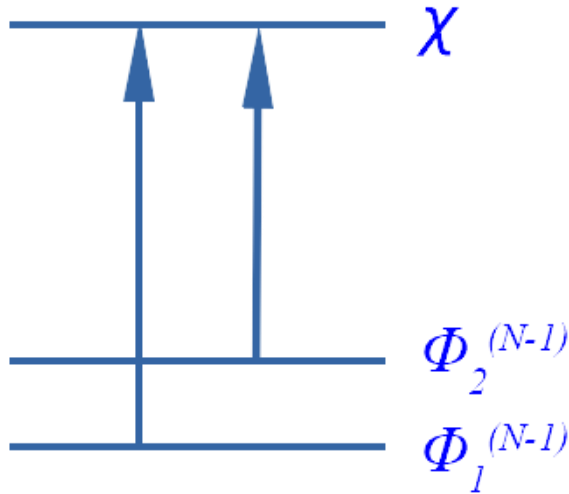


- 1) Preparation of **quantum electronic coherences**
- 2) Probing of **quantum electronic coherences**
- 3) Control of (and by means of) **quantum electronic coherences**

Ionic coherence after photoionization?

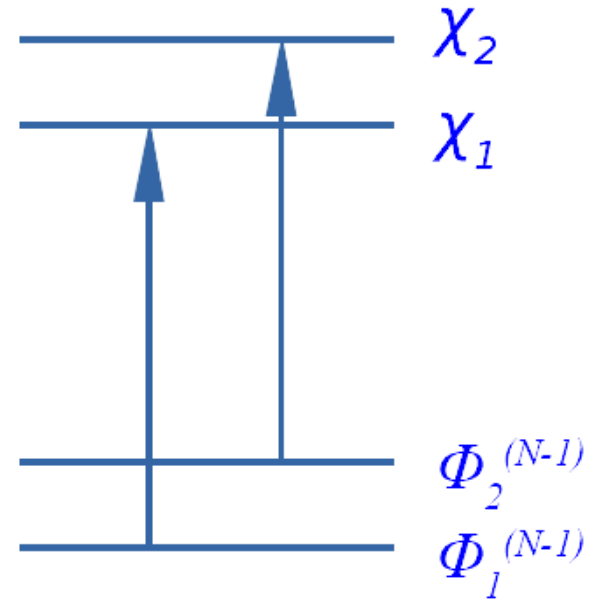


Ionic coherence after ionization?



$$\Psi = [C_1 \Phi_1^{(N-1)} + C_2 \Phi_2^{(N-1)}] \chi$$

(N-1)-electron ion has got a [...] wavefunction – coherent population of ionic states



$$\Psi = C_1 \Phi_1^{(N-1)} \chi_1 + C_2 \Phi_2^{(N-1)} \chi_2$$

Only the full N-electron system has got a wavefunction – incoherent population of ionic states.

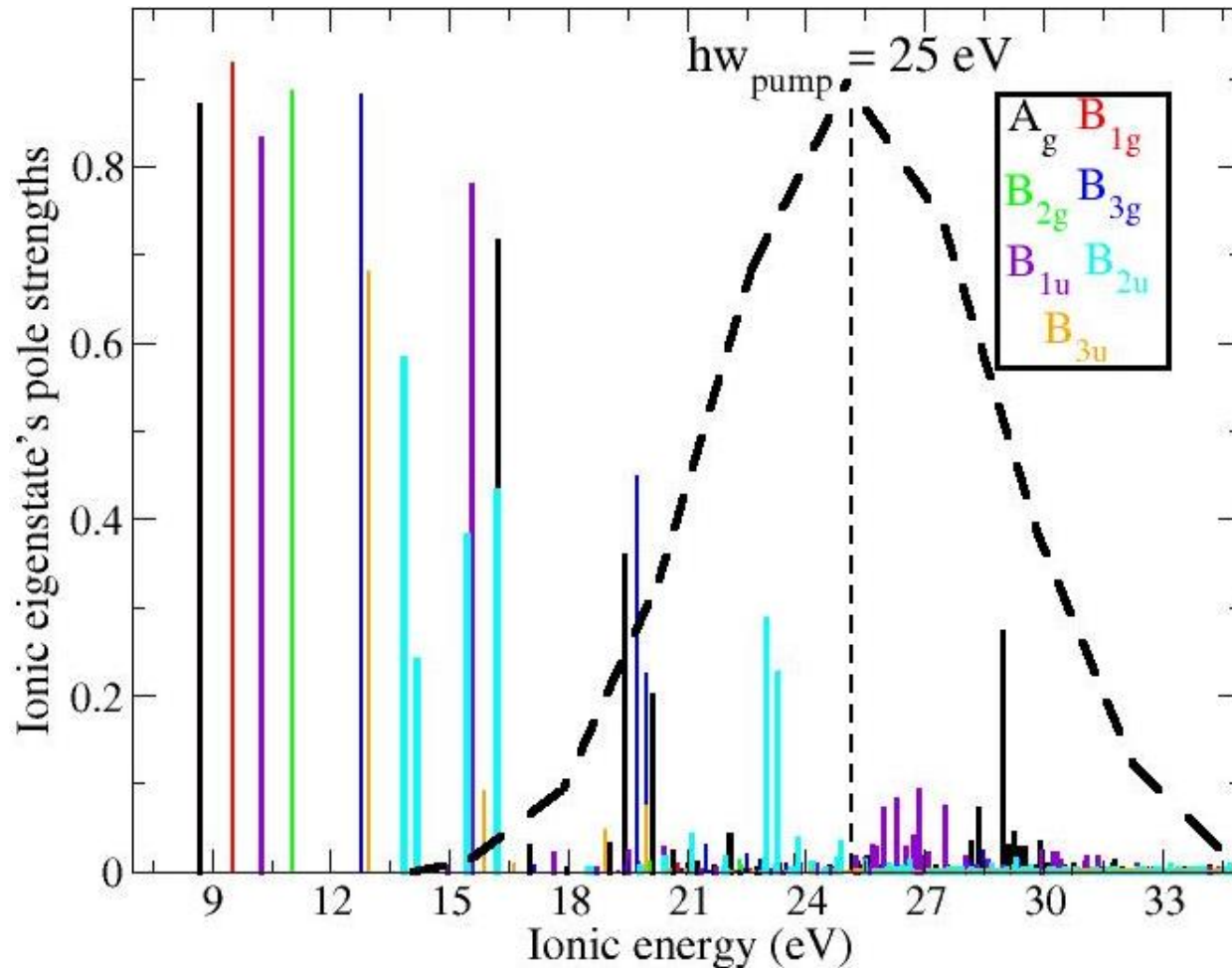
Ion-photoelectron entanglement!

The complete information is contained in the full N-electron wavefunction.

Regimes of Coherent Hole-State Wavepacket Formation: Dependence on Photon Energy & Bandwidth

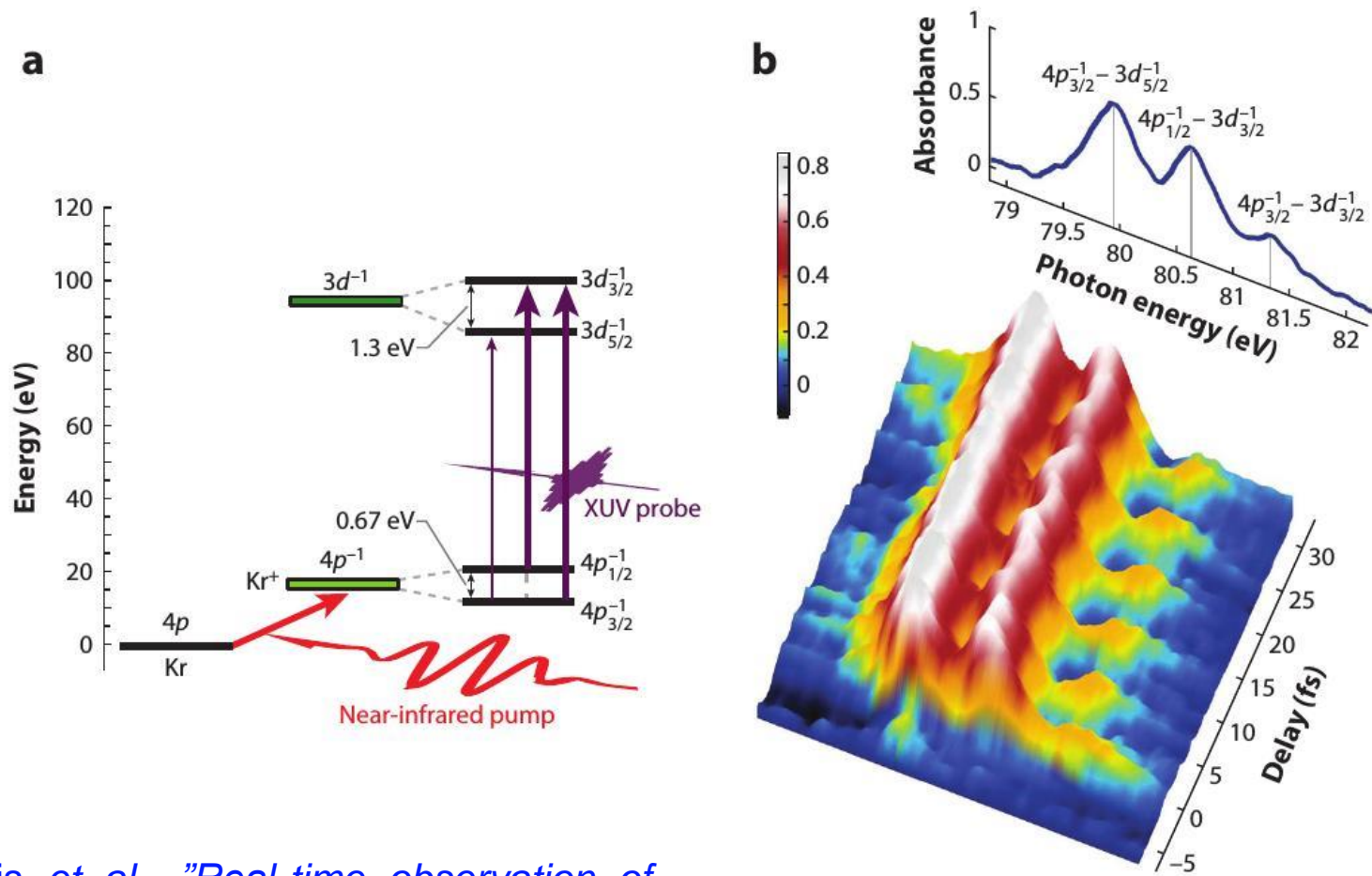
Central Pump photon energy = 25 eV

Bandwidth FWHM = 7 eV



Measuring ionic coherence: the atomic Kr spin-orbit two-level case

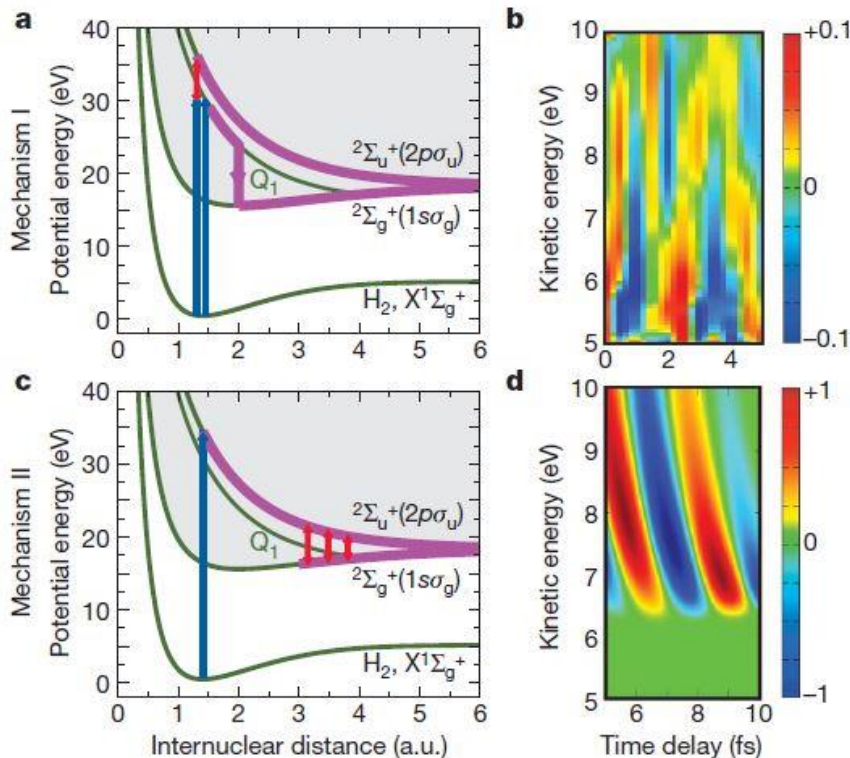
Coherence
between
different
spin-orbit
states
upon
strong-field
IR ionization
of Kr atom



[E. Goulielmakis *et al.*, "Real-time observation of valence electron motion", Nature **466**, 739 (2010)]

Effects of coherence: electron localization in XUV-IR dissociative photoionization of H₂

"H₂ and D₂ are dissociatively ionized by a sequence comprising an isolated attosecond ultraviolet pulse and an intense few-cycle infrared pulse, and a localization of the electronic charge distribution within the molecule is measured that depends—with attosecond time resolution—on the delay between the pump and probe pulses."



Asymmetry in EUV-IR
dissociative ionization of H₂

Electron localization in H₂⁺
**Coherent superposition of
gerade and *ungerade* ionic states**

Electron localization in the ion
is influenced by the interaction
between the IR field
and the photoelectron

Summary

- ❑ Quantum **coherence** and **entanglement** are key concepts in attosecond science. For example, these competing phenomena underpin/destroy electronic hole migration in molecular ions.
- ❑ We are only at the very beginning of appreciating *the role of ultrafast quantum electronic coherences and entanglement in the photoinduced physical & chemical change of matter.*
- ❑ A full quantum-information perspective on attoscience is thus far unexplored and the role of electron-nuclear coupling in this game is even more obscure – *how would an interaction with or a measurement performed on the photoelectron affect the chemical process (dissociation, rearrangement,) in the ion that this electron left behind?*
- ❑ There is a room for a new field to be born: **quantum attoscience.**

Ionic wavefunction...?

The “Sudden Approximation”

Assuming perfect coherence

$$\Psi_{Final}^N(1, 2, 3, \dots, N) \sim A \left[\Psi_{Bound}^{N-1}(1, 2, 3, \dots, N-1) \times \phi^{e^-}(N) \right]$$

$$\Psi_{Bound}^{N-1}(1, 2, 3, \dots, N-1) = \sum_n c_n \Phi_n^{N-1}(1, 2, 3, \dots, N-1)$$

$$c_n = c_n(E_{field}) = ?$$

? Dependence on
laser pulse parameters ?

Sudden approximation
ansatz for the initial state

We need a theoretical tool to model molecular photo-ionization beyond sudden approximation, including electron correlation

$$|\Psi_{initial}^{N-1}\rangle \sim a_i |\Psi_0^N\rangle$$

when is it valid...?

(Ionic wavefunction...?)

Beyond The “Sudden Approximation”: TD B-spline RCS-ADC theory

New *ab initio* method:

Time-dependent (TD) multicentre- **B-spline**
restricted correlation space (**RCS**) –
algebraic diagrammatic construction **ADC**

$$i\hbar \frac{\partial |\Psi^N(t)\rangle}{\partial t} = \hat{H}^N(t) |\Psi^N(t)\rangle$$

Time-dependent
many-electron
Schrodinger equation

$$|\Psi^N(t)\rangle = \sum_m \left\{ \sum_{\mu} c_{m\mu}(t) \hat{c}_{\mu}^{\dagger} |\Psi_m^{(N-1)}\rangle^{[n]} \right\} + \sum_{I_{RCS}} c_{I_{RCS}}(t) |\tilde{\Psi}_{I_{RCS}}^N\rangle^{[n]} + c_0(t) |\Psi_0^{RCS}\rangle^{[n]}$$

continuum/delocalised
bound/localised

Algebraic Diagrammatic Construction (ADC) in the Intermediate State Representation (ISR)

CI-like
expansion

$$\Psi = \sum_{1h1p} Y_i^a \begin{array}{c} \vdots \\ \text{---} c \\ \text{---} b \\ \uparrow a \end{array} \begin{array}{c} \uparrow\uparrow \uparrow\uparrow \uparrow\uparrow \\ \circ \uparrow \\ \uparrow\uparrow \end{array} + \sum_{2h2p} Y_{ij}^{ab} \begin{array}{c} \vdots \\ \text{---} c \\ \uparrow\uparrow b \\ \uparrow\uparrow a \end{array} \begin{array}{c} \uparrow\uparrow \uparrow\uparrow \uparrow\uparrow \\ \circ \uparrow \\ \uparrow\uparrow \\ \uparrow\uparrow \end{array} + \dots$$

$$\Psi_i^{a(0)} = c_a^\dagger c_i \Psi_0$$

$$\Psi_{ij}^{ab(0)} = c_a^\dagger c_b^\dagger c_i c_j \Psi_0$$

$$\Psi_0 = \Phi_0^{\text{HF}} + \Psi_0^{(1)} + \Psi_0^{(2)} + \dots$$

many-body PT for
the ground state

Arnoldi-Lanczos time-propagation scheme

Traditionally, standard *Gaussian* single-electron basis sets are employed, but here I use *B-splines*

Analogous formulation for
singly-ionised states is available

[Cederbaum, Schirmer and co-workers;
Using B-splines: Ruberti, Decleva & Averbukh (2014)]

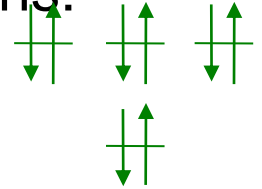
Simple theory of electron holes: Koopmans theorem



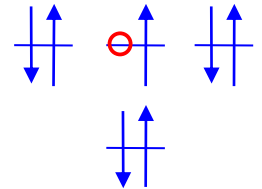
Tjalling C. Koopmans (Nobel prize winner in *economics*, 1975)

If both the ground state of the neutral and the eigenstate of the cation are well approximated by Single HF configurations:

$$\Psi_0^{(N)} = \Phi_0^{\text{HF}}$$



$$\Psi_0^{(N-1)} = \hat{a}_i \Phi_0^{\text{HF}}$$



then the corresponding ionization potential is given by the HF orbital energy:

$$\text{IP} = \mathbf{E}_i(\mathbf{N}-1) - \mathbf{E}_0(\mathbf{N}) = -\boldsymbol{\varepsilon}_i$$

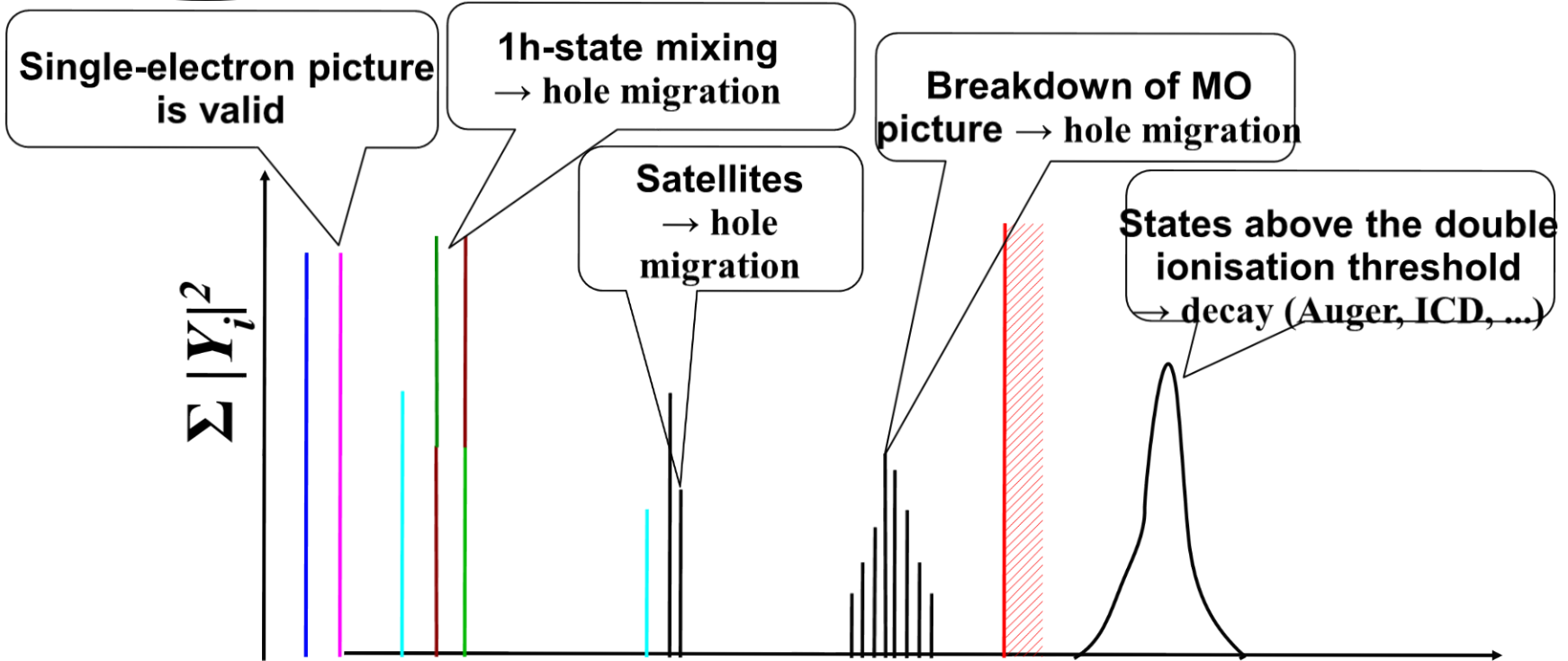
(Ionic wavefunction...?)

