

# Higgs Precision Physics – Decays

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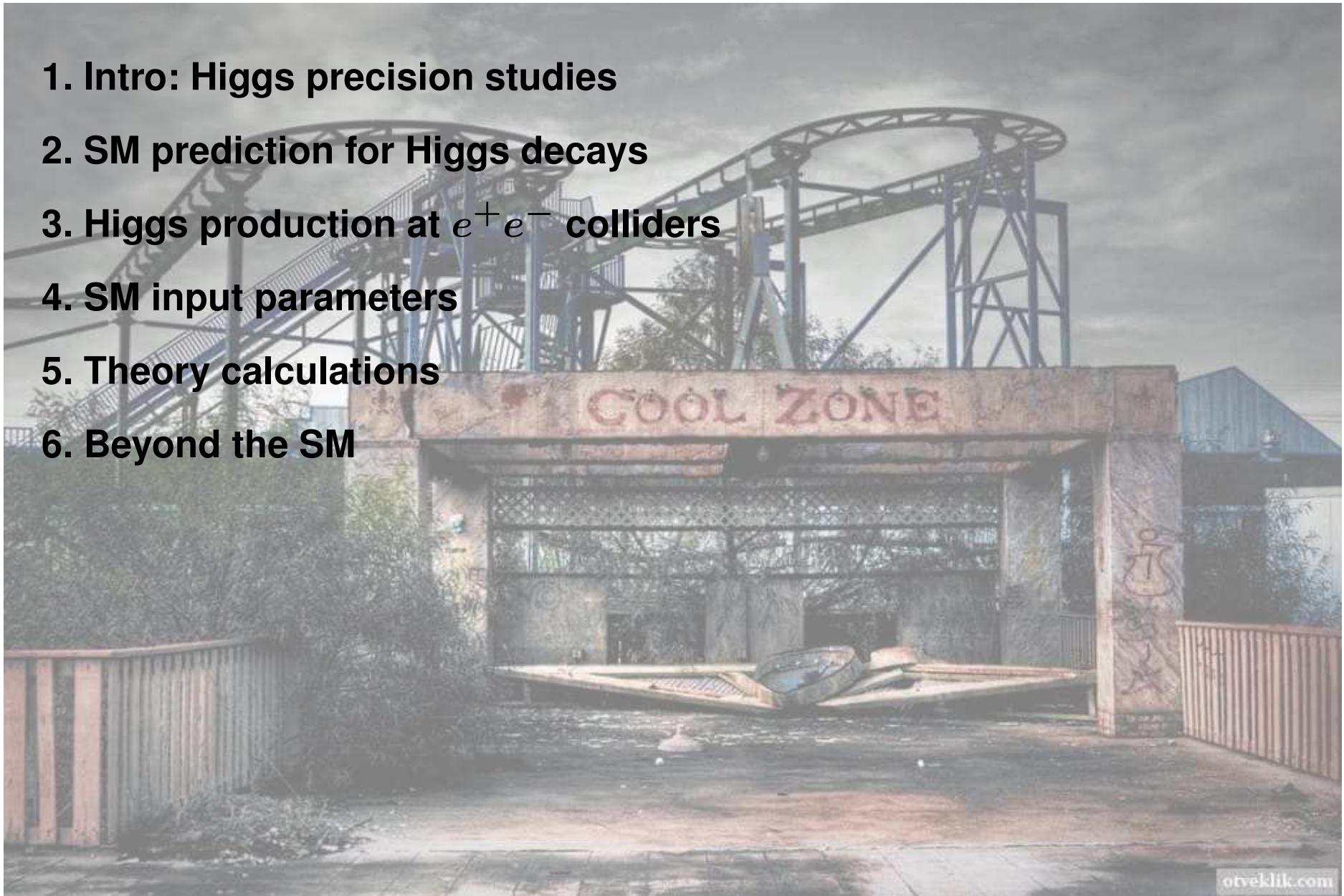


# What does decay have to do with precision?



# What does decay have to do with precision?

1. Intro: Higgs precision studies
2. SM prediction for Higgs decays
3. Higgs production at  $e^+e^-$  colliders
4. SM input parameters
5. Theory calculations
6. Beyond the SM



Higgs physics is entering the precision era:

Determination of Higgs couplings from Higgs production and decay:

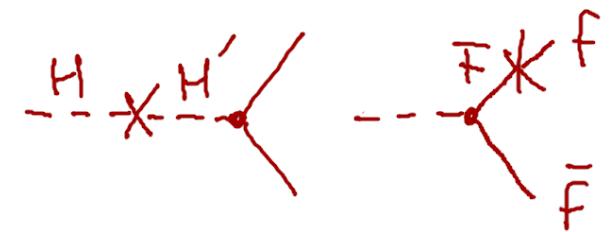
	LHC* (2020)	HL-LHC*	CEPC	FCC-ee
$hbb$	18%	3.7%	1.0%	0.4%
$hcc$	–	–	1.9%	0.7%
$h\tau\tau$	15%	1.9%	1.2%	0.5%
$h\mu\mu$	–	4.3%	5%	6%
$hWW$	9%	1.7%	1.1%	0.2%
$hZZ$	8%	1.5%	0.25%	0.15%
$h\gamma\gamma$	9%	1.8%	1.6%	1.5%
$hgg$	10%	2.5%	1.2%	0.8%