ADC ionic states:

Electron holes beyond Koopmans

$$|\tilde{\Psi}_{\mu,n}^N\rangle = \hat{c}_{\mu}^{\dagger} |\Psi_n^{N-1}\rangle$$

$$\Psi^{(N-1)} = \sum_{1h} Y_i \Psi_i + \sum_{2h1p} Y_{ij}^a \Psi_{ij}^a + \sum_{3h2p} Y_{ijk}^{ab} \Psi_{ijk}^{ab} + \dots$$

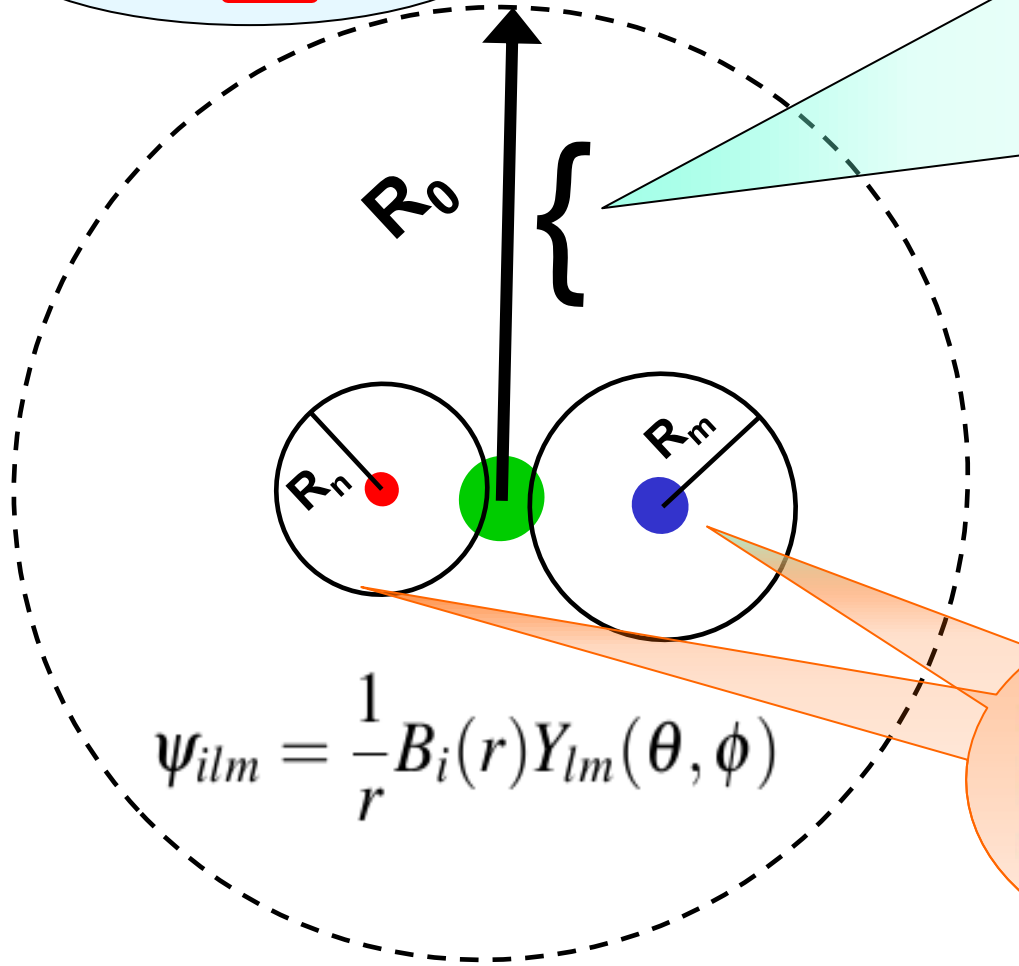


[Cederbaum, Schirmer and co-workers;
Using B-splines: Ruberti, Decleva & Averbukh (2014)]

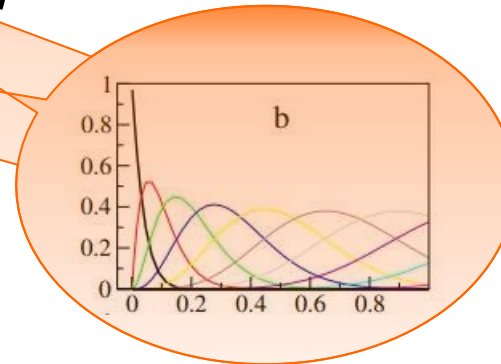
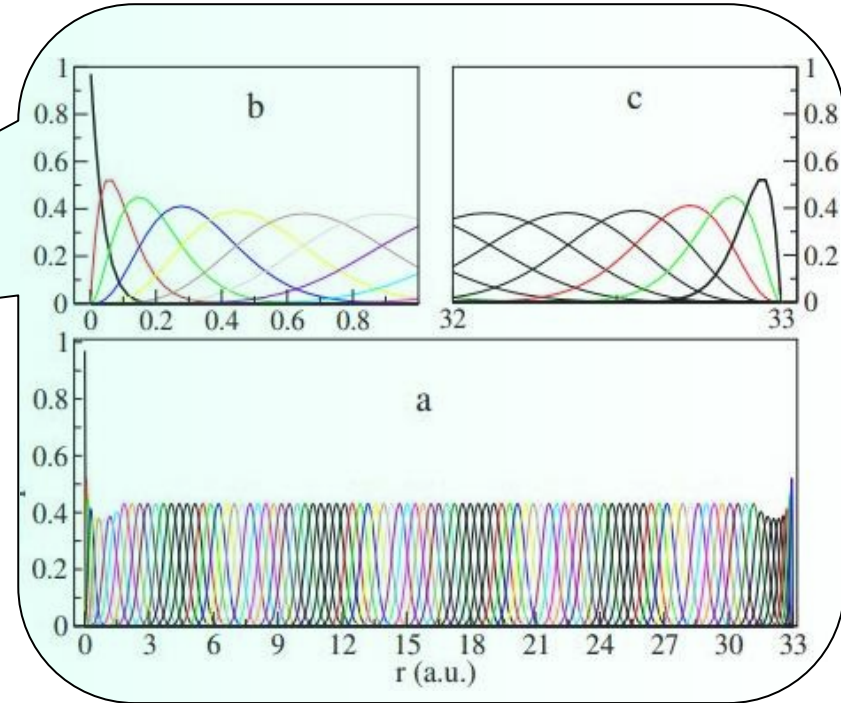
$E_0(N-2)$ $E(N-1)$

Single-electron basis set: Multicentre B-splines

$$|\tilde{\Psi}_{\mu,n}^N\rangle = \hat{C}_{\mu}^{\dagger} |\Psi_n^{N-1}\rangle$$



$$\psi_{ilm} = \frac{1}{r} B_i(r) Y_{lm}(\theta, \phi)$$



[D. Toffoli & P. Decleva
Chem. Phys. (2002)]

TD B-spline RCS-ADC theory

$$H^{\hat{N}}(t) = \hat{H}_{RCS-ADC}^N + \hat{D}_{RCS-ADC}^N E(t) - i\hat{W}$$

Dipole approximation
Length Gauge

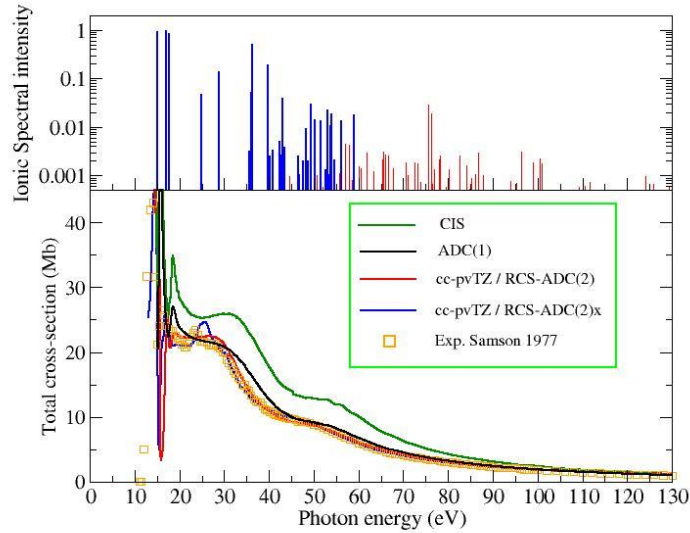
Complex Absorbing Potential
(CAP)

$$\mathbf{H}^{\text{RCS-ADC}(n)} = \begin{pmatrix} \mathbf{H}_{\text{RCS}}^{\text{ADC}(n)} & \langle \tilde{\Psi}_{I_{\text{RCS}}}^N | \hat{H} \hat{c}_\nu^\dagger | \Psi_m^{N-1} \rangle_{\text{ADC}(n)} \\ \langle \Psi_m^{N-1} | \hat{c}_\mu \hat{H} | \tilde{\Psi}_{I_{\text{RCS}}}^N \rangle_{\text{ADC}(n)} & \langle \Psi_m^{N-1} | \hat{c}_\mu \hat{H} \hat{c}_\nu^\dagger | \Psi_n^{N-1} \rangle_{\text{ADC}(n)} \end{pmatrix}$$

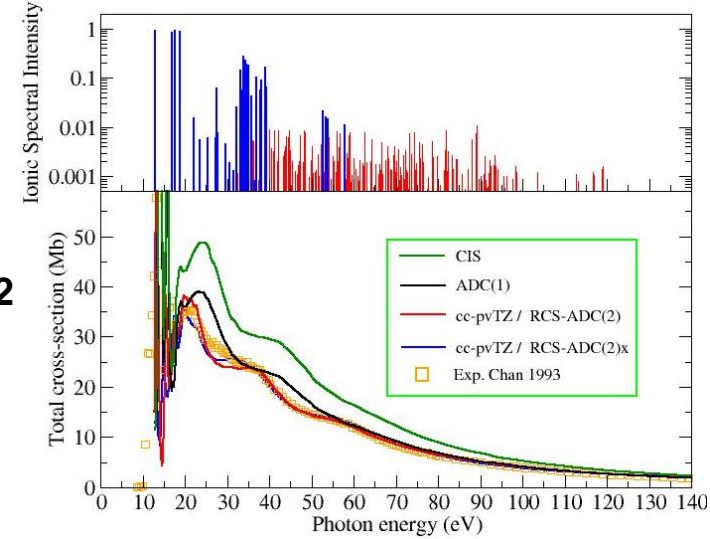
$$\mathbf{D}^{\text{RCS-ADC}(n)} = \begin{pmatrix} \mathbf{D}_{\text{RCS}}^{\text{ADC}(n)} & \langle \tilde{\Psi}_{I_{\text{RCS}}}^N | \hat{D} \hat{c}_\nu^\dagger | \Psi_m^{N-1} \rangle_{\text{ADC}(n)} \\ \langle \Psi_m^{N-1} | \hat{c}_\mu \hat{D} | \tilde{\Psi}_{I_{\text{RCS}}}^N \rangle_{\text{ADC}(n)} & \langle \Psi_m^{N-1} | \hat{c}_\mu \hat{D} \hat{c}_\nu^\dagger | \Psi_n^{N-1} \rangle_{\text{ADC}(n)} \end{pmatrix}$$

$$\langle \Psi_0^{\text{RCS}} | \hat{D} \hat{c}_\mu^\dagger | \Psi_n^{N-1} \rangle = \sum_{\gamma \in \text{RCS}} d_{\gamma\mu} \langle \Psi_0^{\text{RCS}} | \hat{c}_\gamma^\dagger | \Psi_n^{N-1} \rangle = \sum_{\gamma \in \text{RCS}} d_{\gamma\mu} x_{\gamma n}$$

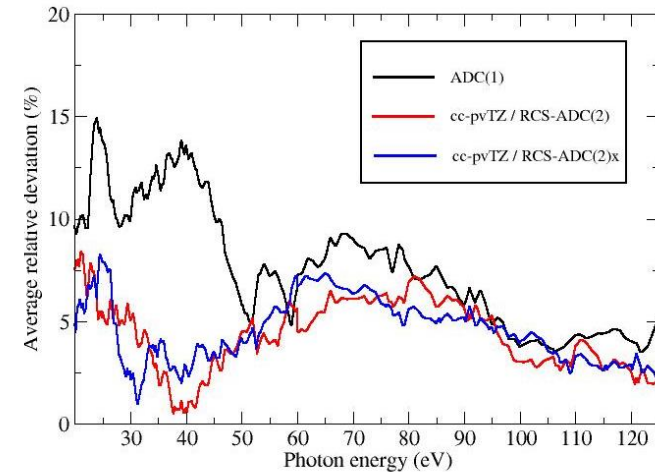
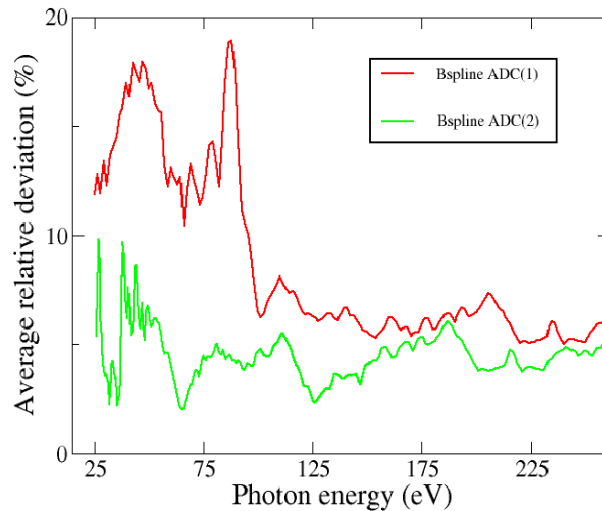
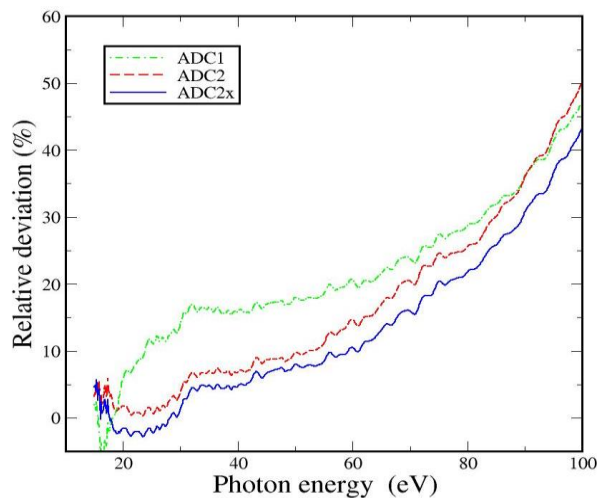
Benchmark calculations: atomic and molecular total photoionization cross-sections



N_2



CO_2



[Ruberti *et al.*,
JCP **139**, 144107 (2013)]

[Ruberti, Decleva and Averbukh,
JCP **141**, 164126 (2014)]

[Ruberti, JCTC **15**, 6, 3635 (2019)]

(Ionic wavefunction...no!) Reduced Ionic Density Matrix

Reduced ionic density matrix (R-IDM):
trace over the **unobserved** photo-electron states

$$\hat{\rho}(t) = |\Psi^N(t)\rangle\langle\Psi^N(t)| \quad \hat{\rho}^{R-IDM}(t) = Tr_{\mu}[\hat{\rho}(t)]$$

$$\hat{\rho}^{R-IDM}(t) = \sum_{m=1,N} \sum_{n=1,N} \rho_{mn}^{R-IDM}(t) |\Psi_m^{N-1}\rangle\langle\Psi_n^{N-1}|$$

Ionic density matrix
from TD B-spline RCS-ADC theory

$$\rho_{mn}^{R-IDM}(t) = \sum_{\mu} \langle\tilde{\Psi}_{\mu,m}^N | \Psi^N(t)\rangle\langle\Psi^N(t) | \tilde{\Psi}_{\mu,n}^N\rangle = \sum_{\mu} c_{m\mu}(t) c_{n\mu}^*(t)$$

[M. Ruberti, P. Decleva and V. Averbukh, JCTC **14**, 4991-5000 (2018)]

[M. Ruberti, Phys. Chem. Chem. Phys., 2019, **21**, 17584-17604 PCCP HOT Article Collection 2019]

Reduced Ionic Density Matrix

Populations

of ionic state n

$$P_n(t) = |\rho_{nn}(t)|$$

Degrees of coherence

between pair of populated ionic states m and n

$$G_{mn}(t) = \frac{|\rho_{mn}(t)|}{\sqrt{P_m(t) * P_n(t)}}$$

Maximum quantum coherence between two states corresponds to $G_{mn} = 1$

Purity

$$Tr(\hat{\rho}^2(t)) \neq Tr(\hat{\rho}(t)) \leq 1 \text{ mixed state,}$$

$$Tr(\hat{\rho}^2(t)) = Tr(\hat{\rho}(t)) = 1 \text{ pure state.}$$

von Neumann

entropy of entanglement

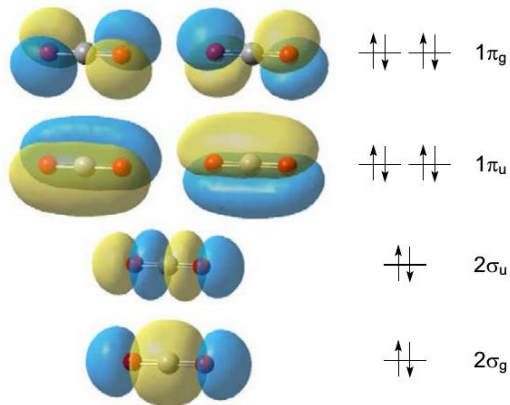
$$s^{Ion}(t) = -\ln\left(\frac{1}{N}\right) \text{ statistical mixture,}$$

$$s^{Ion}(t) = 0 \text{ pure state.}$$

[M. Ruberti, P. Decleva and V. Averbukh, JCTC **14**, 4991-5000 (2018)]

[M. Ruberti, Phys. Chem. Chem. Phys., 2019, **21**, 17584-17604 PCCP HOT Article Collection 2019]

Ionic coherence upon IR strong-field ionization of aligned CO₂ molecules:



State	Experimental I.P.s (eV)	HF B-splines basis set I.P.s (eV)
$ \pi_g\rangle^{-1}$	13.8	15.03
$ \pi_u\rangle^{-1}$	17.3	18.81
$ \sigma_u\rangle^{-1}$	18.1	19.6
$ \sigma_g\rangle^{-1}$	19.36	20.7

Carrier-frequency dependence of the **degree of coherence** and of the **ionic wave-packet parameters**

Ultra-short 2 cycles pulse.

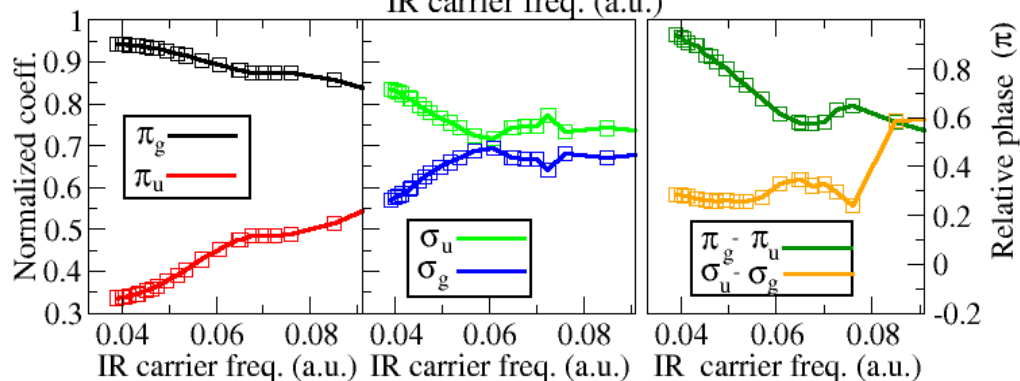
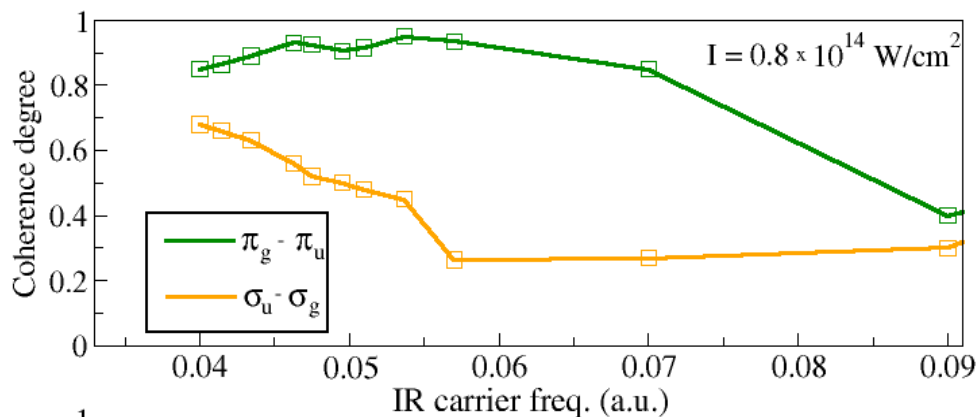
Linear polarization along the molecular axis.