CERN-LPCC-2018-04, PDG '20

\* assuming no exotic decays

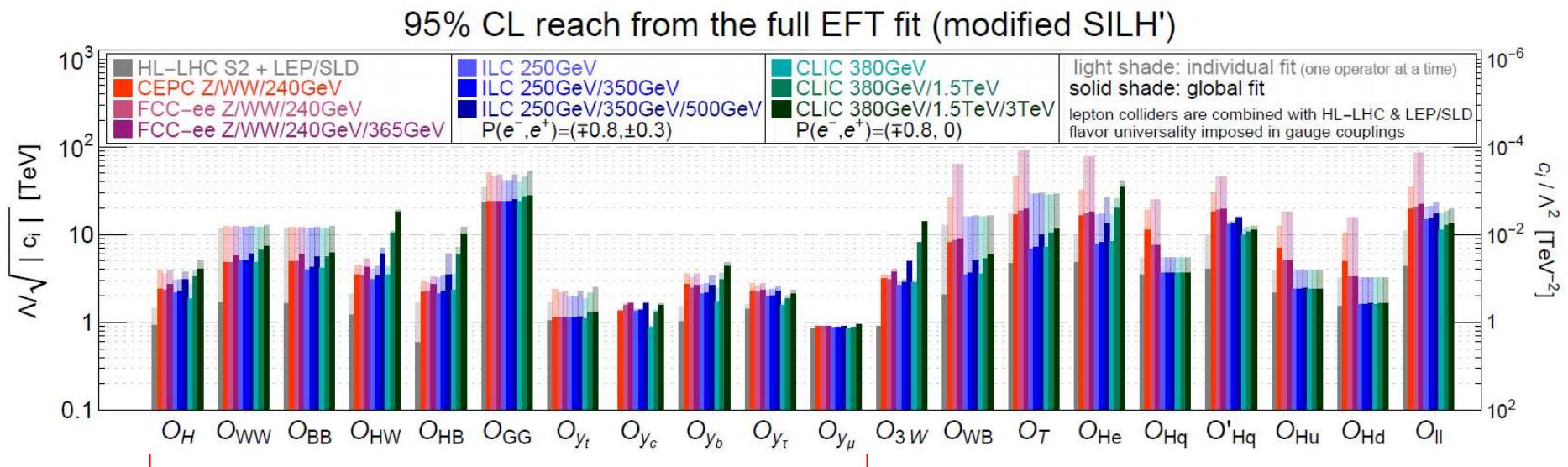
Precision  $\equiv$  large mass scales

- Example: mixing with a heavy particle,  $\theta \sim \frac{m_{\text{SM}}}{M}$



- Model-independent formulation through effective field theory (EFT)

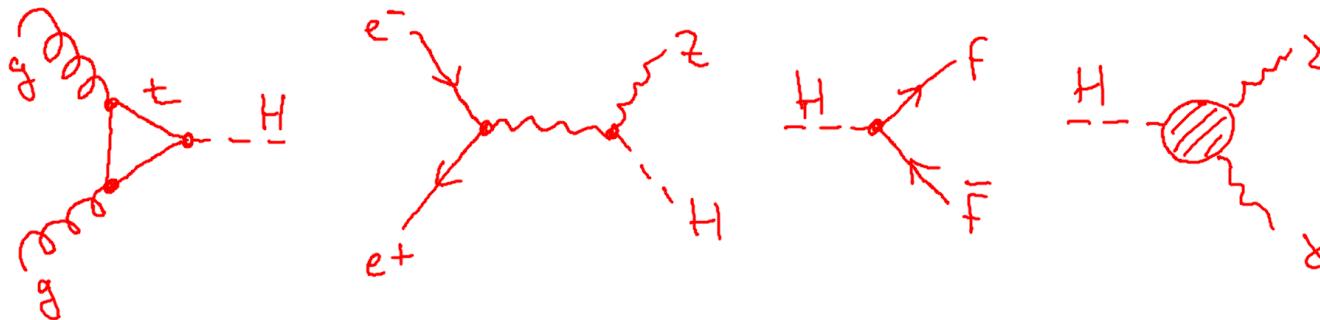
$\rightarrow$  lecture by V. Sanz



de Blas, Durieux, Grojean, Gu, Paul '19

- Access to Higgs interactions through production and decay

→ S Gori, B. Mistlberger



- Most couplings only accessible in decays ( $b, c, \tau, \mu, \gamma, Z$ )
- Decay ratios can reduce systematics, e.g.  $\frac{\sigma[pp \rightarrow H \rightarrow \gamma\gamma]}{\sigma[pp \rightarrow H \rightarrow ZZ]}$
- $\mathcal{O}(\%)$  BSM effects → need SM predictions with higher-order corrections

# SM predictions for Higgs decays

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Reviews: 1404.0319, 1906.05379

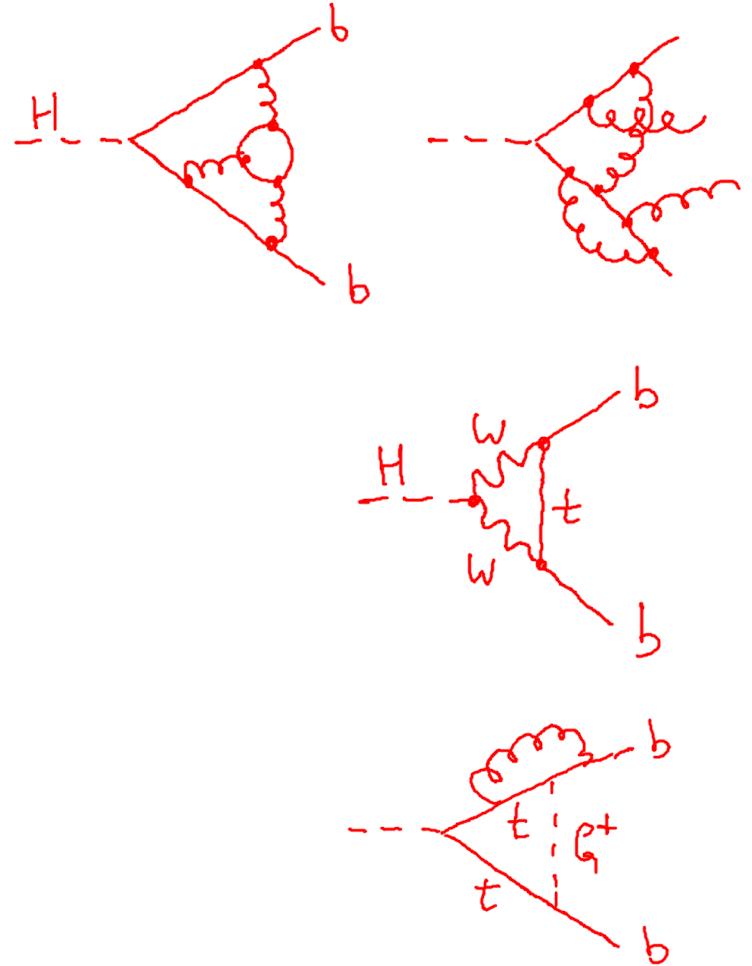
hbb: [HL-LHC: 7%, FCC-ee: 0.8%]

- $\mathcal{O}(\alpha_s^4)$  QCD corrections  
 $21\% + 4\% + 0.2\% - 0.15\%$   
Baikov, Chetyrkin, Kühn '05
- $\mathcal{O}(\alpha)$  QED+EW,  $< 0.5\%$   
Dabelstein, Hollik '92; Kniehl '92
- leading  $\mathcal{O}(\alpha^2)$  and  $\mathcal{O}(\alpha\alpha_s)$  for large  $m_t$   
 $< 0.5\%$  (use for error estimate)  
Kwiatkowski, Steinhauser '94  
Butenschoen, Fugel, Kniehl '07

Current theory error:  $\Delta_{\text{th}} < 0.4\%$

With full 2-loop:  $\Delta_{\text{th}} \sim 0.2\%$

h $\tau\tau$ : [HL-LHC: 3.8%, FCC-ee: 1.1%] → Similar, but no QCD



$hWW^*/hZZ^*$ : [HL-LHC: 3.4%, FCC-ee: 0.4%]

- $m_H < 2m_W \rightarrow$  one or both W/Z must be off-shell

- Need full process  $h \rightarrow 4f$

- Differential distributions contain more info

→ S. Gori

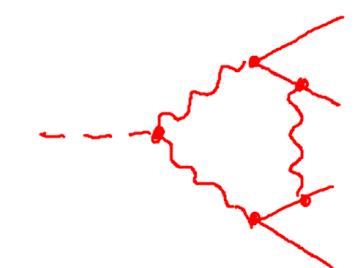
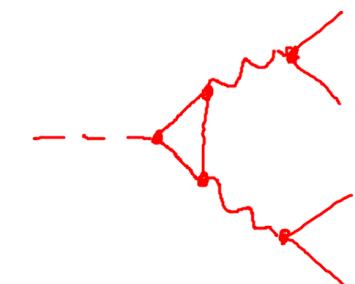
- complete  $\mathcal{O}(\alpha) + \mathcal{O}(\alpha_s)$  for  $h \rightarrow 4f$   
→ MC implementation for distributions

Bredenstein, Denner, Dittmaier, Weber '06

- leading  $\mathcal{O}(\alpha^2)$ ,  $\mathcal{O}(\alpha\alpha_s)$  and  $\mathcal{O}(\alpha\alpha_s^2)$  for large  $m_t$   
→ Small (0.2%) effect

Kniehl, Spira '95; Kniehl, Steinhauser '95

Djouadi, Gambino, Kniehl '97; Kniehl, Veretin '12



Theory error:  $\Delta_{\text{th,EW}} < 0.3\%$ ,  $\Delta_{\text{th,QCD}} < 0.5\%$

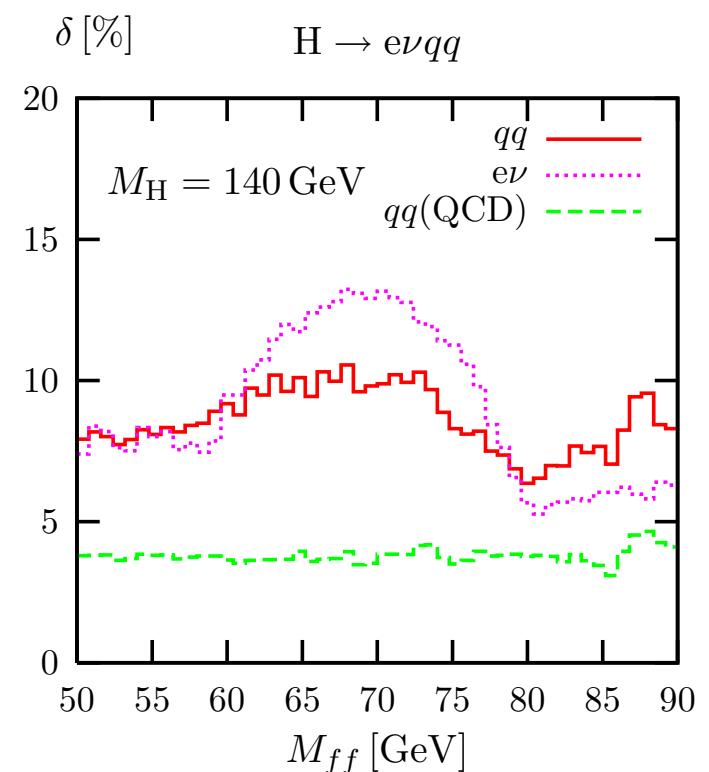
With NNLO final-state QCD corrections:  $\Delta_{\text{th,QCD}} < 0.1\%$

## $hWW^*/hZZ^*$ :

- Non-trivial effects in distributions

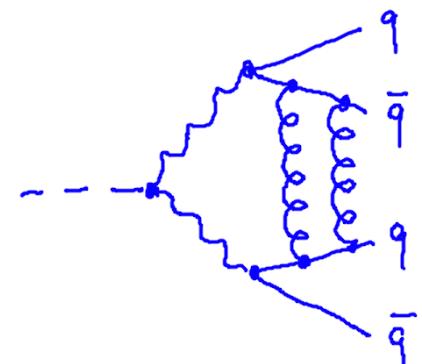
Bredenstein, Denner, Dittmaier, Weber '06

→ Larger theory uncertainty?



- Non-perturbative color reconnection effects

→ No impact on decay rate, but on distributions  
→ Difficult to model, currently done with  
MC programs such as PYTHIA/HERWIG



## $hWW^*/hZZ^*$ :

- Intermediate  $W/Z$  can become on-shell

→ regularize divergencies in  $\frac{1}{p^2 - m_V^2}$  ( $V = W, Z$ )

- One option: *complex mass scheme*

$m_V^2 \rightarrow \mu_V^2 \equiv m_V^2 - i m_V \Gamma'_V$  everywhere,

incl. relations like  $g = \frac{e}{s_W} = \frac{e}{\sqrt{1-m_W^2/m_Z^2}}$

Denner, Dittmaier, Roth, Wieders '05

→ preserves gauge invariance, consistent to all orders

$$\frac{d\Gamma}{d(p^2)} \propto \frac{1}{(p^2 - m_V^2)^2 + m_V^2 \Gamma_V'^2}$$

- In experimental analyses:

$$\frac{d\Gamma}{d(p^2)} \propto \frac{1}{(p^2 - M_V^2)^2 + p^4 M_V^2 / \Gamma_V^2}$$