Peak intensity: 0.8×10^{14} W/cm²

Σ states Energy Gap = **1.26 eV**

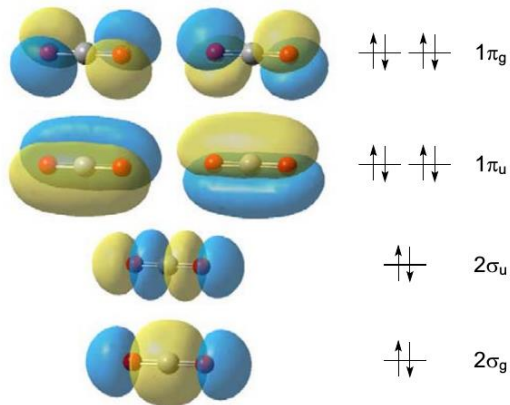
Π states Energy Gap = **3.518 eV**

$$\Psi^\Sigma(t) = C_{\Sigma_u} |\sigma_u^{-1}\rangle + C_{\Sigma_g} e^{-i(E_{\Sigma_g} - E_{\Sigma_u})t} e^{-i\phi_{\Sigma_u - \Sigma_g}} |\sigma_g^{-1}\rangle$$



[M. Ruberti, P. Decleva and V. Averbukh, JCTC 14, 4991-5000 (2018)]

Ionic coherence upon IR strong-field ionization of aligned CO₂ molecules:



State	Experimental I.P.s (eV)	HF B-splines basis set I.P.s (eV)
$ \pi_g\rangle^{-1}$	13.8	15.03
$ \pi_u\rangle^{-1}$	17.3	18.81
$ \sigma_u\rangle^{-1}$	18.1	19.6
$ \sigma_g\rangle^{-1}$	19.36	20.7

IR intensity dependence of the degree of coherence and of the ionic wave-packet parameters

Ultra-short 2 cycles pulse.

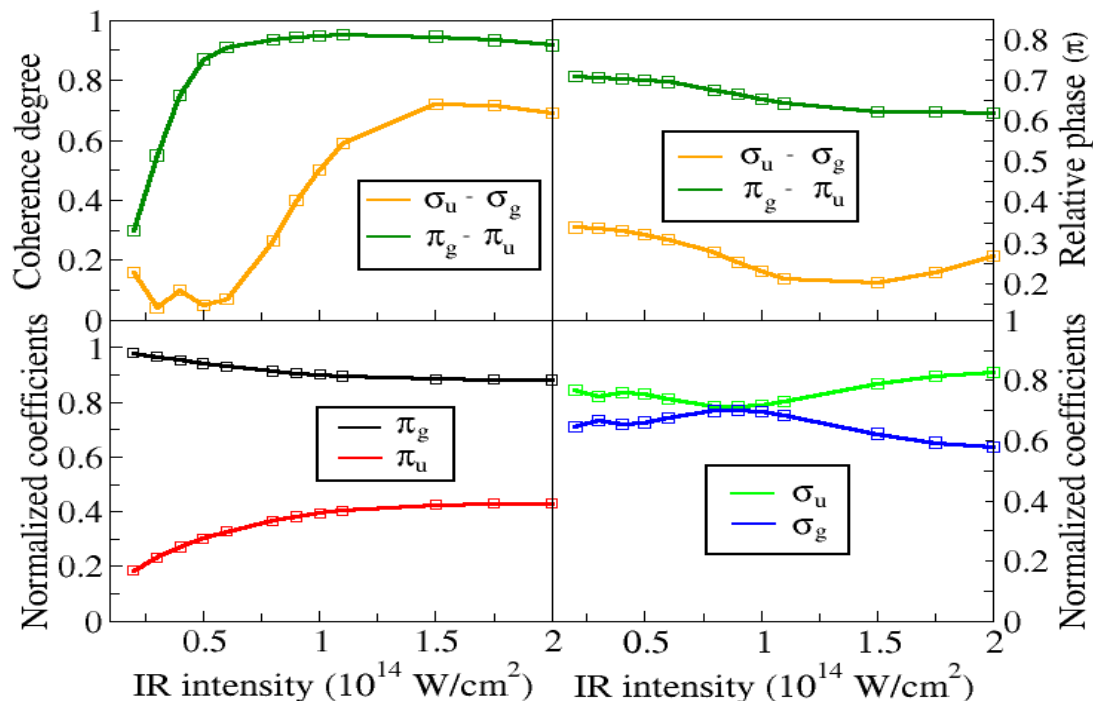
Linear polarization along the molecular axis.

IR wavelength: 800 nm

Σ states Energy Gap = **1.26 eV**

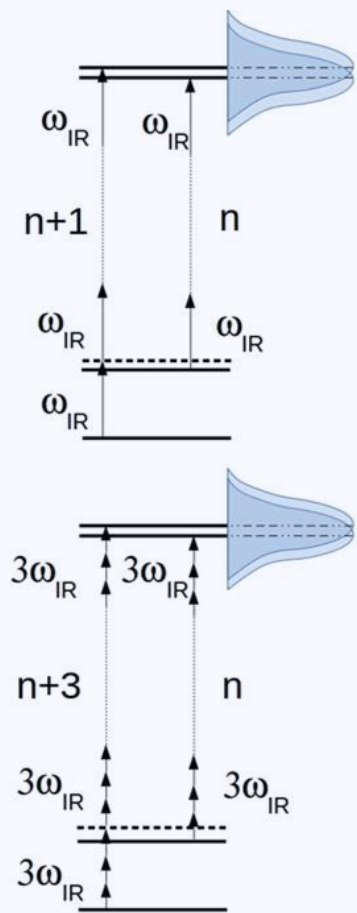
Π states Energy Gap = **3.518 eV**

$$\Psi^\Sigma(t) = C_{\Sigma_u} |\sigma_u^{-1}\rangle + C_{\Sigma_g} e^{-i(E_{\Sigma_g} - E_{\Sigma_u})t} e^{-i\phi_{\Sigma_u - \Sigma_g}} |\sigma_g^{-1}\rangle$$

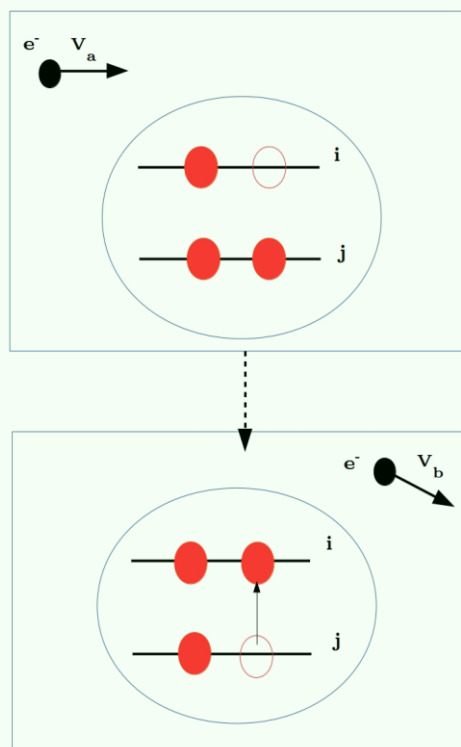


Mechanisms of coherence formation in molecular IR strong-field ionization

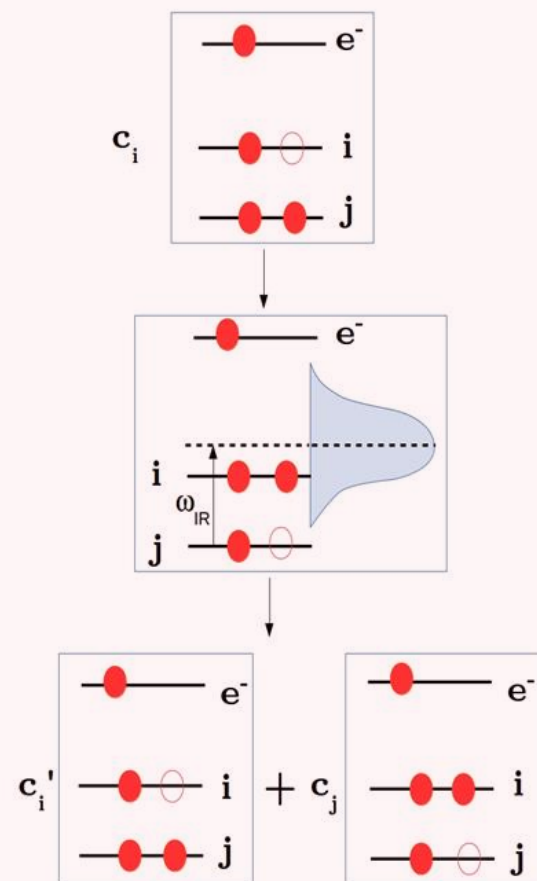
Direct ionisation into the same continuum state



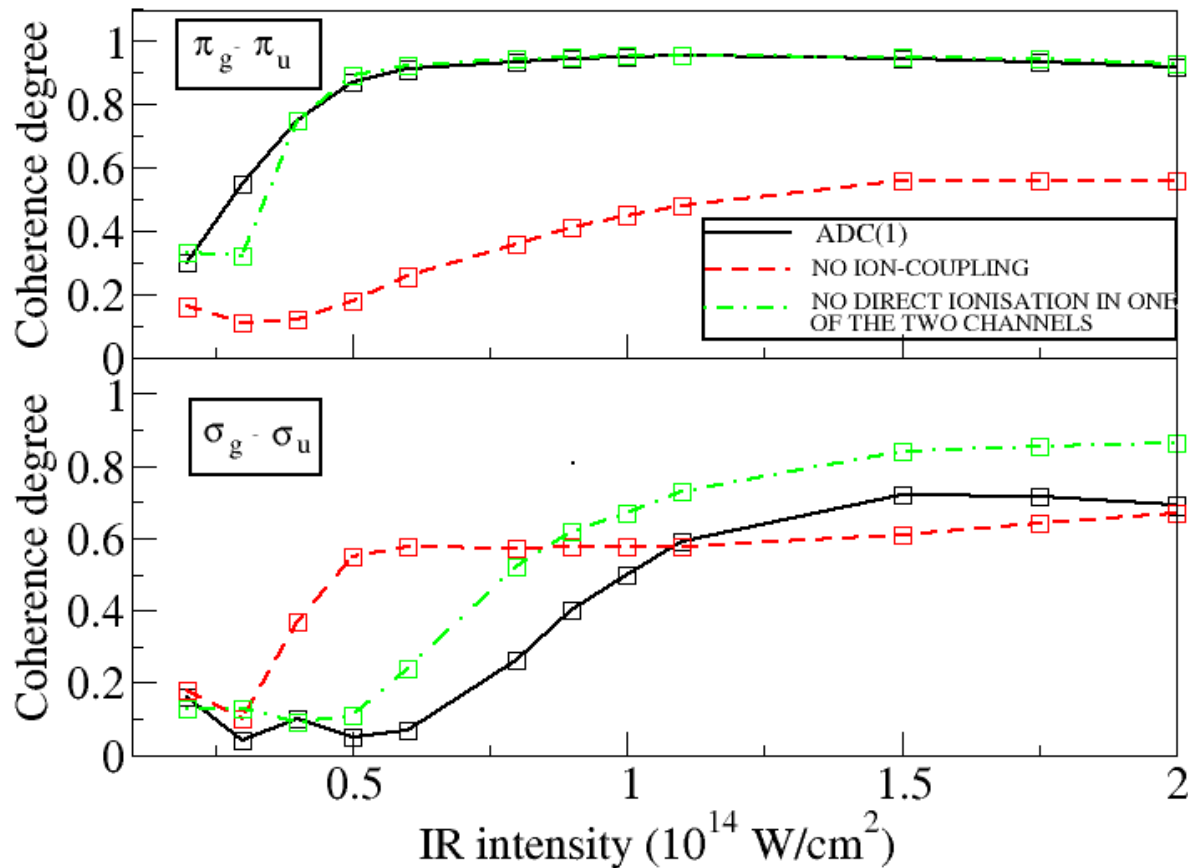
Ion-photoelectron coupling



Dipole coupling of ionised states

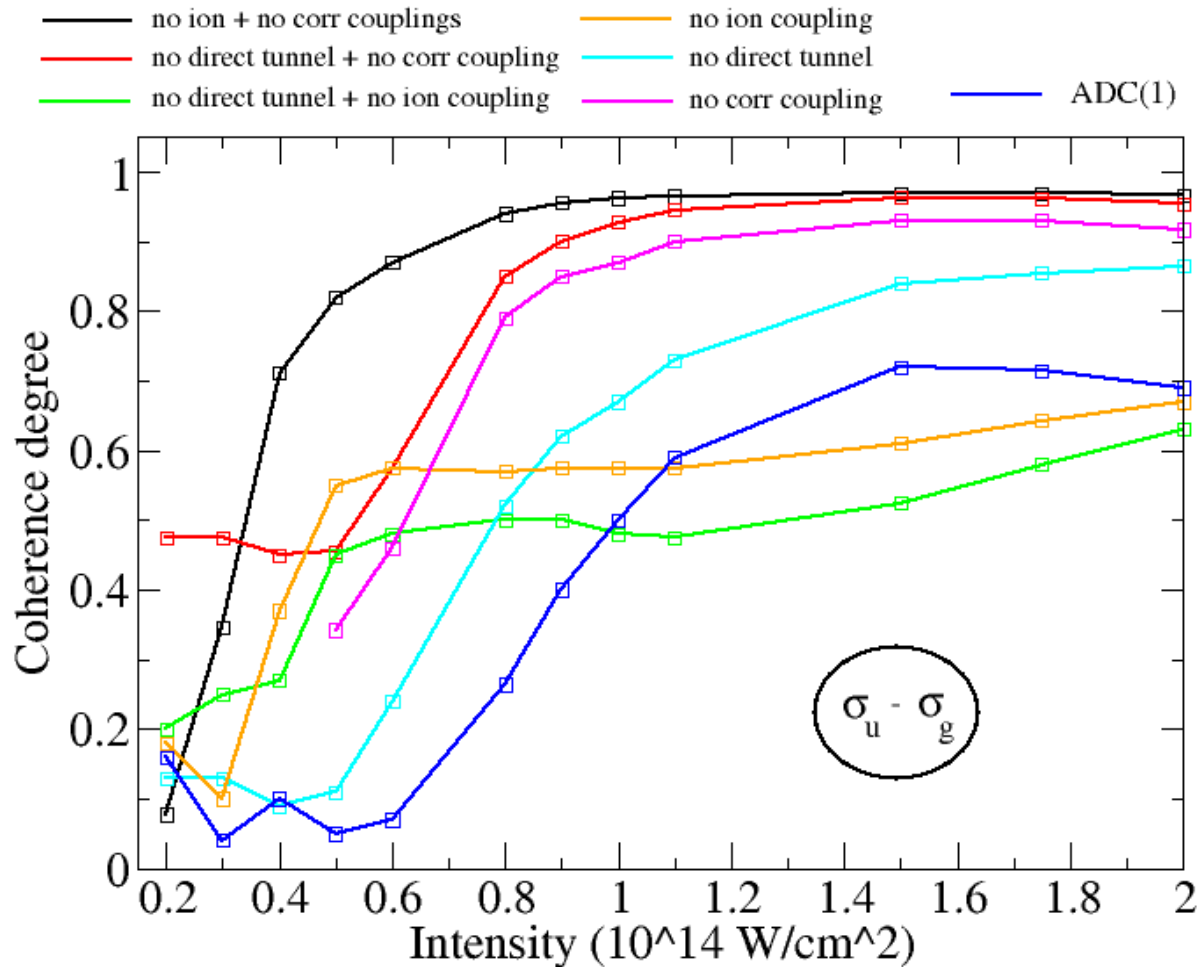


Mechanisms of coherence formation in molecular IR strong-field ionization



★ In some cases, the leading mechanism can be identified

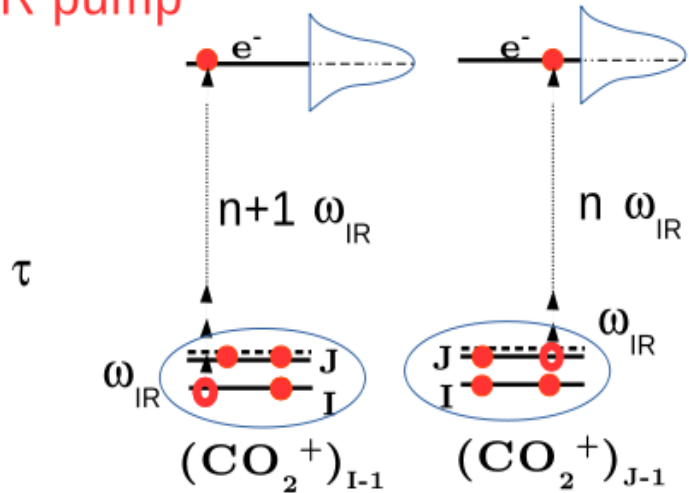
Mechanisms of coherence formation in molecular IR strong-field ionization



★ In general, coherence is formed as non-additive superposition of three different effects

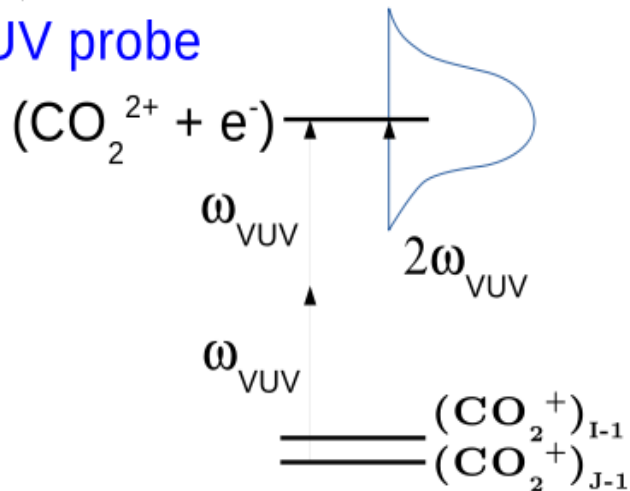
Probing the ionic coherence

IR pump



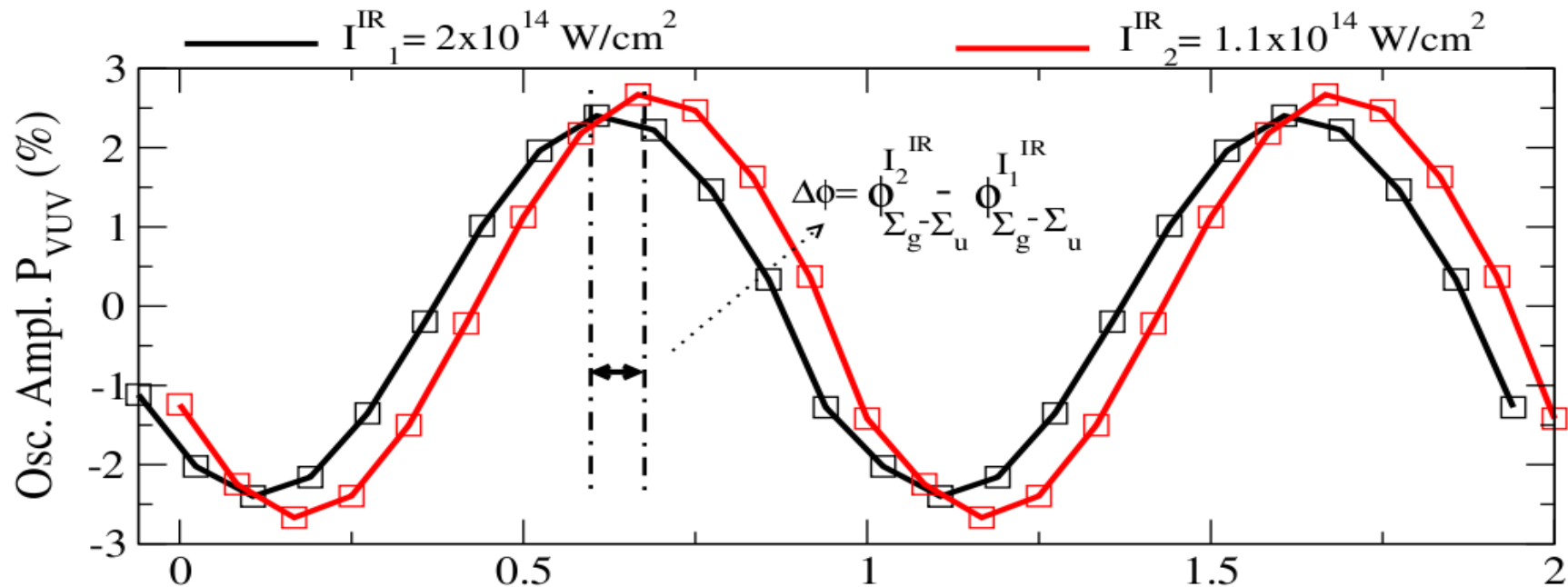
★ Take the ionic state coefficients from the N-electron ADC(1) wavefunction analysis and plug them into ADC(2)x for (N-1)-electron state to simulate the probe