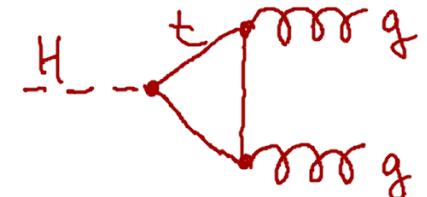
$$m_Z = M_Z / \sqrt{1 + \Gamma_Z^2 / M_Z^2} \approx M_Z - 34 \text{ MeV}$$

$$\Gamma'_Z = \Gamma_Z / \sqrt{1 + \Gamma_Z^2 / M_Z^2} \approx \Gamma_Z - 0.9 \text{ MeV}$$

hgg: [HL-LHC: 5%, FCC-ee: 1.6%]

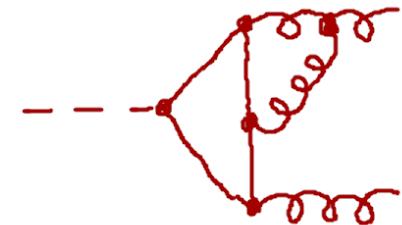
- Loop-induced process in the SM

$$\begin{aligned}\Gamma[h \rightarrow gg] &= \frac{y_t^2 \alpha_s^2 m_t^2}{4\pi^3 m_H} \left[ 1 + \left( 1 - \frac{4m_t^2}{m_H^2} \right) \arcsin^2 \frac{m_H}{2m_t} \right]^2 \\ &\approx \frac{\alpha_s^2 m_H^3}{72\pi^3 v^2} \quad \left( m_H \ll 2m_t, \ y_t = \frac{\sqrt{2}}{v} m_t \right)\end{aligned}$$



For heavy BSM quark:

$$\delta\Gamma[h \rightarrow gg] \approx \frac{y_{hQQ}^2 \alpha_s^2 m_H^3}{144\pi^3 M_Q^2}$$



- $N^n$ LO corrections require  $(n+1)$ -loop diagrams

- $\mathcal{O}(\alpha_s)$  exact result

Djouadi, Graudenz, Spira, Zerwas '95

- Use large  $m_t$  expansion, used for  $\mathcal{O}(\alpha_s^2)$

Schreck, Steinhauser '07

→ fast convergence ( $m_t^{-4}$  term is  $\sim 1\%$  of  $m_t^0$  term)

- $\mathcal{O}(\alpha_s^3)$  with low-energy theorem (equiv. to 1st term in large  $m_t$  exp.)

Baikov, Chetyrkin '06

hgg:

- **Low-energy theorem:** for  $m_t \rightarrow \infty$  top loop becomes effective  $hgg$  interaction:

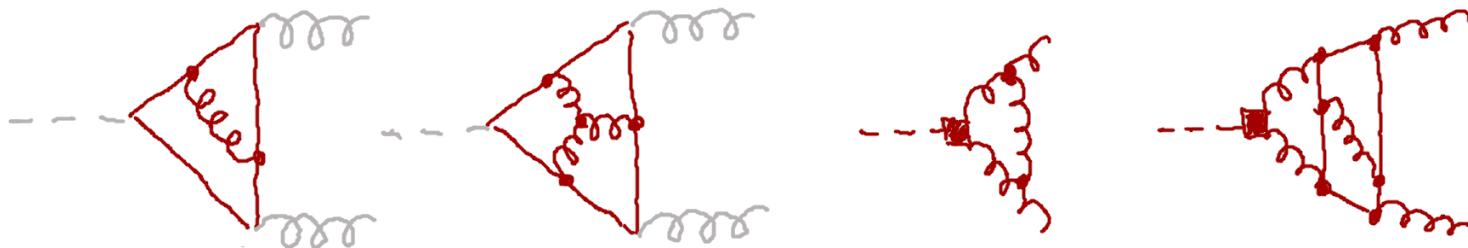
$$\mathcal{L}_{\text{eff}} = \frac{C_{hg}}{v} h G^{\mu\nu} G_{\mu\nu} \quad \text{LO SM res.: } C_{hg} = -\frac{\alpha_s}{12\pi} \quad (1)$$

$$\Gamma[h \rightarrow gg] = \frac{2C_{hg}^2 m_H^3}{\pi v^2} \quad (2)$$

Gauge-invariant form:

$$\mathcal{L}_{\text{eff}} = \frac{C_{hg}}{v^2} (\phi^\dagger \phi) G^{\mu\nu} G_{\mu\nu}, \quad \phi = \left( G^+, \frac{v+h+G^0}{\sqrt{2}} \right)^\top \quad (3)$$

- Higher orders:
  - corrections to  $C_{hg}$  (zero external momentum)
  - corrections to (2) with effective  $hgg$  cpl.

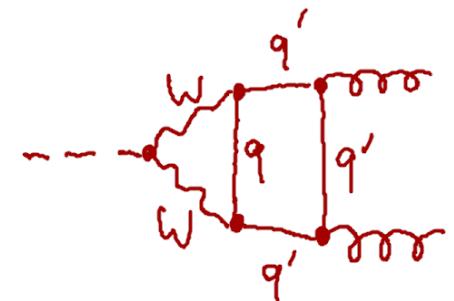
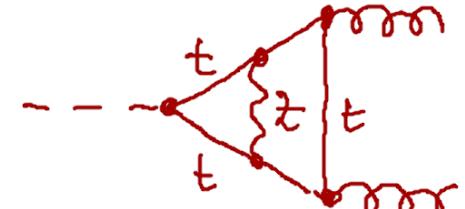


hgg:

■ **EW corrections:**

- Can use expansion or low-energy theorem for top-quark contributions

Degrassi, Maltoni '04



- Direct 2-loop calculation for (dominant) light-quark contributions

Aglietti, Bonciani, Degrassi, Vicini '04; Degrassi, Maltoni '04

QCD corrections: 65% + 20% + 2%

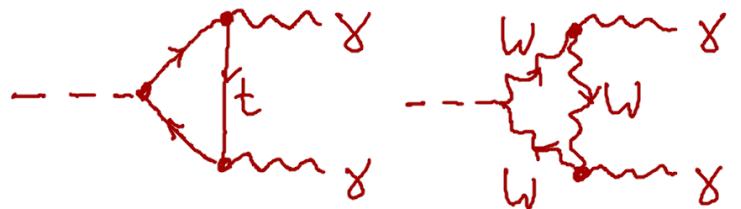
EW corrections:  $\sim 5\%$

Theory error (dominated by QCD):  $\Delta_{\text{th}} \approx 3\%$

With  $\mathcal{O}(\alpha_s^4)$  in large  $m_t$ -limit (4-loop massless QCD diags.):  $\Delta_{\text{th}} \approx 1\%$

$h\gamma\gamma$ : [HL-LHC: 3.6%, FCC-ee: 3.0%]

- Top+W loop at LO

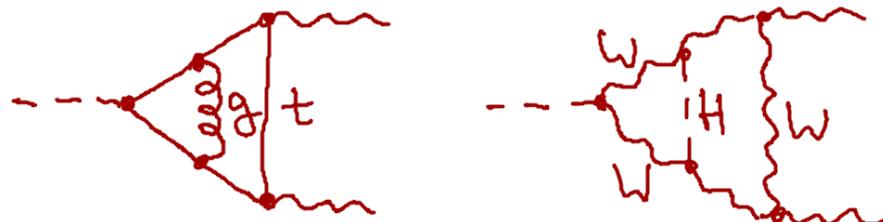


- $\mathcal{O}(\alpha_s^2)$  QCD corrections (with large  $m_t$  expansion)

Zheng, Wu '90; Djouadi, Spira, v.d.Bij, Zerwas '91  
 Dawson, Kauffman '93; Maierhöfer, Marquard '12

- NLO EW

Aglietti, Bonciani, Degrassi, Vicini '04; Degrassi, Maltoni '04  
 Actis, Passarino, Sturm, Uccirati '08



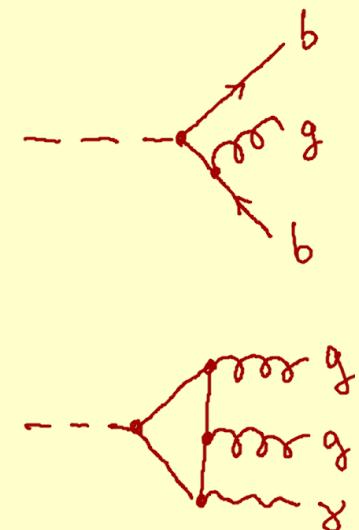
Theory error:  $\Delta_{\text{th}} < 1\%$

- $\Gamma \rightarrow XX$  is not a real observable
- Effect of detector acceptance and isolation criteria, in particular for extra gluon (jet) and photon radiation
- Effect of selection cuts to reduce backgrounds
- Monte-Carlo tools needed to simulate these effects

■ **Current LHC studies:**  
rad. corrections in decay, only parton shower for  $\gamma, g$  emission;  
differential rates corrected (reweighted) with fixed-order results

■ **More accurate procedure:**  
Implement fixed-order rad. corr. in MC program,  
**matching** to avoid double counting with parton shower

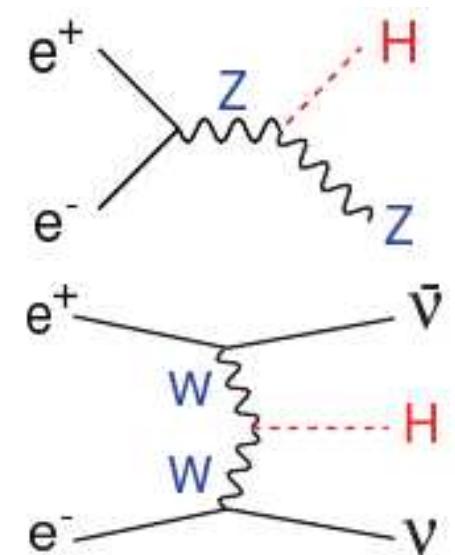
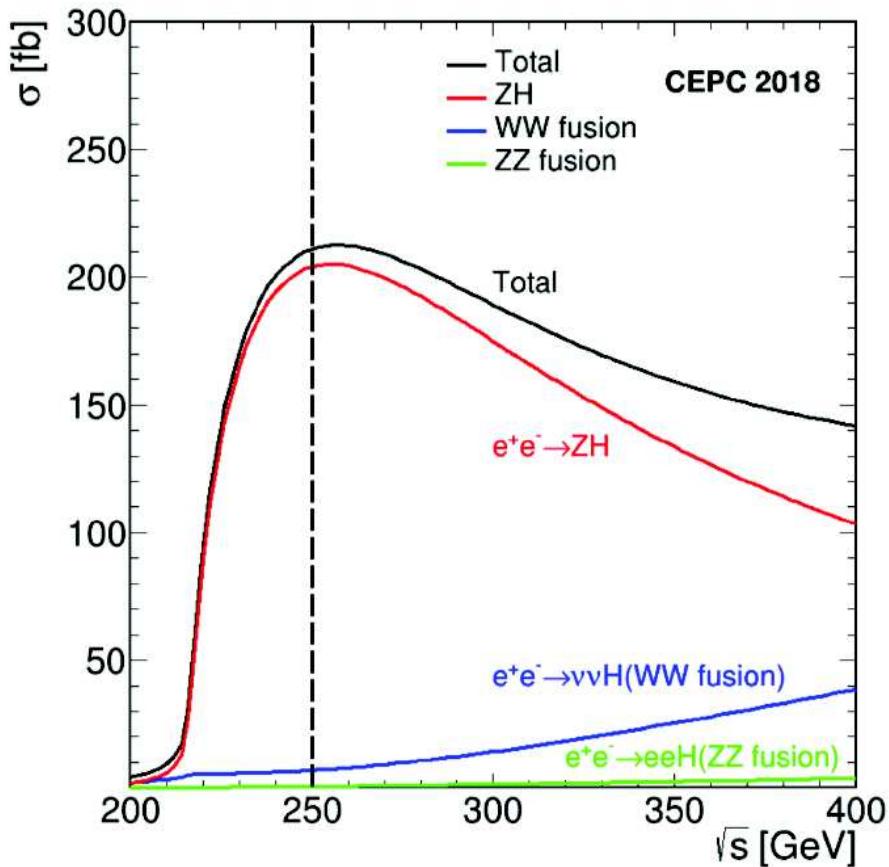
Catani, Krauss, Kuhn, Webber '01  
Nason '04; Frixione, Nason, Oleari '07