VUV probe



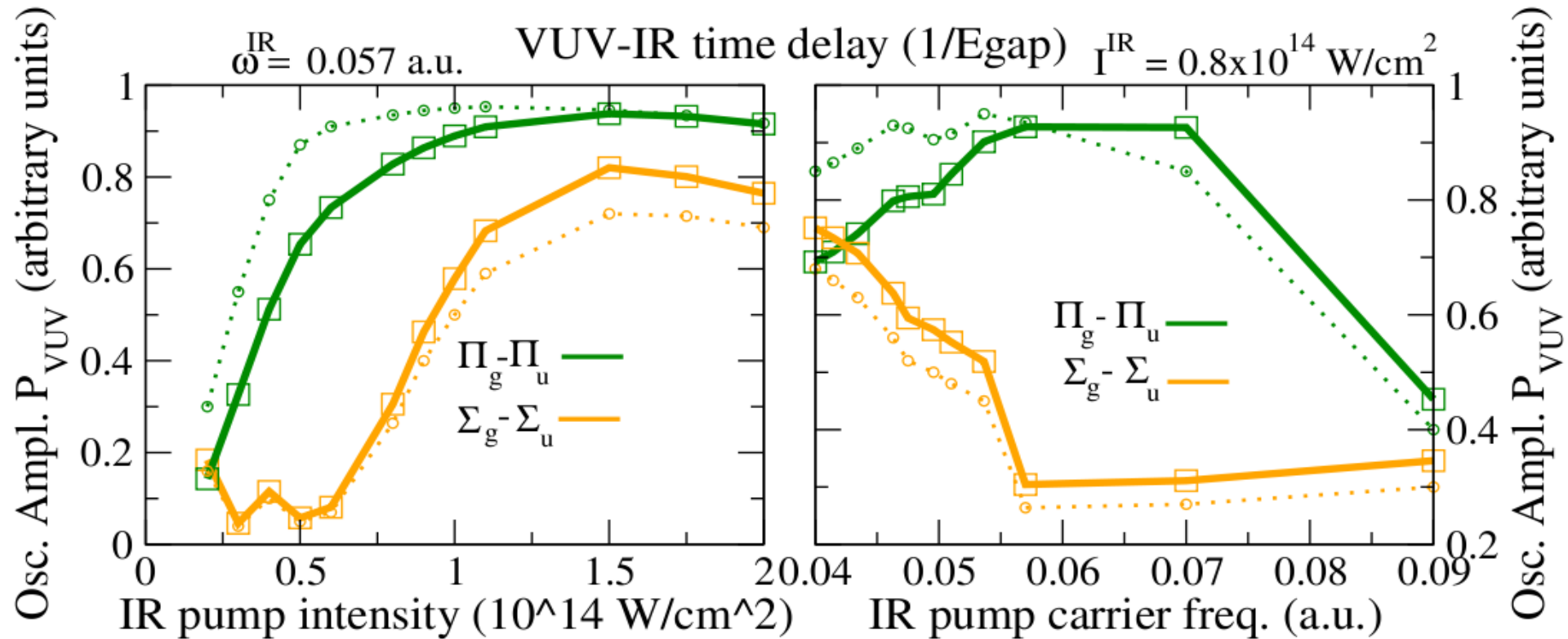
★ ω - 2ω probe is required for symmetry reasons, for a general molecule a single-colour probe would work

Probing the ionic coherence



★ Total yield of CO_2^{2+} follows the two-state ($\Sigma_g - \Sigma_u$) oscillations
Relative phase of the oscillations at two different pump intensities gives relative phase of the coherent superposition

Probing the ionic coherence



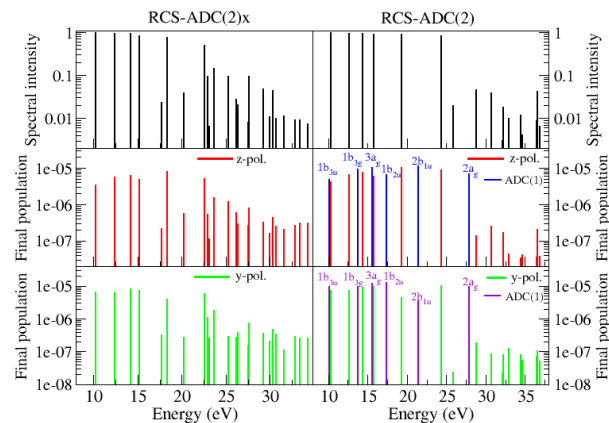
★ Degree of coherence achieved at the pump stage (dotted lines) can be inferred from the amplitude of the CO_2^{2+} yield oscillations

First-principles simulation of IR-pump XUV-probe attosecond spectroscopy of molecules: summary

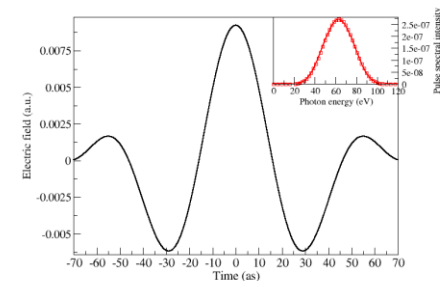
- ✓ Molecular IR pump - XUV probe experiment is simulated fully, with both pump and probe described within the same *ab initio* framework
- ✓ Mechanisms of ionic coherence formation identified and their interplay is revealed through numerical experiments
- ✓ Time-dependent B-spline ADC is a powerful *ab initio* tool ready for further attosecond spectroscopy applications

Reduced ionic density matrix upon attosecond XUV ionization of C_2H_4

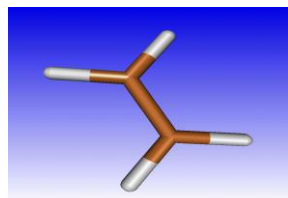
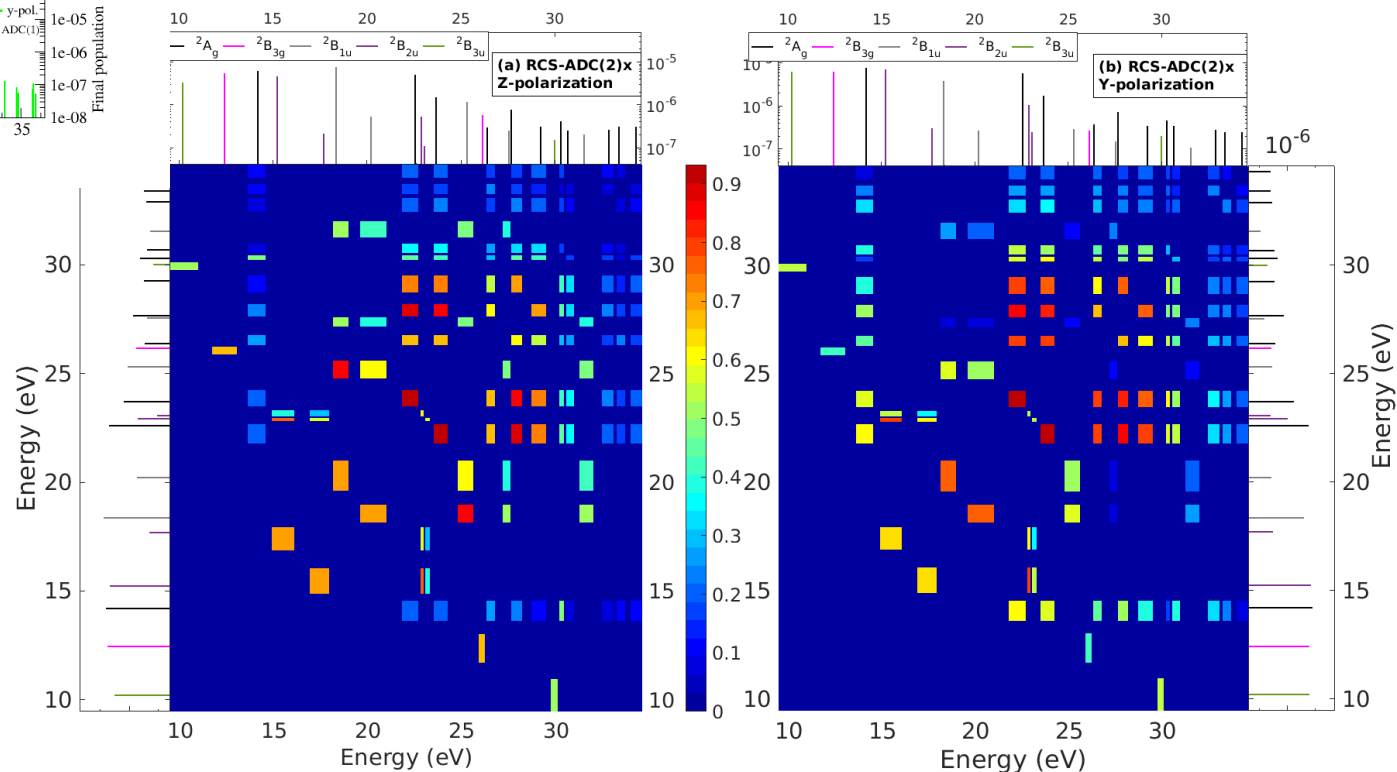
Ionic populations



Mean photon energy = 63 eV
Intensity = 3×10^{11} W/cm²

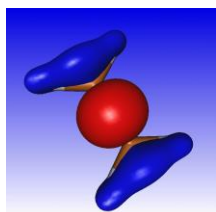
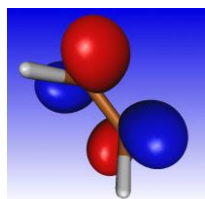
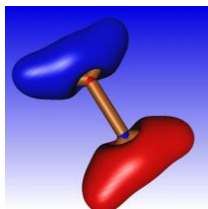
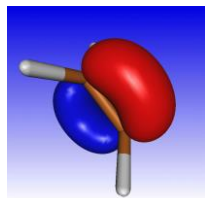


Ionic coherences G_{mn}



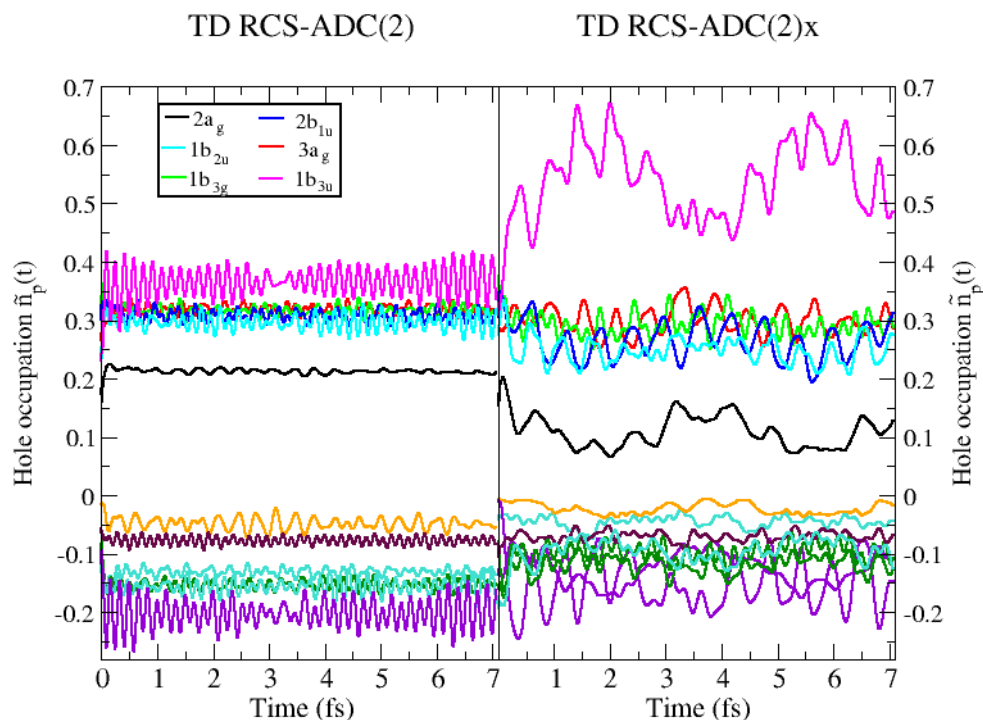
Charge dynamics upon attosecond ionization of C_2H_4

Natural charge orbitals



$$\tilde{Q}(r,t) = \langle \Psi_0^N | \hat{Q}(r) | \Psi_0^N \rangle - Tr(\hat{Q}(r) \hat{\rho}(t))$$

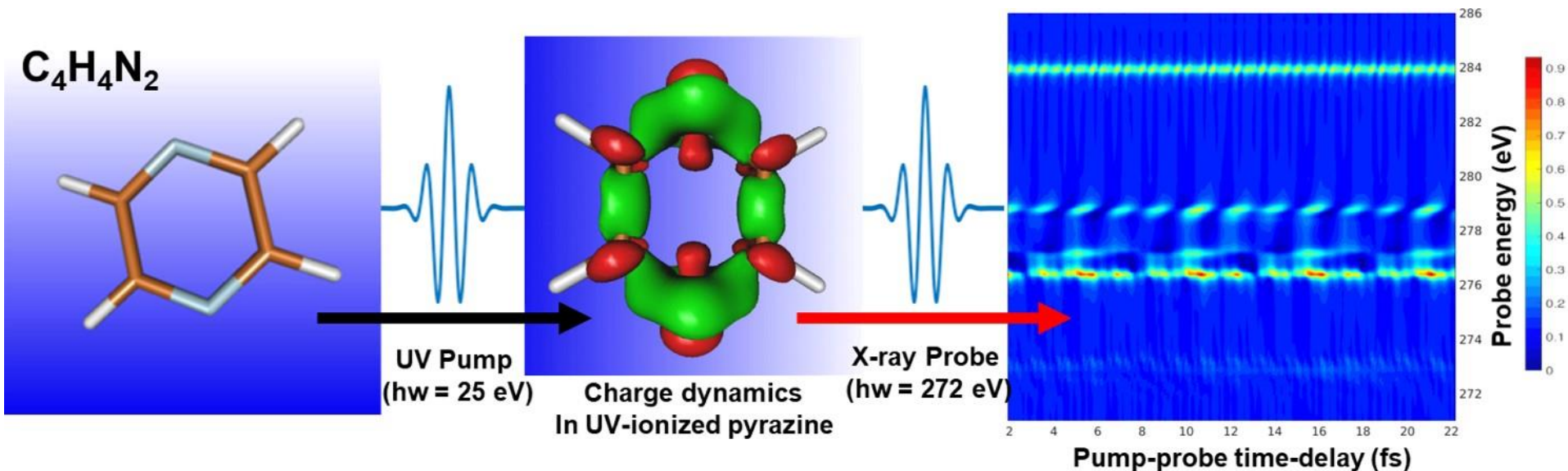
$$\tilde{Q}(r,t) = \sum_p |\tilde{\phi}_p(r)|^2 \tilde{n}_p(t)$$



Dominance of
correlation-driven
charge dynamics

TD B-spline RCS-ADC Theory Applied to Electron Hole Wavepacket Coherence in Pyrazine:

Complete simulation of full attosecond XUV-pump - X-ray-probe experiment in a polyatomic molecule



PUMP: Characterization of ionized system prepared by the XUV pump: Ionic density matrix, coherence and charge dynamics upon attosecond XUV ionization.

PROBE: X-ray attosecond transient absorption spectra.

PUMP attosecond ionization of pyrazine

Populations
of ionic state n

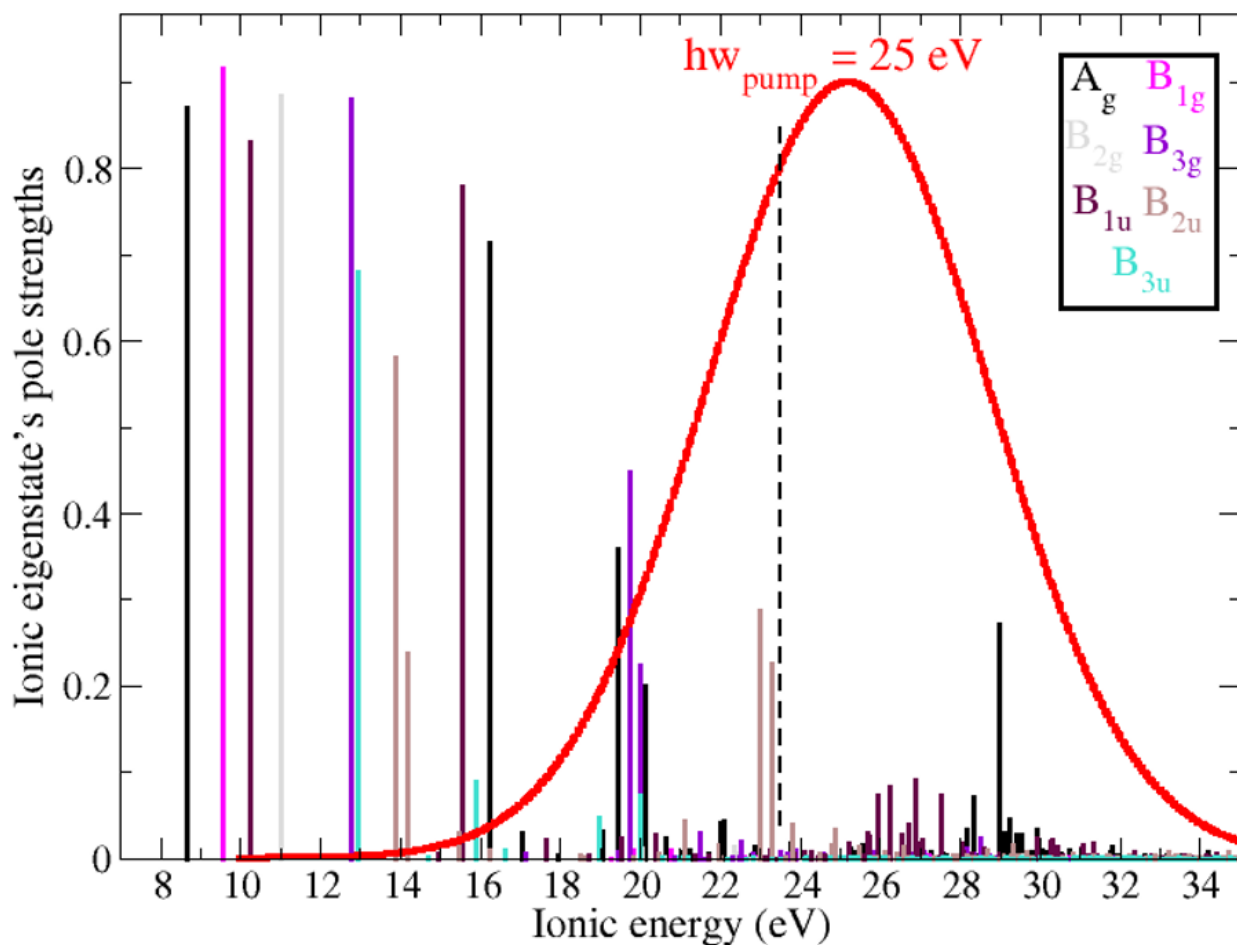
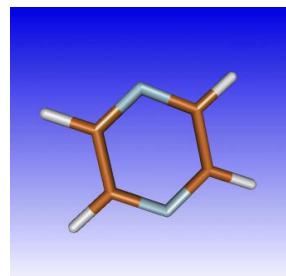
$$P_n(t) = |\rho_{nn}(t)|$$