# Higgs production at $e^+e^-$ colliders

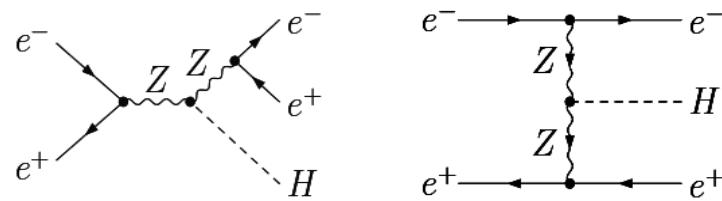
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- hZ production: dominant at  $\sqrt{s} \sim 240$  GeV
- WW fusion: sub-dominant but useful for constraining  $h$  width Han, Liu, Sayre '13



hZ production: [CEPC: 0.5%, FCC-ee: 0.3%]

- $\mathcal{O}(\alpha)$  corr. to  $hZ$  production and  $h, Z$  decay      Kniehl '92; Denner, Küblbeck, Mertig, Böhm '92  
Consoli, Lo Presti, Maiani '83; Jegerlehner '86  
Akhundov, Bardin, Riemann '86
- $\Gamma_H/m_H \approx 4 \times 10^{-5}$ ,  $\Gamma_Z/m_Z \approx 0.025$      $\Rightarrow$  include off-shell Z effects  
Technology for  $\mathcal{O}(\alpha)$  corr. to  $h f \bar{f}$  production available      Boudjema et al. '04  
Denner, Dittmaier, Roth, Weber '03

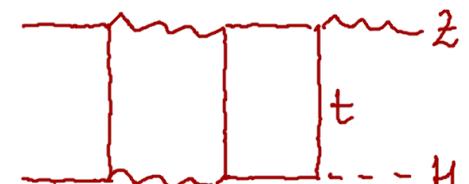
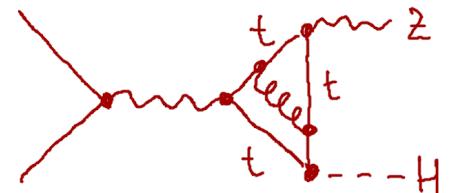


- $\mathcal{O}(\alpha\alpha_s)$  corrections  
Gong et al. '16; Chen, Feng, Jia, Sang '18

Theory error:  $\Delta_{th} \sim \mathcal{O}(1\%)$

With full 2-loop corrections for  $ee \rightarrow HZ$ :

$\Delta_{th} \lesssim \mathcal{O}(0.3\%)$



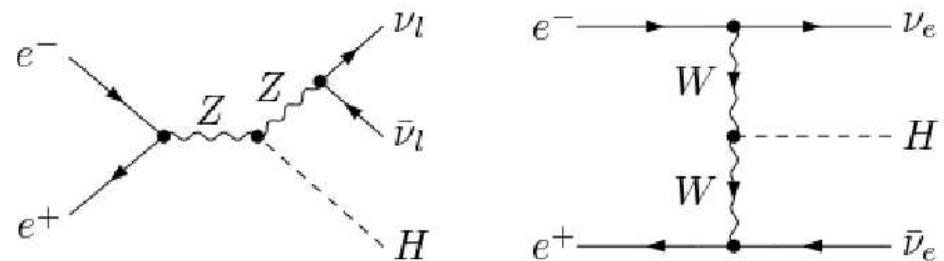
Freitas, Song '21; Li, Wang, Wu '21

WW fusion:

- $\mathcal{O}(\alpha)$  corrections

Belanger et al. '02; Denner, Dittmaier, Roth, Weber '03

Theory error:  $\Delta_{\text{th}} \sim \mathcal{O}(1\%)$ ?



Full NNLO calc. for  $2 \rightarrow 3$  process is very challenging, but may not be needed

SM predictions for Higgs decays need measured input parameters

Reviews: [1906.05379](#), [2012.11642](#)

- $M_Z$ ,  $M_W$ : current precision  $< 0.1\%$  → negligible impact
- $m_t$ : Most precise measurement at LHC:  $\delta m_t \sim 0.3$  GeV

Additional theory error from scheme translation

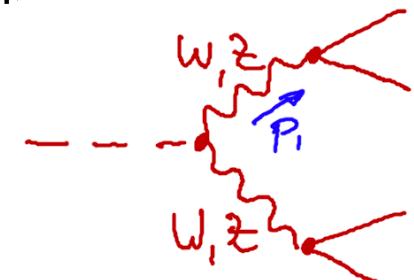
[Hoang, Plätzer, Samitz '18](#)

→ Total uncertainty  $\delta m_t \sim 0.5$  GeV

→ Negligible impact because  $m_t$  only appears in loops

SM predictions for Higgs decays need measured input parameters

Reviews: [1906.05379](#), [2012.11642](#)



- $M_H$ : high precision important for  $h \rightarrow WW^*, ZZ^*$

$$\text{amplitude} \propto \frac{1}{p_1^2 - m_V^2}, \quad p_1 \sim m_H - m_V \text{ (both } V, V^* \text{ at rest)}$$

$$\Gamma_{VV^*} = [\text{energy}]$$

$$\Rightarrow \Gamma_{VV^*} \propto \frac{[\text{energy}]^5}{|p_1^2 - m_V^2|^2} \sim \frac{m_H^5}{[(m_H - m_V)^2 - m_V^2]^4}$$

$$\text{Uncertainty } \delta m_H \quad \Rightarrow \quad \frac{\delta \Gamma_{VV^*}}{\Gamma_{VV^*}} = \frac{m_H - 6m_V}{m_H - 2m_V} \frac{\delta m_H}{m_H} \approx 10 \frac{\delta m_H}{m_H}$$

$$\delta m_H = 0.2 \text{ GeV} \quad \Rightarrow \quad \frac{\delta \Gamma_{VV^*}}{\Gamma_{VV^*}} = 1.6\%$$

$$\text{CEPC/FCC-ee/ILC can achieve } \delta m_H \lesssim 20 \text{ MeV} \quad \Rightarrow \quad \frac{\delta \Gamma_{VV^*}}{\Gamma_{VV^*}} \lesssim 0.2\%$$

- $\alpha_s$ : important for  $h \rightarrow gg$  (also  $h \rightarrow q\bar{q}$ )

$$\delta\alpha_s = 0.001 \quad \Rightarrow \quad \frac{\delta\Gamma_{gg}}{\Gamma_{gg}} \approx 3\%$$

Methods for  $\alpha_s$  determination:

d'Enterria, Skands, et al. '15

- Most precise determination using Lattice QCD:

$$\alpha_s = 0.1184 \pm 0.0006 \quad \text{HPQCD '10}$$

$$\alpha_s = 0.1185 \pm 0.0008 \quad \text{ALPHA '17}$$

$$\alpha_s = 0.1179 \pm 0.0015 \quad \text{Takaura et al. '18}$$

$$\alpha_s = 0.1172 \pm 0.0011 \quad \text{Zafeiropoulos et al. '19}$$

→ Difficulty in evaluating systematics

- $e^+e^-$  event shapes and DIS:  $\alpha_s \sim 0.114$

Alekhin, Blümlein, Moch '12; Abbate et al. '11; Gehrmann et al. '13

→ Subject to sizeable non-perturbative power corrections

→ Systematic uncertainties in power corrections?

- $\alpha_s$ :

- Hadronic  $\tau$  decays:  $\alpha_s = 0.119 \pm 0.002$  PDG '18

→ Non-perturbative uncertainties in OPE and from duality violation

Pich '14; Boito et al. '15,18

- Electroweak precision ( $R_\ell = \Gamma_Z^{\text{had}} / \Gamma_Z^\ell$ ):

$\alpha_s = 0.120 \pm 0.003$  PDG '18

→ No (negligible) non-perturbative QCD effects

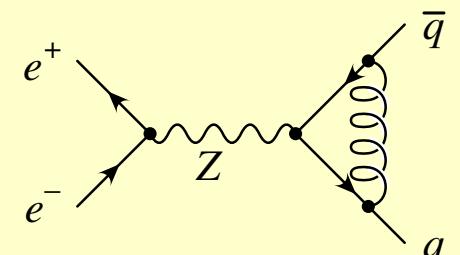
FCC-ee:  $\delta R_\ell \sim 0.001$

$$\Rightarrow \delta \alpha_s < 0.0001$$

Theory input: N<sup>3</sup>LO EW corr. + leading N<sup>4</sup>LO

to keep  $\delta_{\text{th}} R_\ell \lesssim \delta_{\text{exp}} R_\ell$

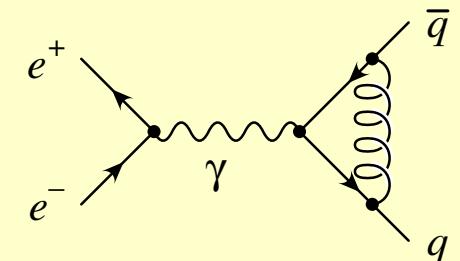
**Caviat:**  $R_\ell$  could be affected by new physics



- $\alpha_s$ :

- $R = \frac{\sigma[ee \rightarrow \text{had.}]}{\sigma[ee \rightarrow \mu\mu]}$  at lower  $\sqrt{s}$

e.g. CLEO ( $\sqrt{s} \sim 9$  GeV):  $\alpha_s = 0.110 \pm 0.015$   
 Kühn, Steinhauser, Teubner '07



- dominated by  $s$ -channel photon, less room for new physics
- QCD still perturbative

naive scaling to  $50 \text{ ab}^{-1}$  (BELLE-II):  $\delta\alpha_s \sim 0.0001$

- $m_b, m_c$ : From quarkonia spectra using Lattice QCD

$\delta m_b^{\overline{\text{MS}} \rightarrow 0} \sim 30 \text{ MeV}, \delta m_c^{\overline{\text{MS}} \rightarrow 0} \sim 25 \text{ MeV}$

LHC HXSWG '16

$$\Rightarrow \frac{\delta\Gamma_{bb}}{\Gamma_{bb}} \approx 1.4\%, \quad \frac{\delta\Gamma_{cc}}{\Gamma_{cc}} \approx 4.0\%$$

→ estimated improvements  $\delta m_b^{\overline{\text{MS}} \rightarrow 0} \sim 13 \text{ MeV}, \delta m_b^{\overline{\text{MS}} \rightarrow 0} \sim 7 \text{ MeV}$

$$\Rightarrow \frac{\delta\Gamma_{bb}}{\Gamma_{bb}} \approx 0.4\%, \quad \frac{\delta\Gamma_{cc}}{\Gamma_{cc}} \approx 0.4\%$$

Lepage, Mackenzie, Peskin '14

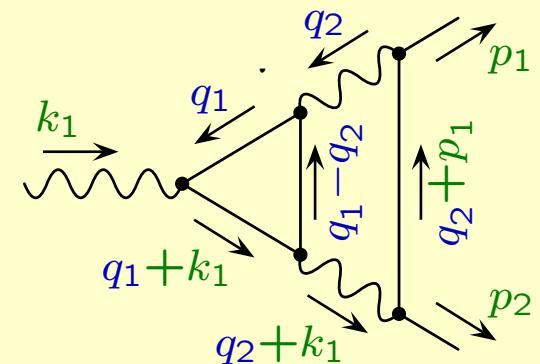
## Theory uncertainties

- Theory error estimate is not well defined, ideally  $\Delta_{\text{th}} \ll \Delta_{\text{exp}}$
- Common methods:
  - Count prefactors ( $\alpha, N_c, N_f, \dots$ )
  - Extrapolation of perturbative series
    - e.g. assume  $\frac{\Gamma_{\text{NNNLO}}}{\Gamma_{\text{NNLO}}} \sim \frac{\Gamma_{\text{NNLO}}}{\Gamma_{\text{NLO}}}$
  - Renormalization scale dependence
    - (for  $\overline{\text{MS}}$  renormalization, widely used for QCD)
  - Renormalization scheme dependence
    - e.g. compare  $\overline{\text{MS}}$  and OS renormalization

Experimental precision requires inclusion of **multi-loop corrections** in theory

Integrals over loop momenta:

$$\int d^4 q_1 d^4 q_2 f(q_1, q_2, p_1, k_1, \dots, m_1, m_2, \dots)$$



Computer algebra tools:

- Generation of diagrams,  $\mathcal{O}(1000) - \mathcal{O}(10000)$
  - Lorentz and Dirac algebra
  - Integral simplification (e.g. symmetries)
- } not a limiting factor

Evaluation of loop integrals:

- Analytical
- Approximate (expansions)
- Numerical

- Mostly used for diagrams with few mass scales
- Reduce to **master integrals** with integration-by-parts and other identities  
[Chetyrkin, Tkachov '81](#); [Gehrmann, Remiddi '00](#); [Laporta '00](#); ...