Pump central photon energy

25 eV

Intensity **$6 \times 10^{11} \text{ W/cm}^2$**

Bandwidth = **7 eV**



PUMP attosecond ionization of pyrazine

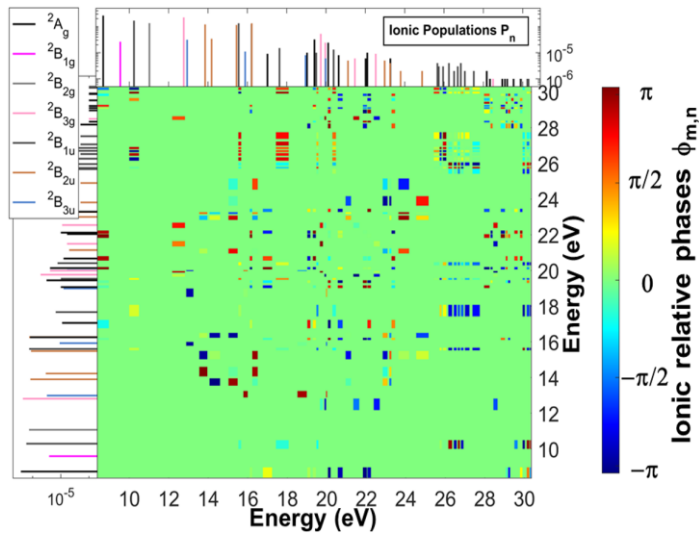
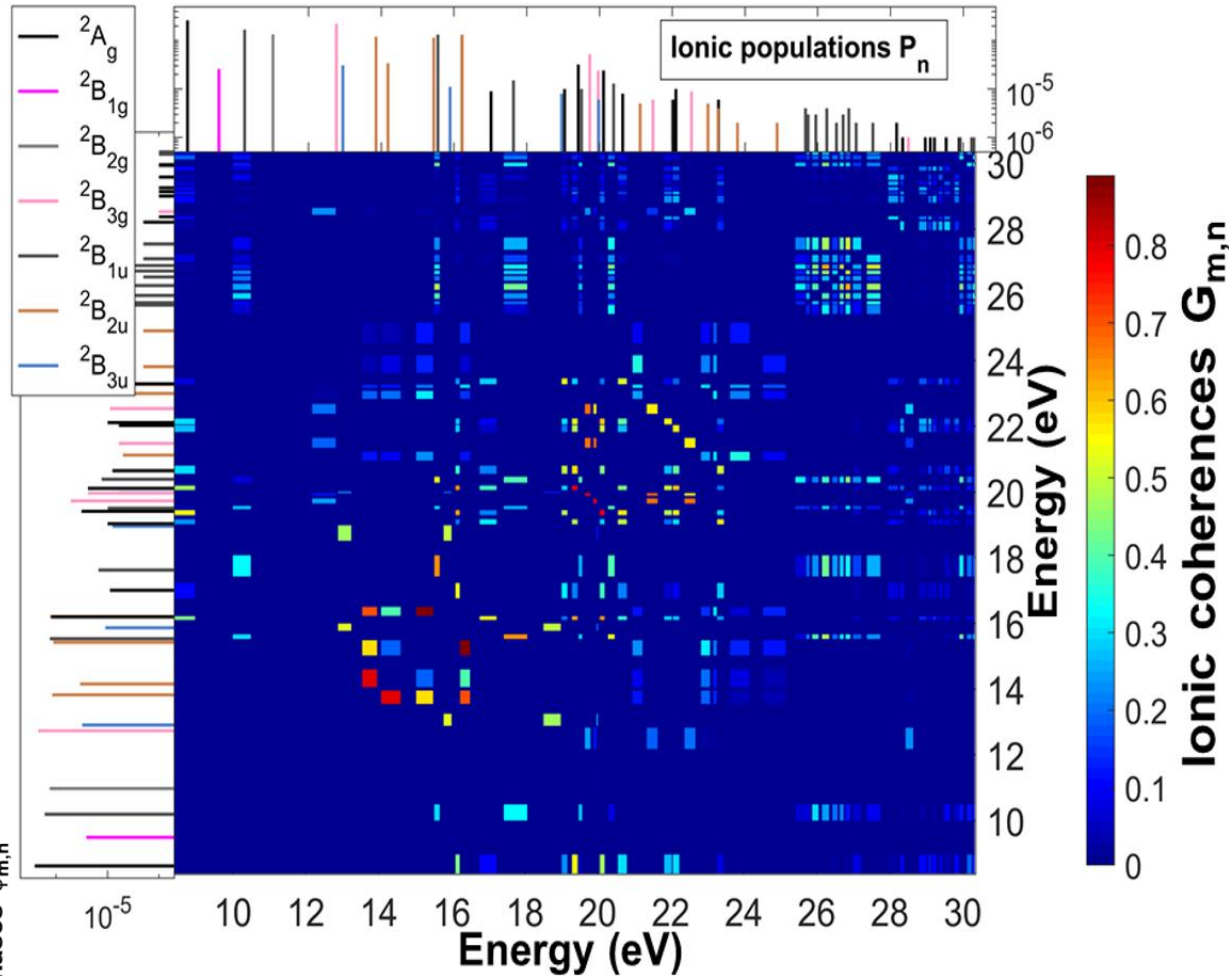
Degrees of coherence

between pair of populated ionic states m and n

$$G_{mn}(t) = \frac{|\rho_{mn}(t)|}{\sqrt{P_m(t) * P_n(t)}}$$

Maximum quantum coherence corresponds to $G_{mn} = 1$

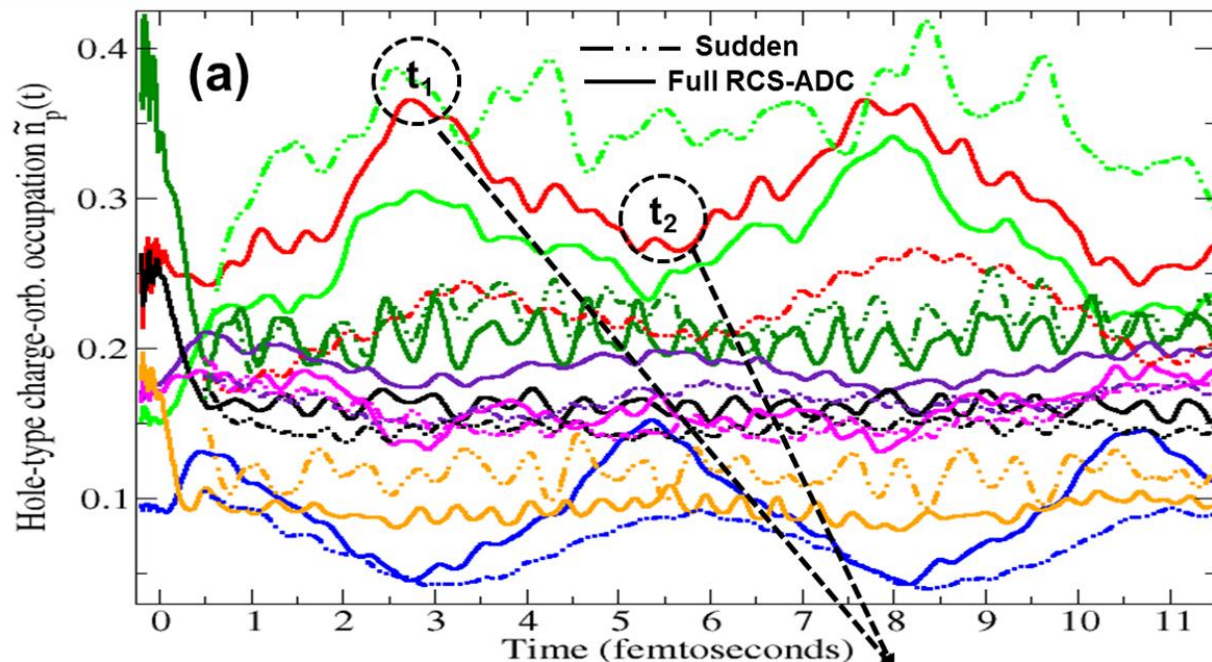
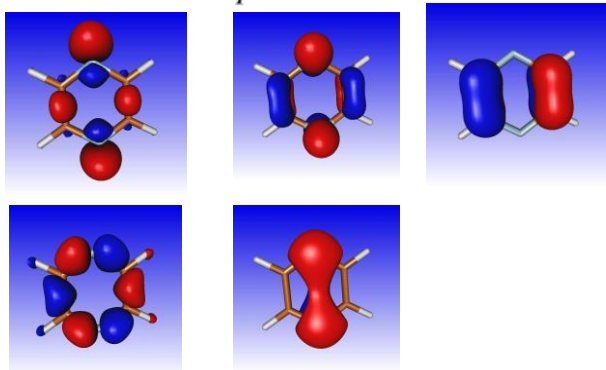
Ionic relative phases



Deviation from sudden approximation for smaller photon energies: the case of pyrazine

Natural charge orbitals

$$\tilde{Q}(r,t) = \sum_p |\tilde{\phi}_p(r)|^2 \tilde{n}_p(t)$$



Sudden approximation ansatz

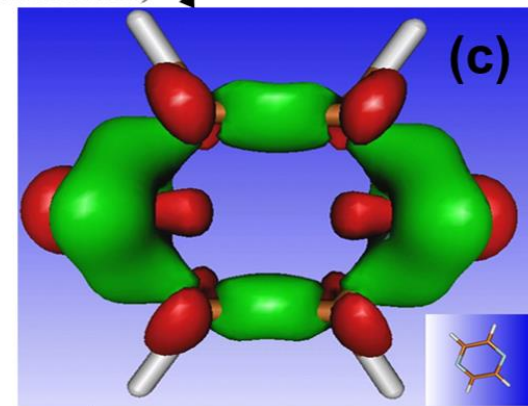
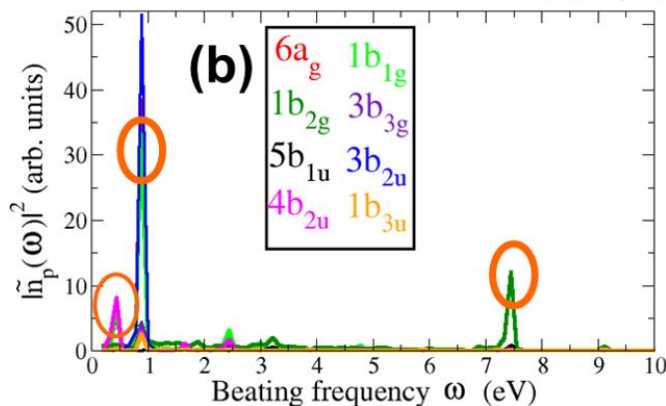
for the initial state

? Dependence on

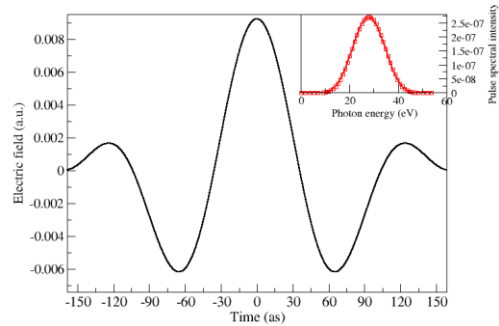
laser pulse parameters ?

$$|\Psi_{initial}^{N-1}\rangle \sim a_i |\Psi_0^N\rangle$$

when is it valid...?



Ionic coherence evolution upon XUV ionization of pyrazine



Ionic decoherence due to coupling to the emitted photoelectron

Purity

$Tr(\hat{\rho}^2(t)) \neq Tr(\hat{\rho}(t)) \leq 1$ mixed state,

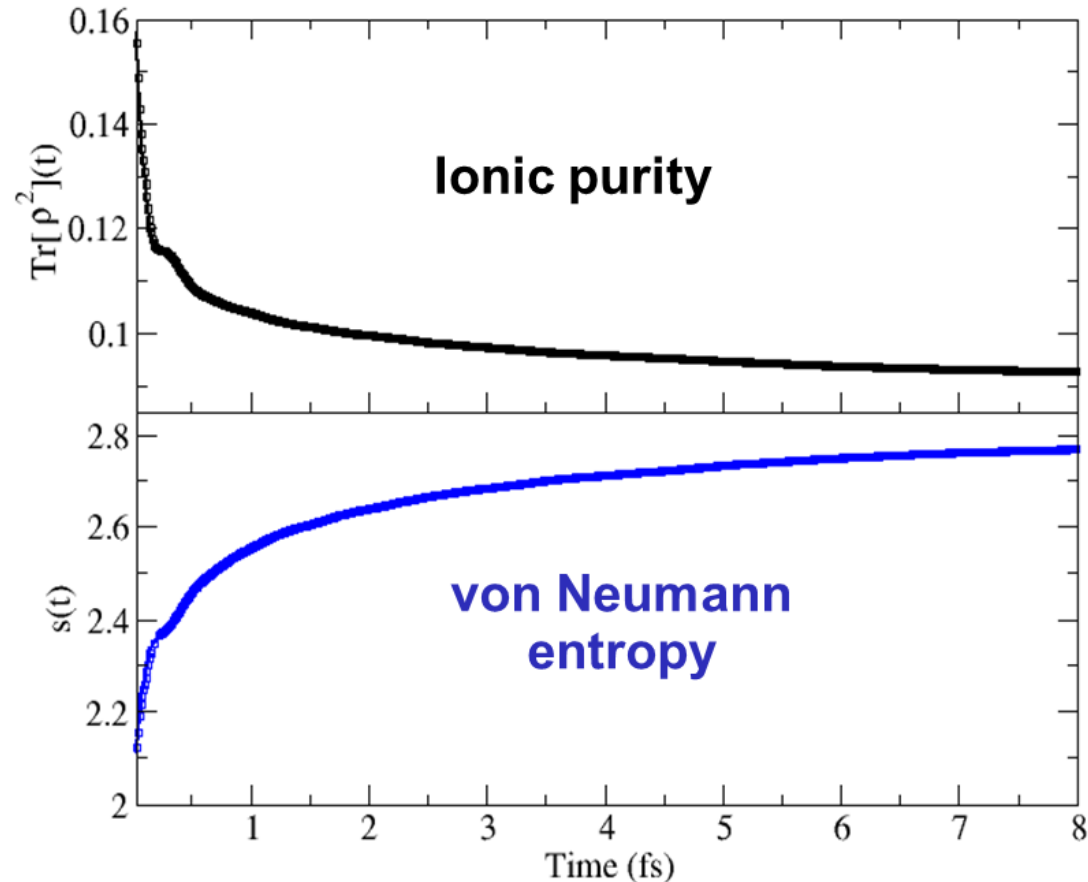
$Tr(\hat{\rho}^2(t)) = Tr(\hat{\rho}(t)) = 1$ pure state.

von Neumann

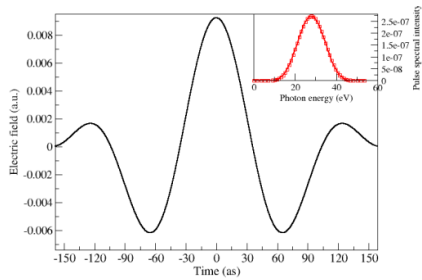
entropy of entanglement

$s^{Ion}(t) = -\ln\left(\frac{1}{N}\right)$ statistical mixture,

$s^{Ion}(t) = 0$ pure state.

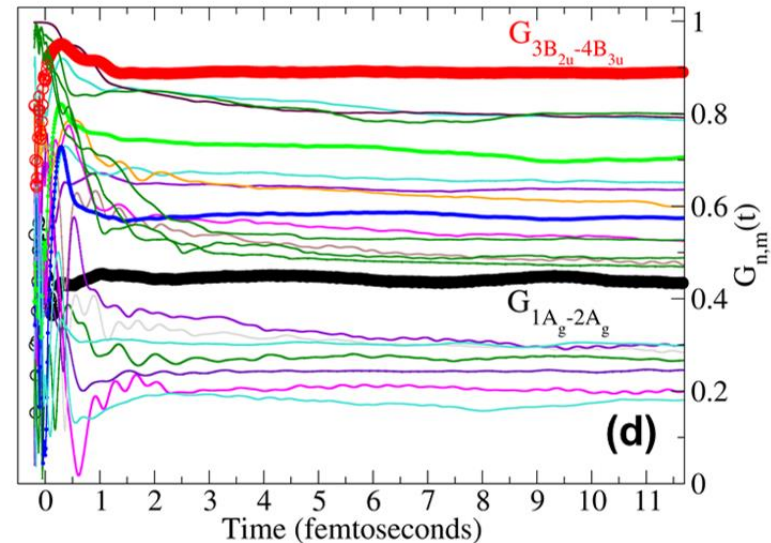
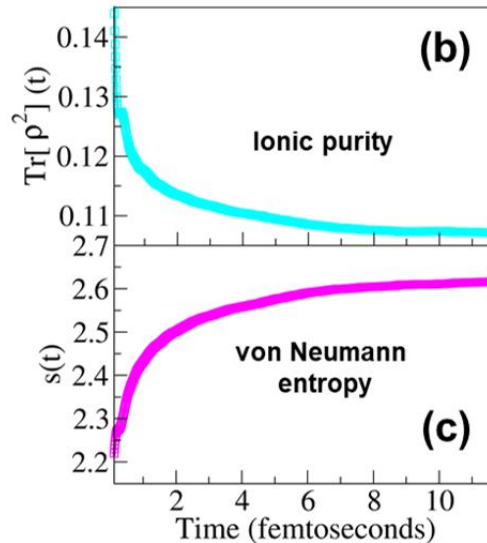
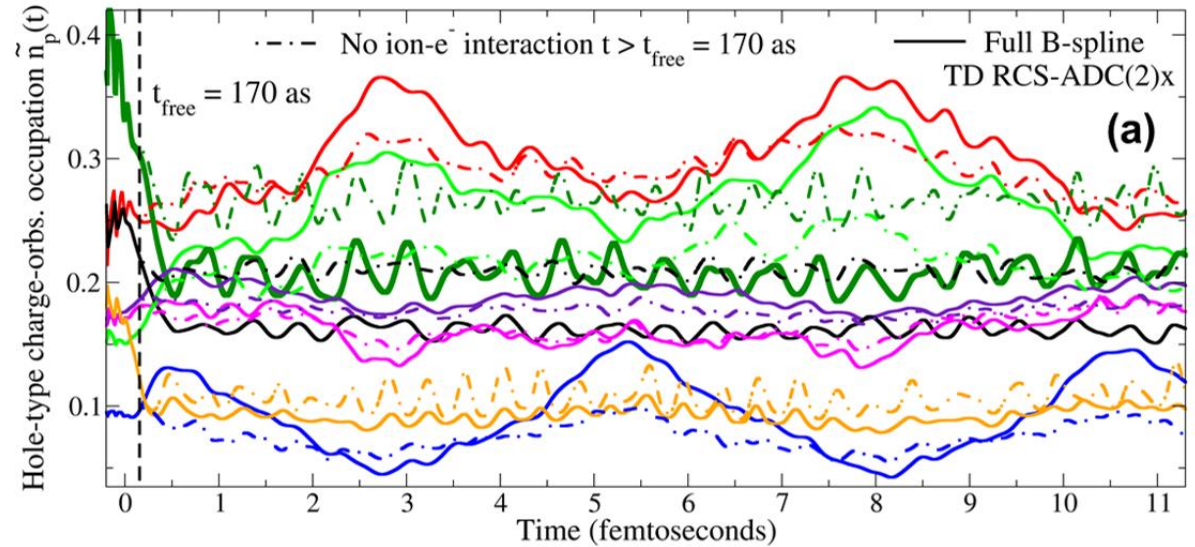


Ionic coherence evolution upon XUV ionization of pyrazine



Central Pump
 photon energy = 25 eV
 Intensity = 6×10^{11} W/cm²

Ionic decoherence due to coupling to the emitted photoelectron

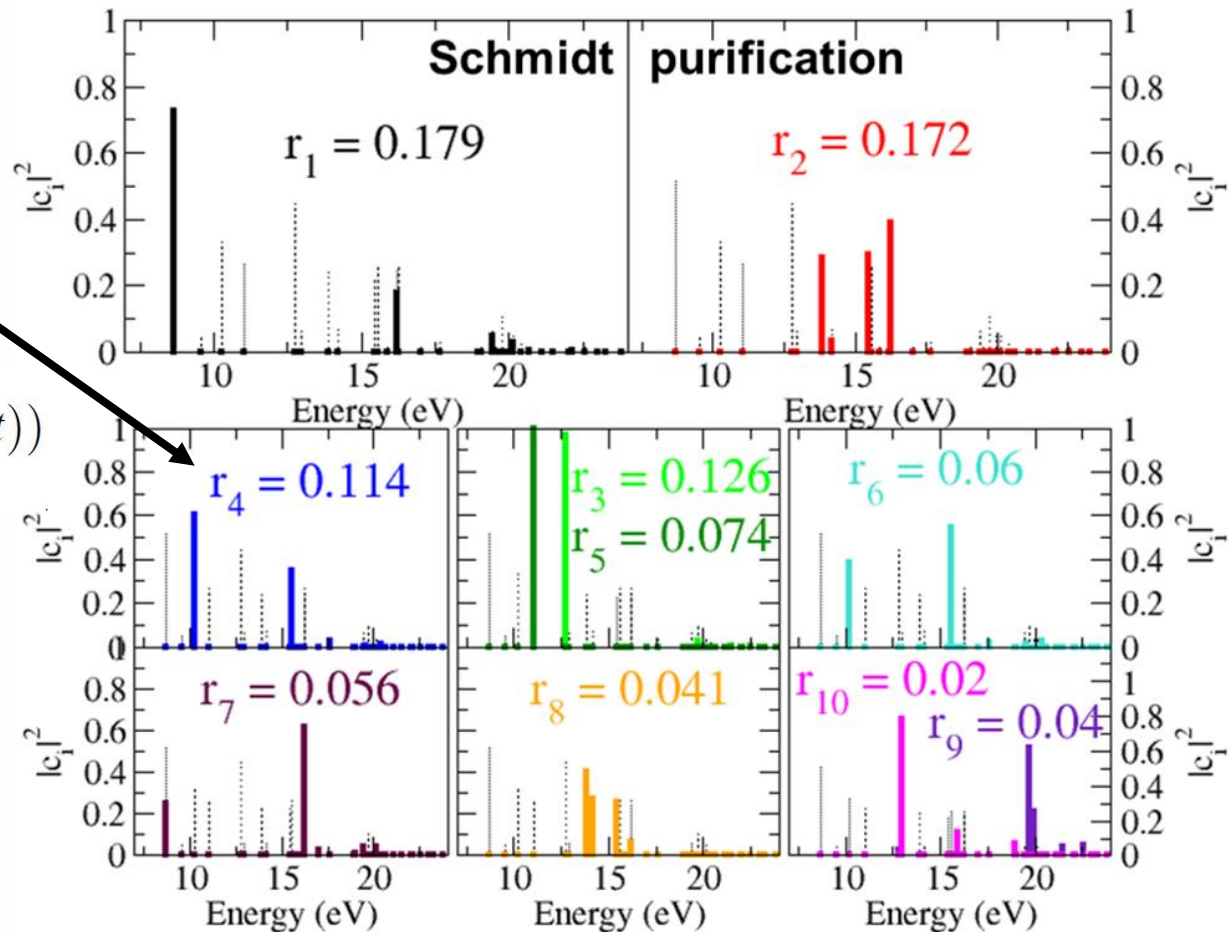


Schmidt decomposition of ionic density matrix

Schmidt decomposition of ionic density matrix:
incoherent sum of coherent ionic channels –
each wavefunction has weight r_i

$$\hat{\rho}(t) = \sum_{i=1,N} r_i(t) \hat{P}_i(t)$$

$$S^{\text{Ion}}(t) = - \sum_{i=1,N} r_i(t) \times \ln(r_i(t))$$



PROBE-STEP

X-ray attosecond transient absorption spectra of XUV-ionized pyrazine

First-principles simulation of the **probe step**:

Time-Dependent von Neumann equations for the ionic density matrix interacting with the X-ray probe pulse

$$\frac{d}{dt}\hat{\rho}(t) = -\frac{i}{\hbar}\left[\hat{H}^{\text{Ion}}(t), \hat{\rho}(t)\right]$$

X-ray probe duration = 180 as; Intensity = 10^{14} W/cm²

Photon energies:

N K-edge absorption window
photon energy = 393 eV

C K-edge absorption window
photon energy = 272 eV

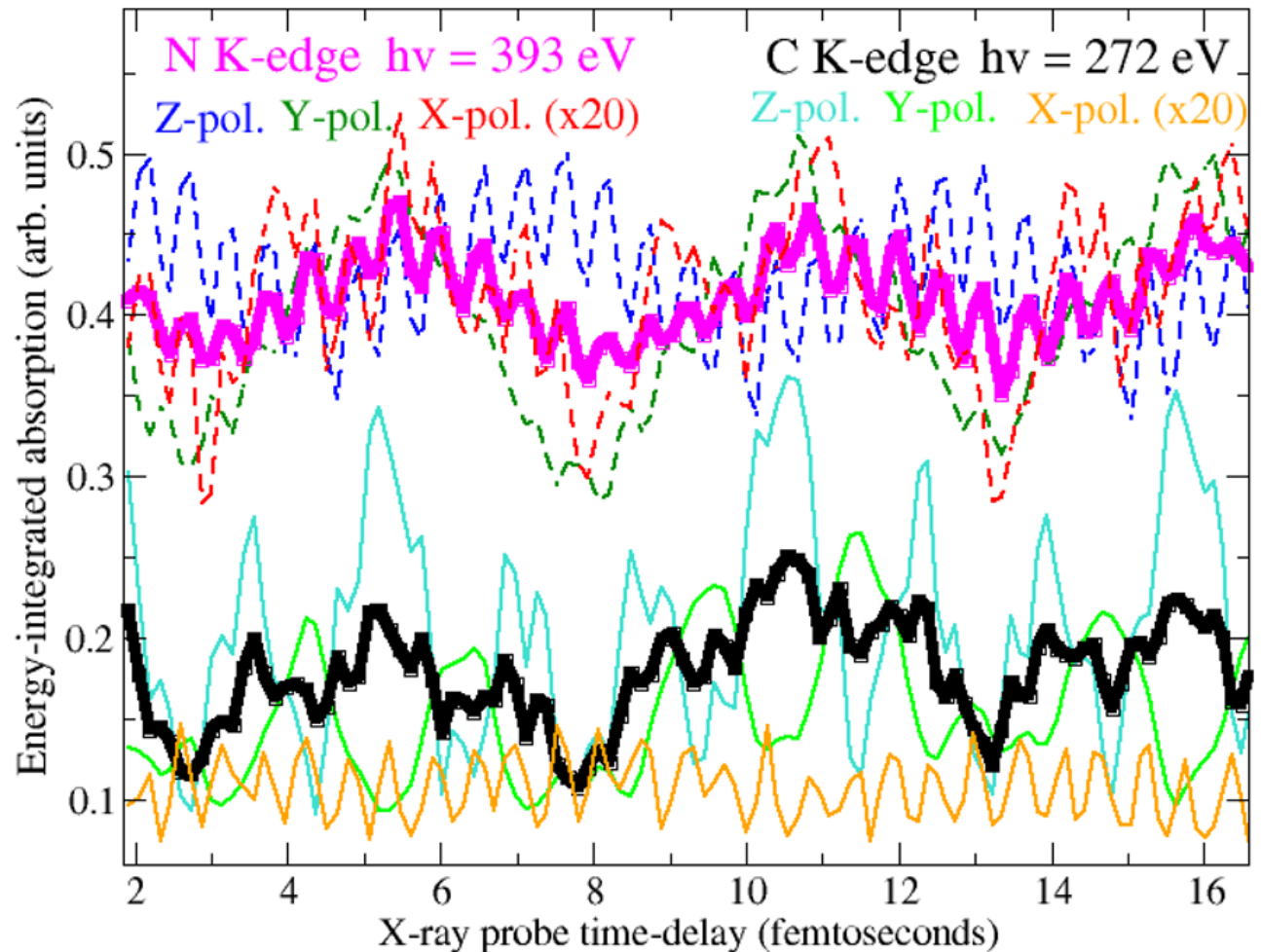
PROBE-STEP

X-ray energy-integrated transient absorption of XUV-ionized pyrazine

Attosecond
(XUV) pump -
(X-ray) probe
spectroscopy
of charge dynamics:

- high temporal resolution;
- chemical selectivity;
- spatial resolution.

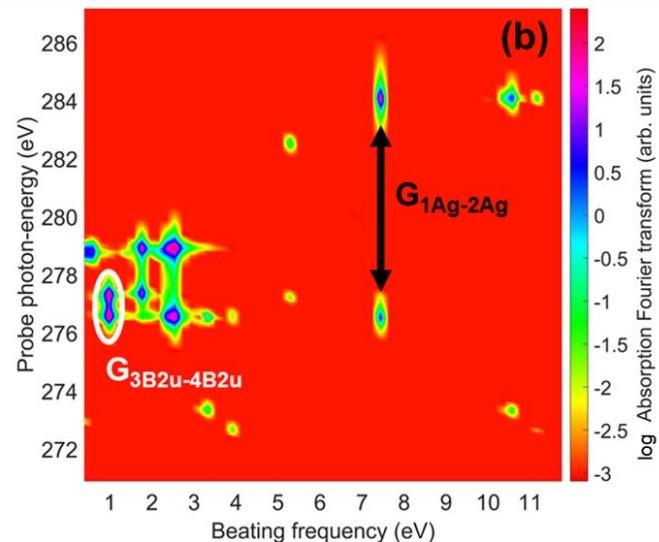
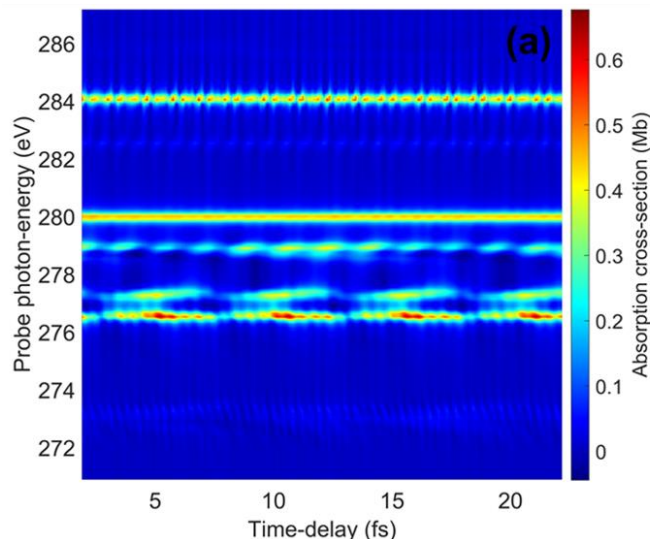
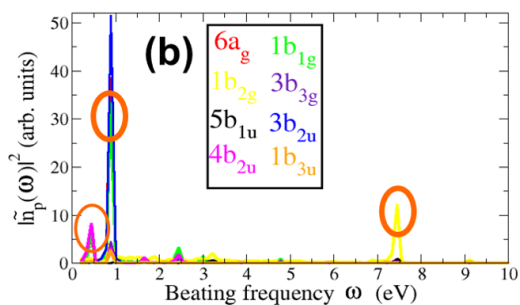
**C-N charge oscillation
period $T \sim 0.54$ fs**



Quantum electronic coherences

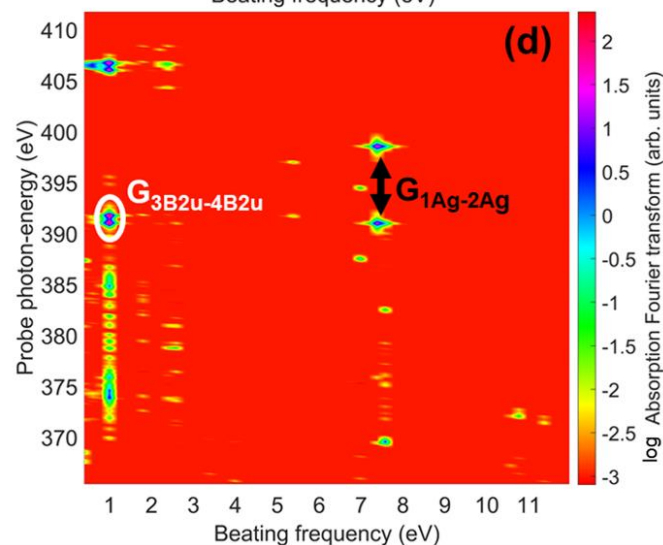
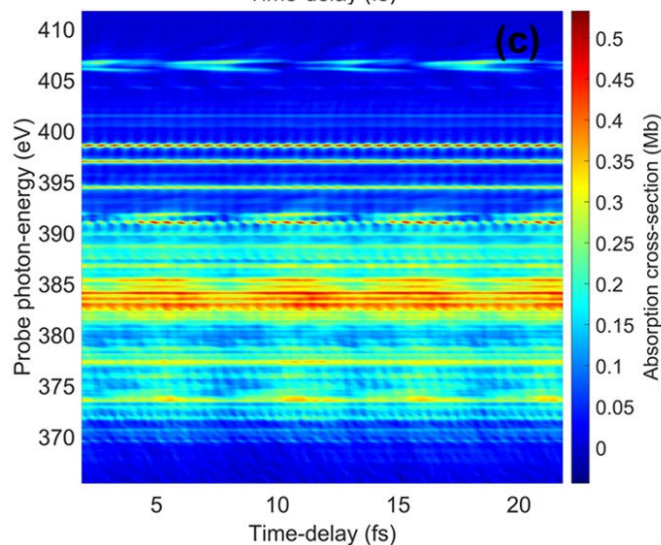
by attosecond transient absorption spectroscopy

Carbon K-edge photon energy = 272 eV



Nitrogen K-edge photon energy = 393 eV

All (1h and 2h1p)
final core-ionised
states included

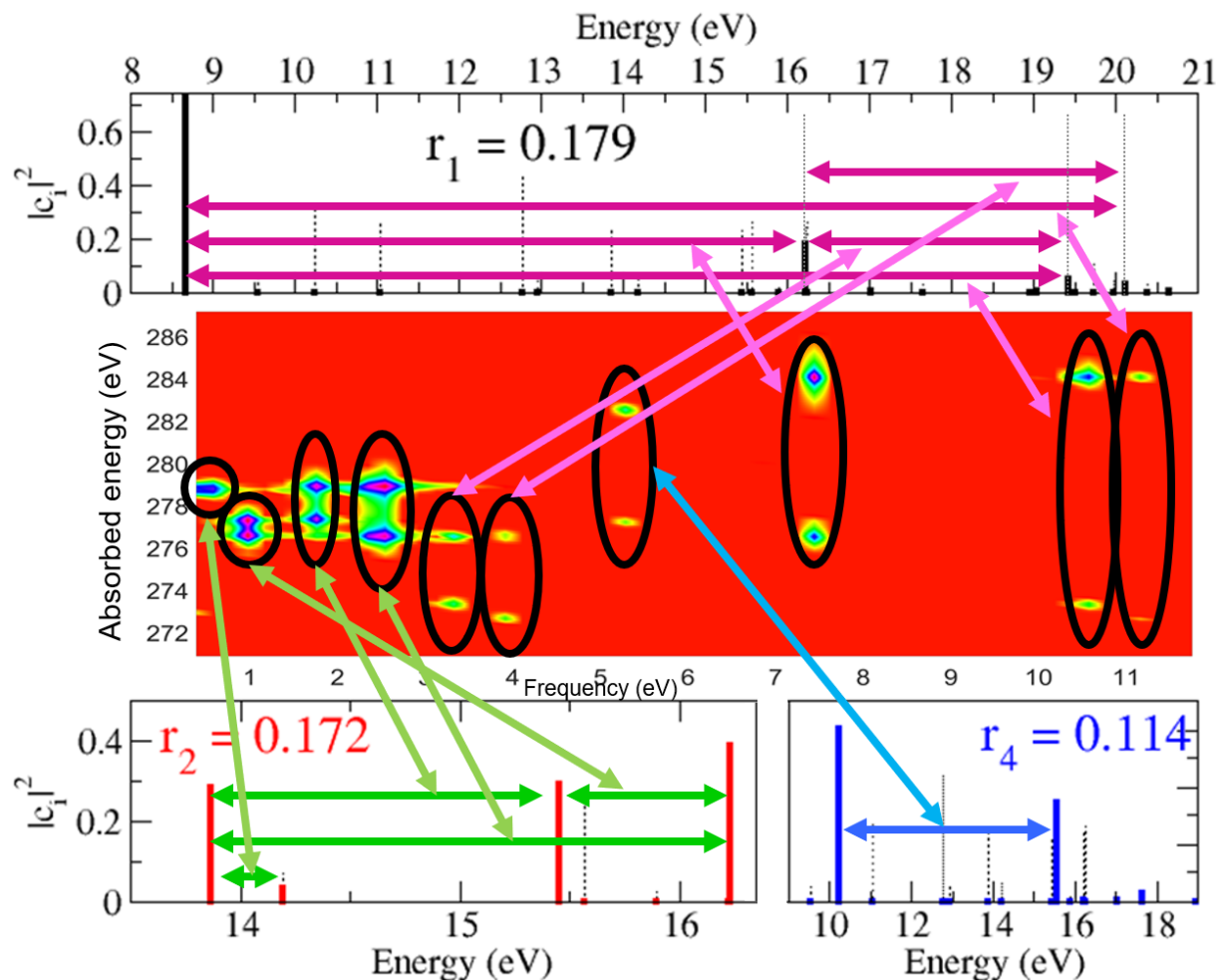


M. Ruberti,
Faraday Discussions (2020)

PROBE-STEP

Mapping ATAS observable onto ionic coherences

Quantum electronic coherences by attosecond transient absorption spectroscopy



Peaks in the Fourier-transformed ATAS

$$\hat{\rho}(t) = \sum_{i=1,N} r_i(t) \hat{P}_i(t)$$

Schmidt decomposition of ionic density matrix

M. Ruberti,
Faraday Discussions (2020)

B-spline ADC, RCS-ADC Publications

B-spline ADC theory

2014 -Atomic **photoionization cross-section** and **high harmonic generation spectra**, J. Chem. Phys.

2016 -**XUV-IR transient absorption** in helium New J. Phys.

2018 -Book chapter in “Attosecond Molecular Dynamics”, RSC Theoretical and Computational Chemistry series.

2018 -Multi-channel dynamics in **high harmonic generation** of CO₂, Phys. Chem. Chem. Phys.

2018 -**IR-pump – XUV-probe attosecond experiment** in CO₂, J. Chem. Theory Comput.

B-spline RCS-ADC theory

2019 -Total **photoionization cross-sections** of molecules, J. Chem. Theory Comput.

2019 -**Ionic coherence and charge dynamics** in XUV attosecond ionization of C₂H₂ and C₂H₄, Phys. Chem. Chem. Phys.

2019 -**Coherent control** of single-photon Laser Enabled Auger Decay in Neon, New J. Phys.

2020 -XUV pump – X-ray probe **attosecond transient absorption spectroscopy** in pyrazine, Faraday Discussions.

2021 -**Time-Resolved photoelectron spectroscopy**: Femtosecond X-ray pump – X-ray probe in glycine, [arXiv:2012.04852](https://arxiv.org/abs/2012.04852), under review in Science.

First-principles simulation of XUV-pump X-ray probe attosecond spectroscopy of molecules: summary

- Time-dependent B-spline RCS-ADC *ab initio* method describes many-electron ionization dynamics in polyatomic molecules with inclusion of relaxation & correlation effects (shake-up) and full description of photoelectron.
- Quantum electronic coherence/entanglement formation can be calculated *ab initio*, in the molecular case, as a function of the ionizing pulse parameters.
- Going beyond the sudden approximation is necessary especially in the lower photoelectron energy regime.
- Fully *ab initio* simulations of complete attosecond pump - attosecond probe experiments possible in polyatomic molecules.
- X-ray ATAS set-up is a promising and powerful scheme to retrieve the many-electron quantum ionic coherences arising in attosecond photoionization and pre-determining the subsequent charge-directed reactivity.
- Quantum coherence decay mechanisms: residual interaction with slow photoelectrons, Auger decay, coupling to nuclear motion, vibrational ground-state spatial spread. Unravelling the interplay between these mechanisms is the next theoretical, computational and experimental challenge...

Quantum State Tomography for Attosecond Ionization of Entangled Many-Electron Systems

- ❖ Fundamental understanding of the physics underlying photochemical & photophysical transformations of matter at attosecond to femtosecond timescales
- ❖ *Extend and apply concepts of Quantum Information to the field of attosecond dynamics*

**Full characterisation
+
tomographic reconstruction
+
coherent quantum control
of
X-ray atto-ionised
quantum state**

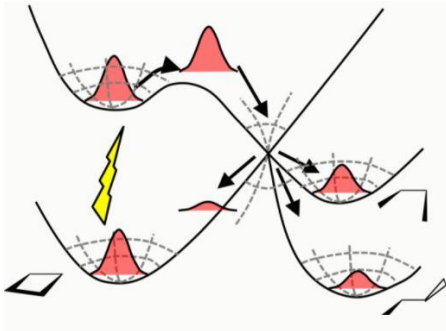
Research Vision- Theoretical “ingredients”

- ✓ Description of **electron correlation**: shakeup states, correlation satellite states.
- ✓ Description of **electronic relaxation**: relaxation satellite states.
- ✓ **Accurate partial cross-sections** for the many-electron ionic states of the system as a function of the ionising laser pulse parameters.
- ✓ **Time-dependent description** of the ionisation process: ultrafast formation and loss of coherence and effect of different laser field parameters.
- ✓ Description of the **N-electron wavefunction**, ion-photoelectron **entanglement**.
- ✓ Effect of the residual **inter-channel couplings** between the created ionic system and the **emitted photoelectron. Correlation effects in the continuum.**
- ❖ Description of **non-radiative relaxation processes**: Auger decay, Coster-Kronig. It can be done sequentially, or semi-empirically (introducing extra parameters). The model describes one electron in the continuum.

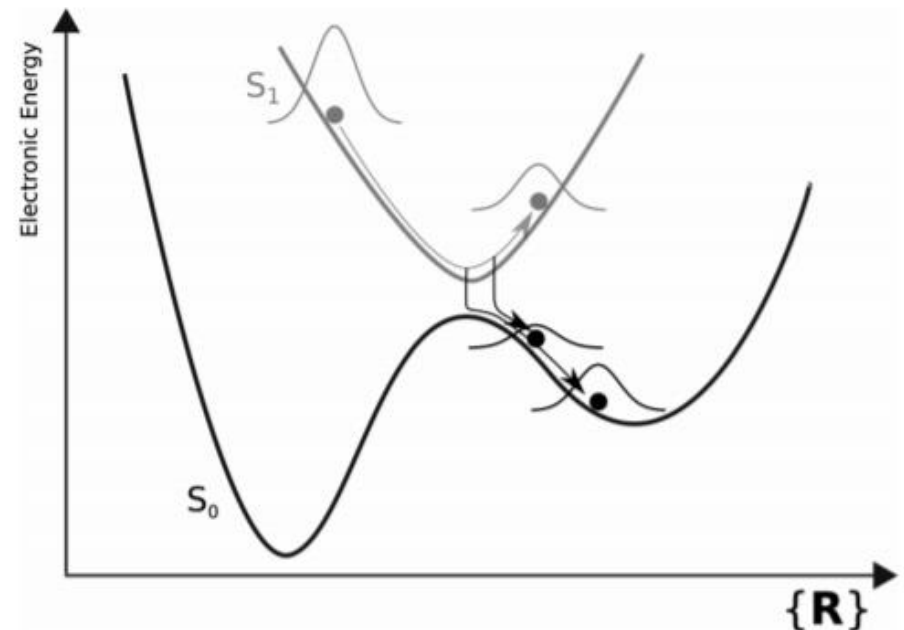
Research Vision - Theoretical “ingredients”

❑ **Coupling to nuclear motion**, multidimensional problem:
electronic coherences among several electronic states →
non-adiabatic nuclear dynamics across several potential energy surfaces

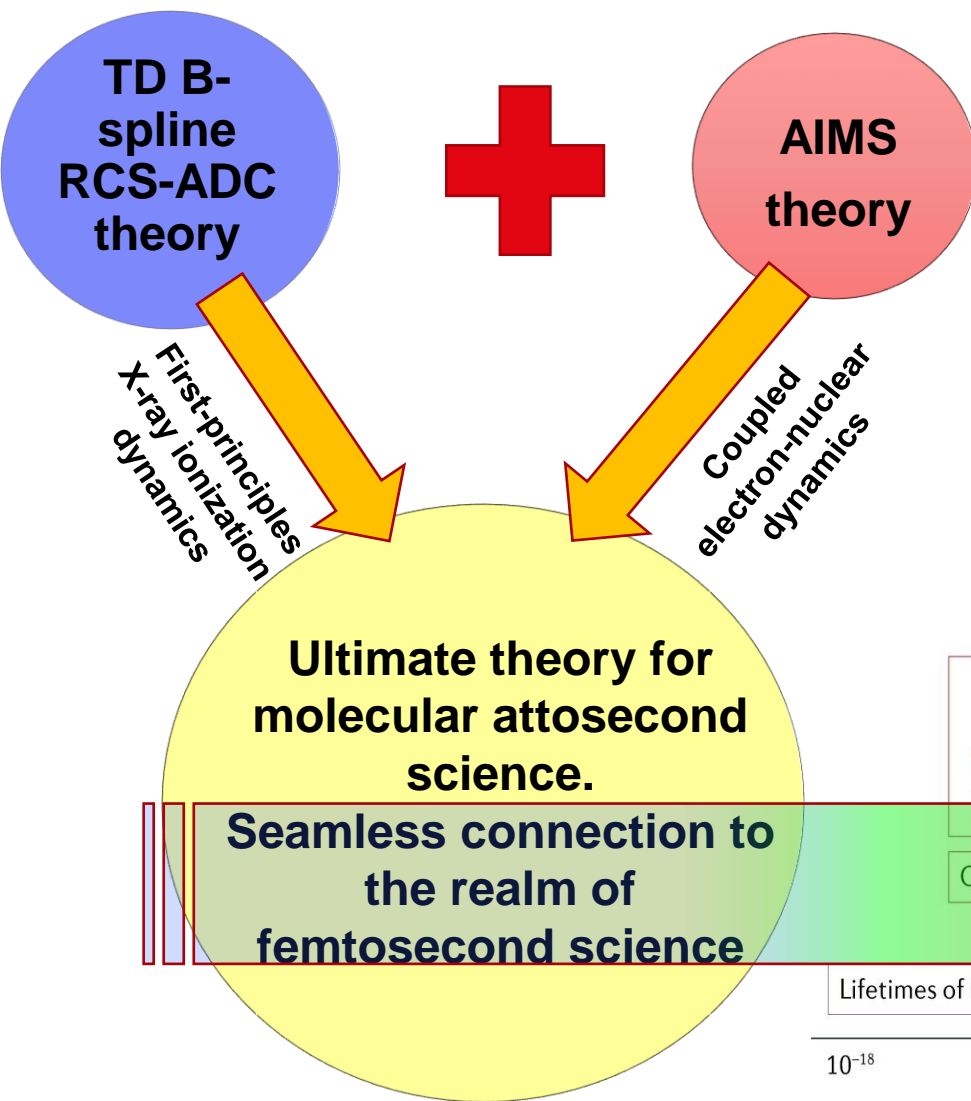
❑ **Complete description of the QUANTUM many-body state of the molecular system.**



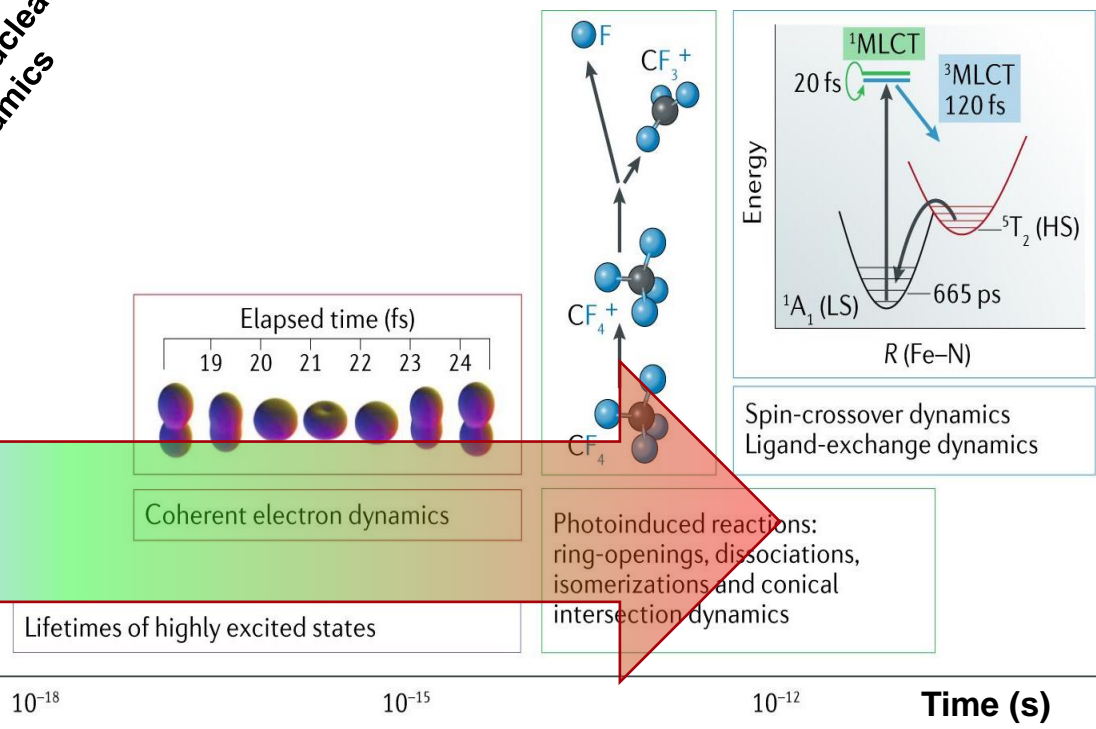
✓ *Ab initio* multiple spawning (**AIMS**)
scheme for coupled electron-nuclear
dynamics, developed by
Prof. T. Martinez (PULSE Institute).



Fully quantum theory for ultrafast coupled electron-nuclear dynamics triggered by X-ray attosecond laser pulses



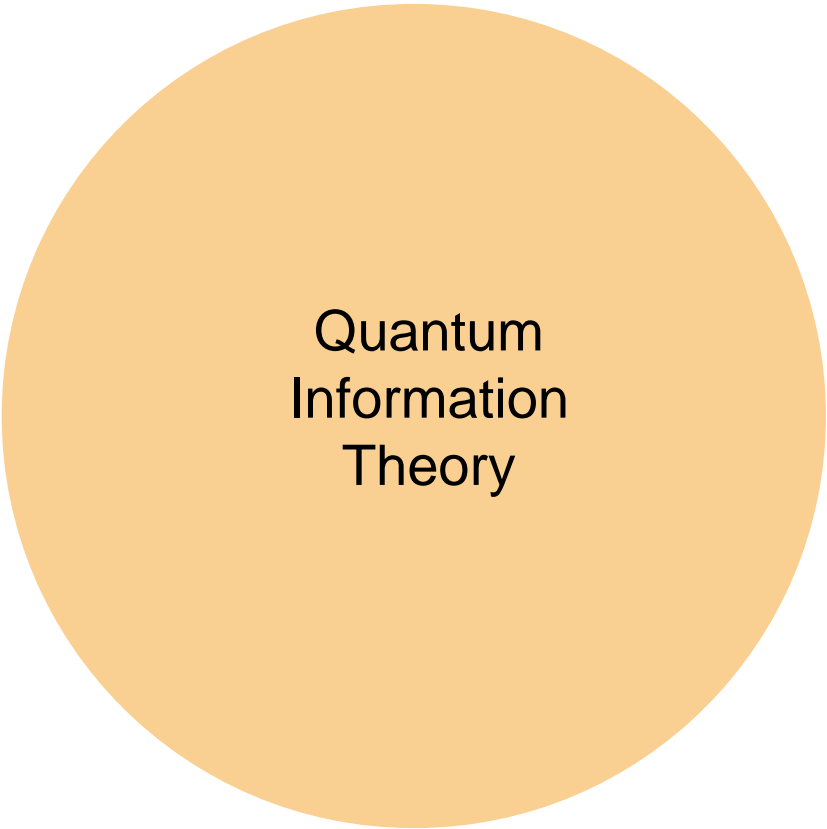
Collaboration with the **Ultrafast Theory and Simulation Group** of the Chemical Science Division at PULSE Institute (Prof. Todd Martinez)



Research Vision

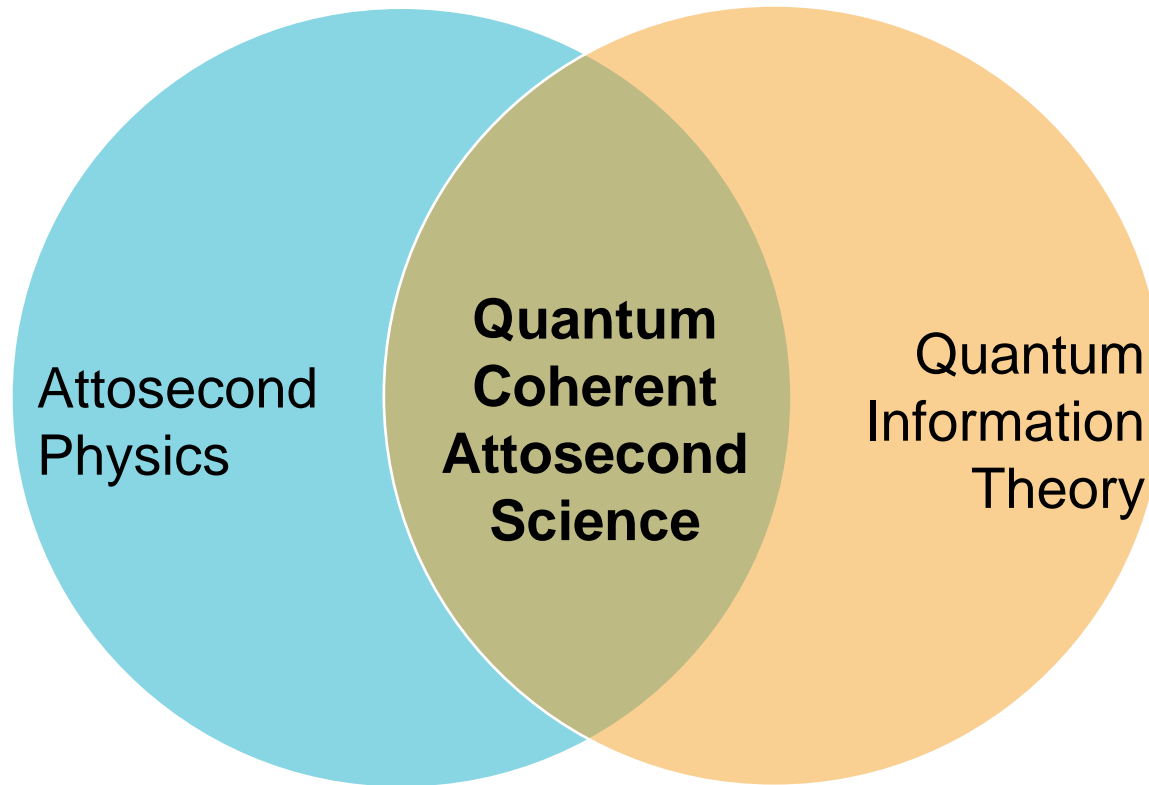


Attosecond
Physics



Quantum
Information
Theory

Research Vision



My goal is to lay the foundations of a new field of study at the intersection of Attosecond Physics and Quantum Information

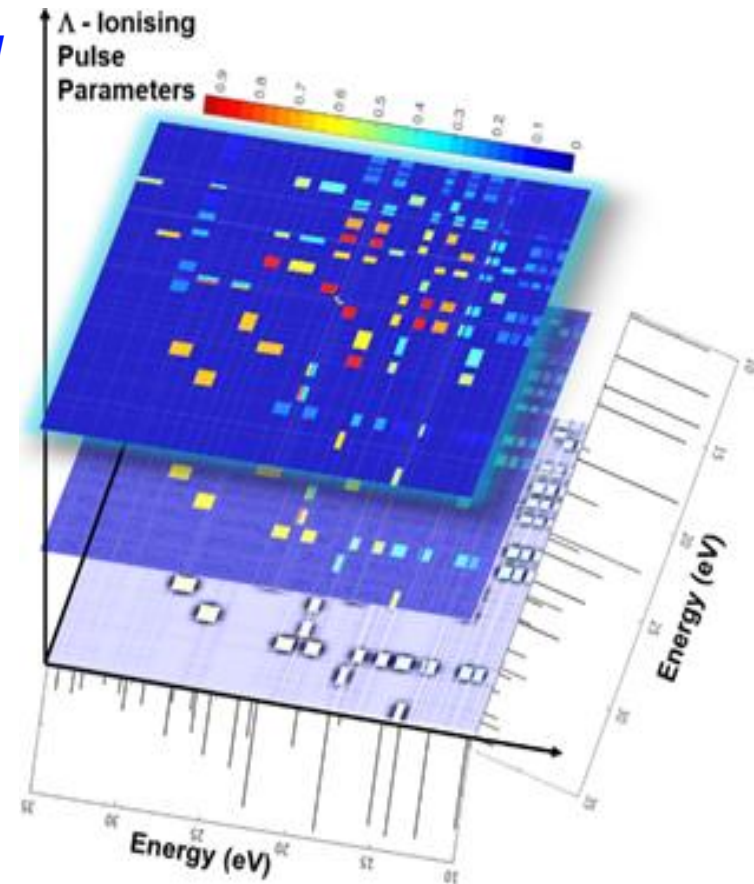
Research Vision

Quantum characterization of the X-ray ionized state in different light-matter interaction regimes

Research goals:

- ❖ *Physical mechanisms of coherence formation upon X-ray attosecond ionization*
- ❖ *Coherent many-electron dynamics on time scales down to the attosecond regime.*
- ❖ *Competition between nuclear motion and Auger decay.*

Laser parameters: frequency, intensity, polarization, stochasticity, etc...



Collaboration with the **strong-field AMO Science Group** of the Chemical Science Division at PULSE Institute (Prof. Philip Bucksbaum)

Research Vision

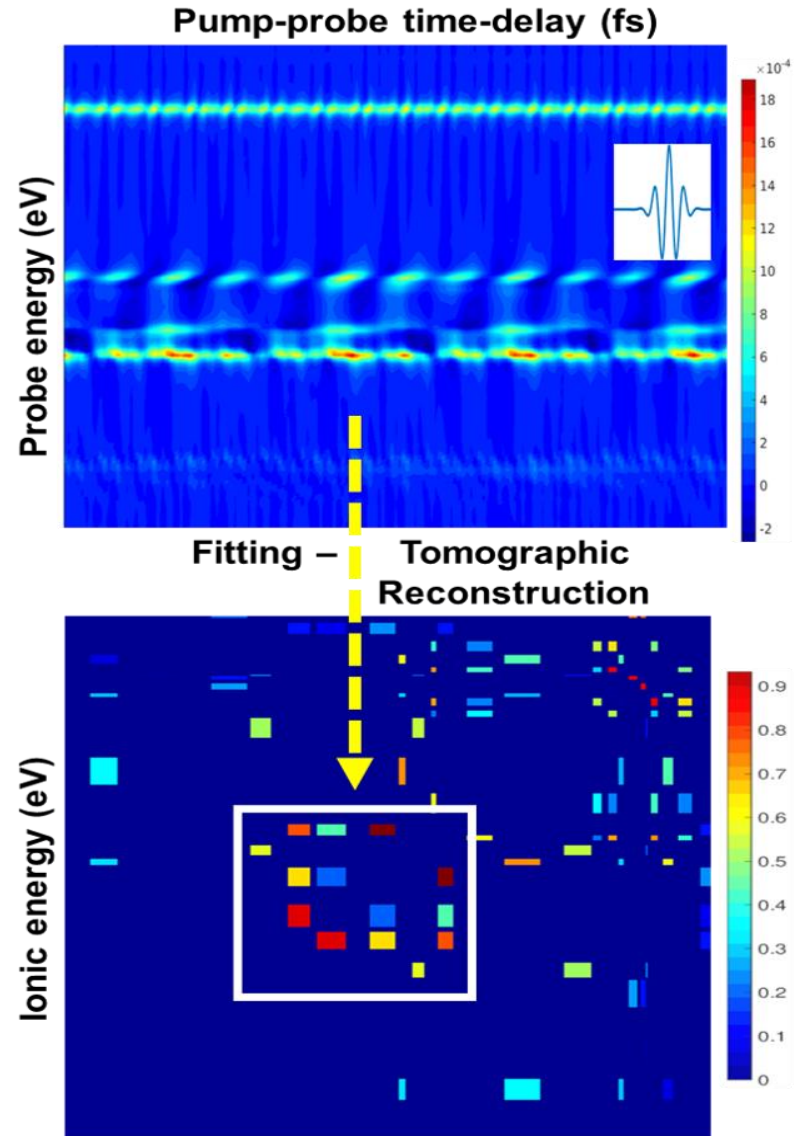
*Quantum state tomography of mixed ionic state prepared by attosecond photoionization **beyond the qubit case***

Research goal:

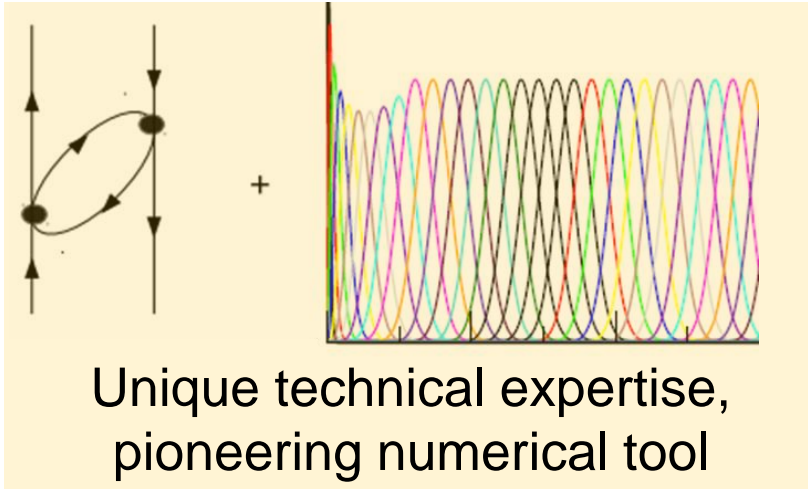
- ❖ *Develop theoretically quantum protocols for the **tomographic reconstruction of the attosecond ionized quantum states***

Extend spectroscopic technique to

- ❖ Time-Resolved Photoelectron Spectroscopy
- ❖ Auger Spectroscopy
- ❖ Stimulated X-ray Raman
- ❖ Multidimensional X-ray Spectroscopy



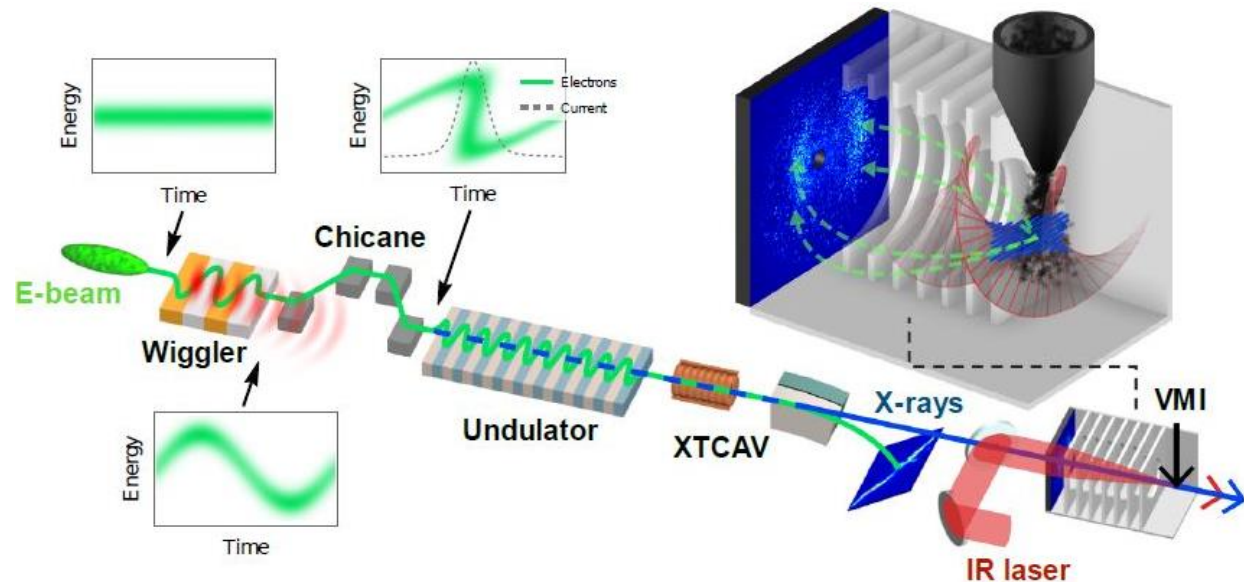
Research Vision



Collaboration with the SLAC **Attosecond Science group** of the Chemical Science Division at PULSE Institute:
(James Cryan, Ago Marinelli, Peter Walter)

Guide experimental efforts at LCLS towards the observation and characterization of new attosecond physical phenomena

“Real-time Observation of Ultrafast Electron Motion using Attosecond XFEL Pulses” **Attosecond Campaign** at LCLS II using **XLEAP**

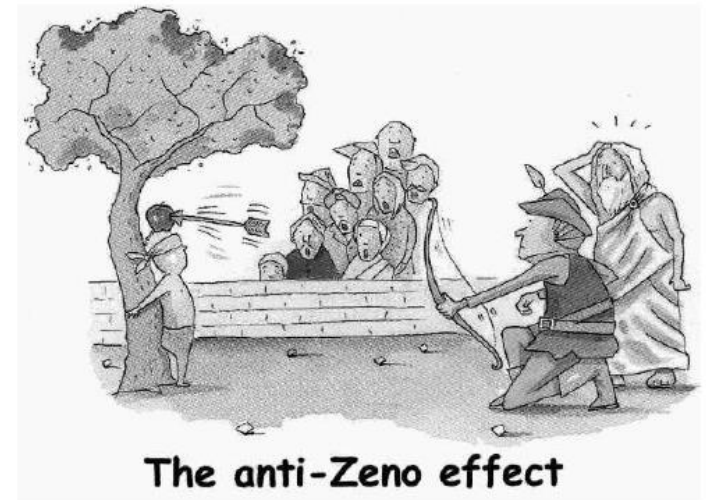
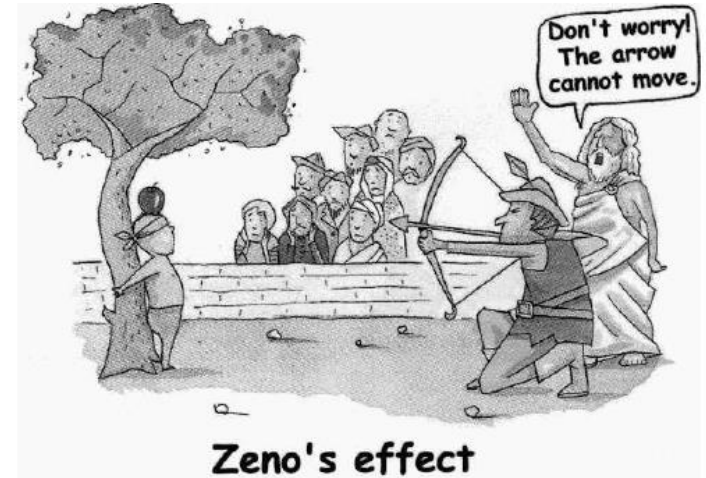
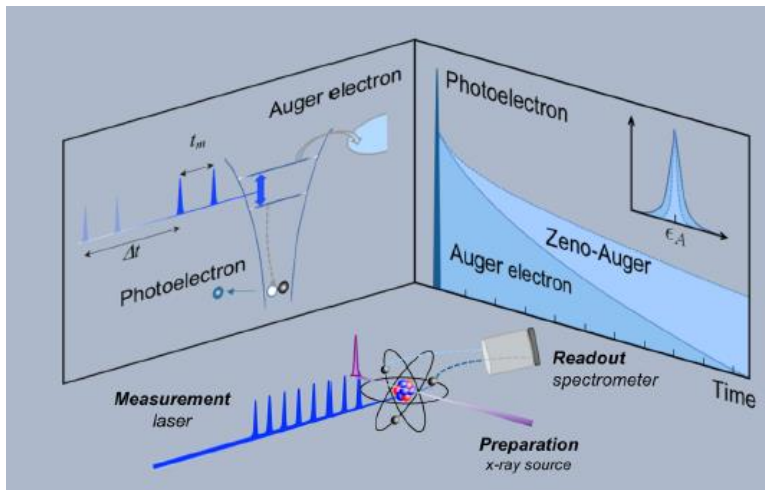


Research Vision

Quantum control of X-ray triggered ultrafast electron-nuclear dynamics

Research goals:

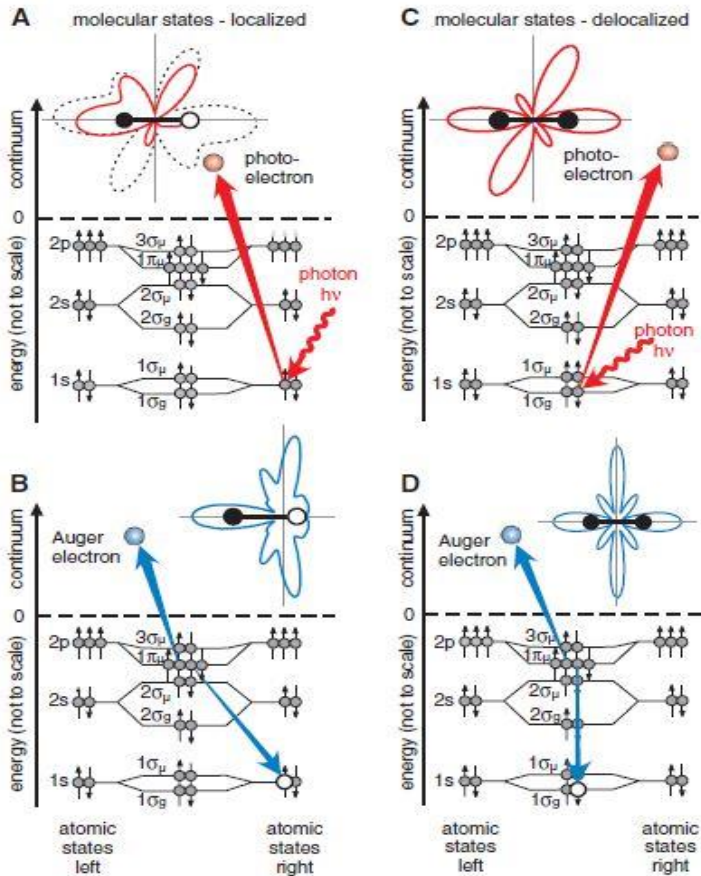
- ❖ **Coherent quantum optimal control schemes** for the manipulation of the electronic coherence and attosecond hole migration
- ❖ **Quantum-Zeno based control schemes for Auger-decay** using trains of attosecond pulses



Investigating electronic entanglement via coincidence measurements

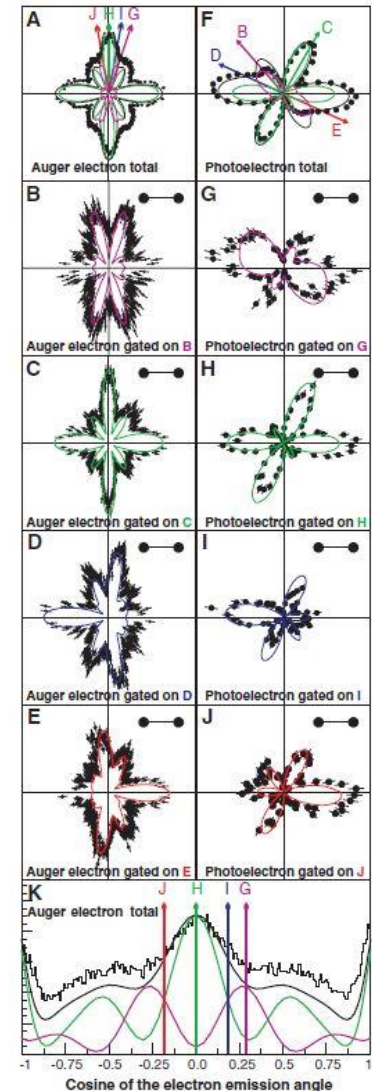
Pioneering experiment [Schöffler *et al.*, *Ultrafast Probing of Core Hole Localization in N₂*, *Science* **320**, 920 (2008)]:

Localization of the core vacancy upon N₂ X-ray photoionization



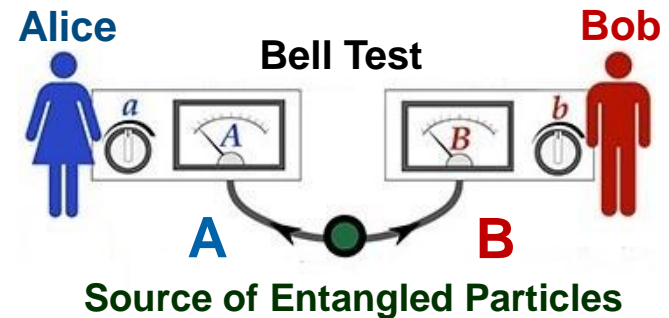
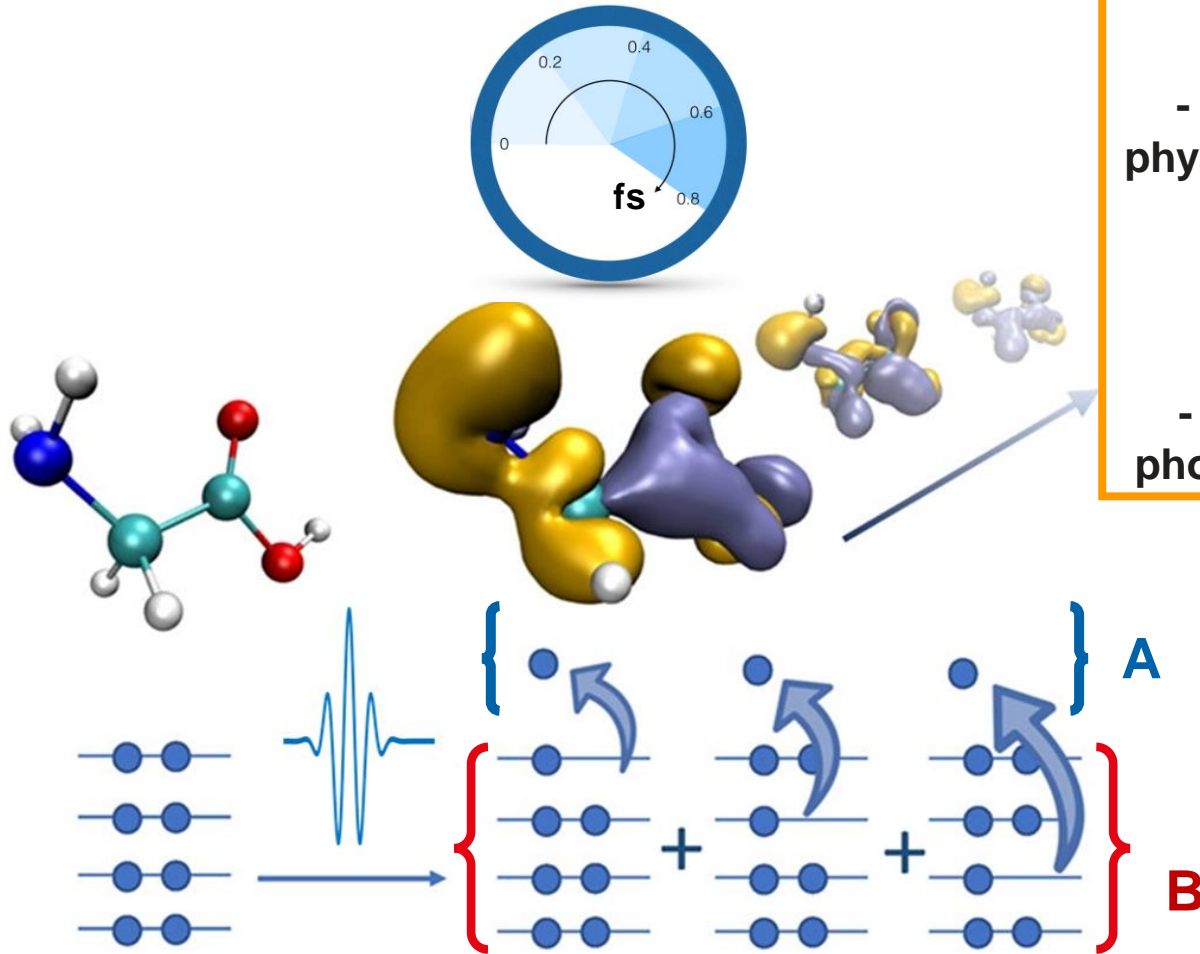
Auger electron angular emission patterns after inner-shell ionization as an ultrafast probe of hole localization:

“Observation of symmetry breaking (localization) or preservation (delocalization) depends on how the quantum entangled Bell state created by Auger decay is detected by the measurement.”



Research Vision

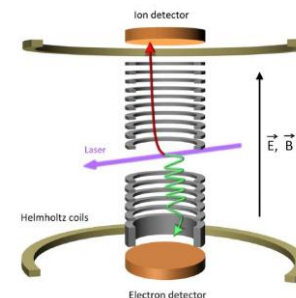
- Quantum electronic coherence & entanglement in molecular photo-ionisation
- Attosecond quantum control of physical & chemical change of matter
- New attosecond quantum phenomena
- Quantum information theory of photo-chemistry & radiation damage



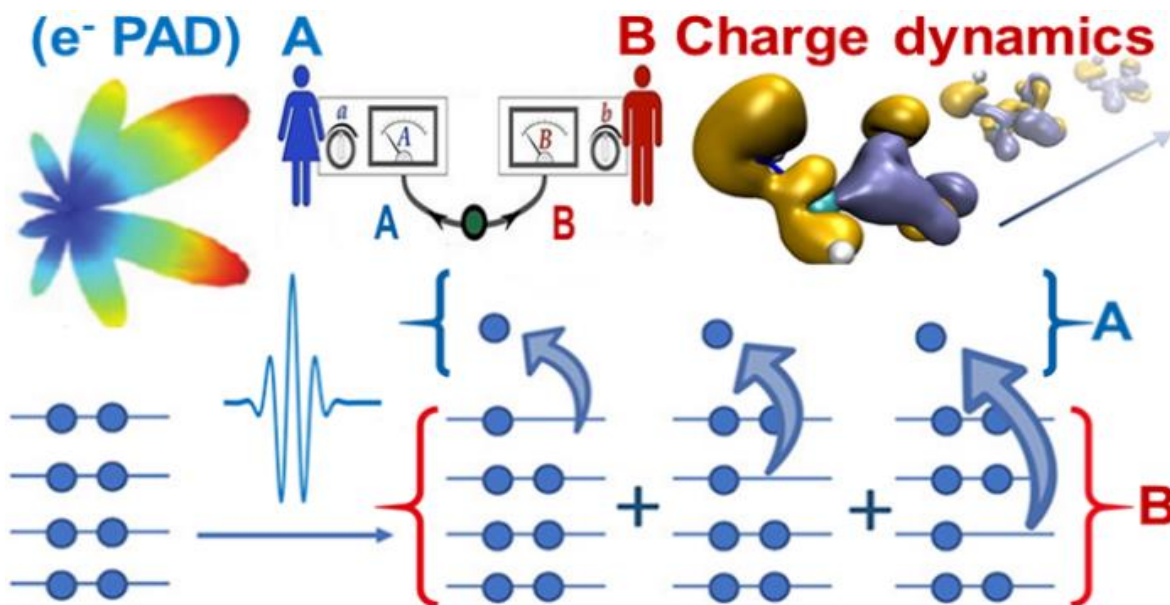
Research Vision

Research goals:

- ❖ **Unravel the role of entanglement in ultrafast charge dynamics**
- ❖ **Design Bell tests and entanglement witnesses** in attosecond photoionization by **coincidence techniques**



High (>100 kHz) rep. rate attosecond experiment with controlled time-delay and COLTRIMS detection.



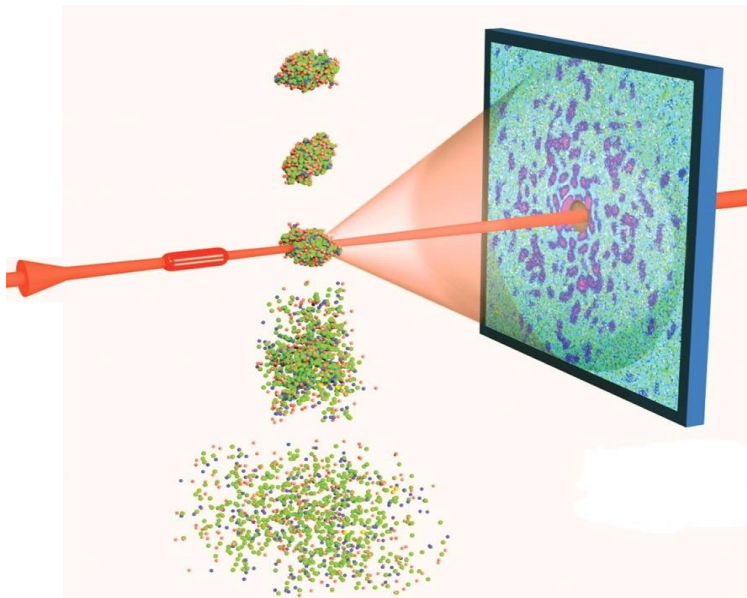
Measuring entanglement between photoelectrons and ions: measure in coincidence photoelectrons & photoproducts from the partner ion fragments (by further ionising the system with a second pulse), i.e. examine in coincidence the two consecutively emitted photoelectrons.

New physics for the DREAM end-station at LCLS-II

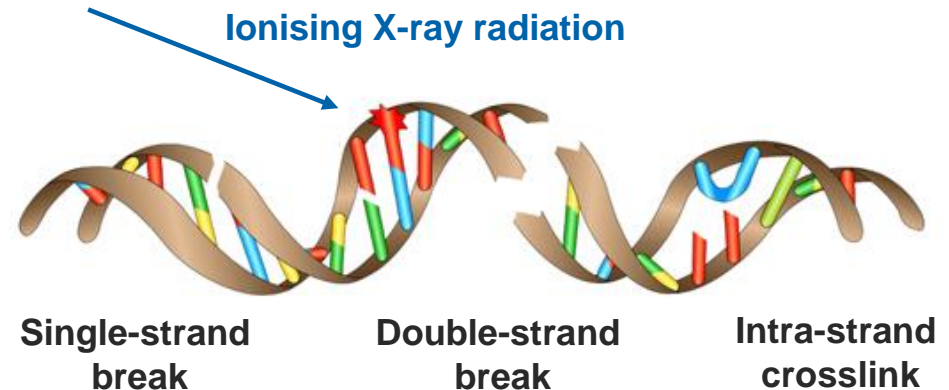
- Better statistics - higher rep rate with LCLS II and higher collection efficiency of electrons and ions with DREAM
- Increased temporal resolution (XLEAP)
- Probing several atomic sites to track charge migration using different probe energies using two-color XLEAP
- Suppression of unwanted background by coincidence measurement allows higher fidelity measurement
- Entanglement measurements correlating the emitted electrons with DREAM.

Research vision: longer-term impacts

On enabling
drug discovery
by single-biomolecule
X-ray “diffract before destroy”
measurement



On more efficient
radiotherapy techniques
for cancer treatment
by understanding & control of
radiation damage to DNA



Research Vision: summary

With my fundamental quantum attoscience program, I will:

- Equip SLAC with the world leading core capability in the many-body attosecond quantum theory
- Harness the power of LCLS-II, DREAM, XLEAP to obtain seminal quantum information insights into X-ray photoionization

Thank you
for your attention

Quantum State Tomography for
Attosecond Ionisation of Entangled
Many-Electron Systems

Dr. Marco Ruberti

Imperial College
London