Public programs:	Reduze	<a href="#">von Manteuffel, Studerus '12</a>
	FIRE	<a href="#">Smirnov '13,14</a>
	LiteRed	<a href="#">Lee '13</a>
	KIRA	<a href="#">Maierhoefer, Usovitsch, Uwer '17</a>

- Large need for computing time and memory
- Evaluate master integrals with differential equations or Mellin-Barnes rep.  
[Kotikov '91](#); [Remiddi '97](#); [Smirnov '00,01](#); [Henn '13](#); ...
  - Result in terms of Goncharov polylogs / multiple polylogs
  - Some problems need iterated elliptic integrals / elliptic multiple polylogs  
[Broedel, Duhr, Dulat, Trancredi '17,18](#)  
[Ablinger et al. '17](#)
  - Even more classes of functions needed in future?

# Asymptotic expansions

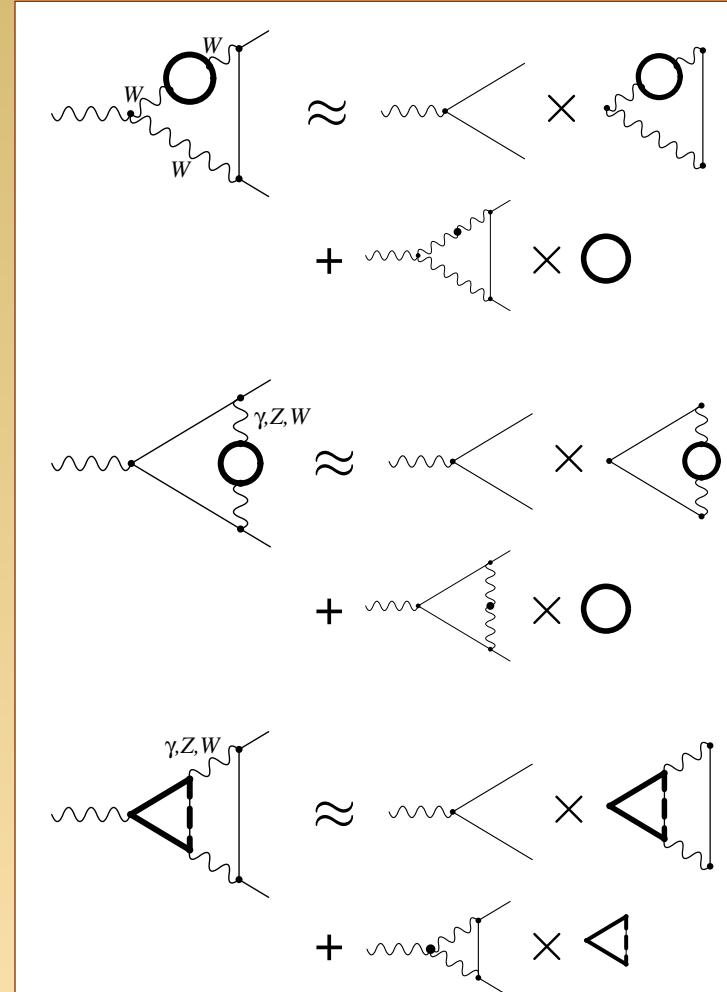
- Exploit large mass/momentum ratios,  
e. g.  $M_Z^2/m_t^2 \approx 1/4$
- Evaluate coeff. integrals analytically
- Fast numerical evaluation

→ Public programs:

exp      [Harlander, Seidensticker, Steinhauser '97](#)  
asy      [Pak, Smirnov '10](#)

→ Possible limitations:

- no appropriate mass/momentum ratios
- bad convergence
- impractical if too many mass/moment scales



## Challenge 1: presence of UV/IR divergencies

- Remove through subtraction terms

$$\underbrace{\int d^4 q_1 d^4 q_2 (f - f_{\text{sub}})}_{\text{finite}} + \underbrace{\int d^4 q_1 d^4 q_2 f_{\text{sub}}}_{\text{solve analytically}}$$

Cvitanovic, Kinoshita '74  
Levine, Park, Roskies '82  
Bauberger '97  
Nagy, Soper '03  
Awramik, Czakon, Freitas '06  
Becker, Reuschle, Weinzierl '10  
Sborlini et al. '16  
...

- Remove through variable transformations:

- a) Sector decomposition

Public programs: (py) SecDec  
FIESTA

Carter, Heinrich '10; Borowka et al. '12,15,17  
Smirnov, Tentyukov '08; Smirnov '13,15

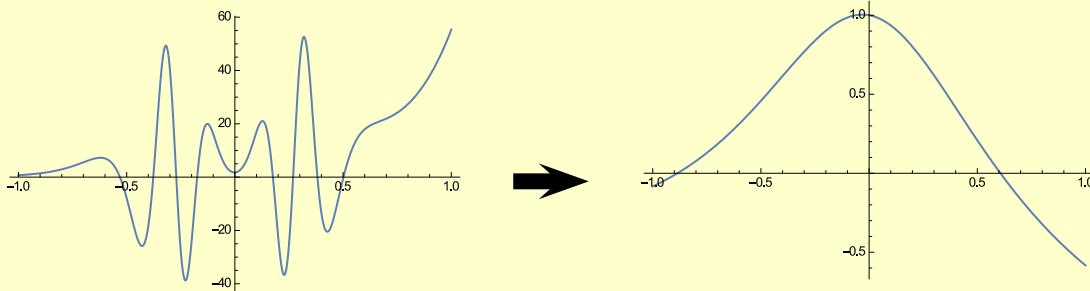
- b) Mellin-Barnes representations

Public programs: MB/MBresolve  
AMBRE/MBnumerics

Czakon '06; Smirnov, Smirnov '09  
Gluza, Kajda, Riemann '07  
Dubovyk, Gluza, Riemann '15  
Usovitsch, Dubovyk, Riemann '18

## Challenge 2: stability and convergence

- Integration in momentum space:  $4L$  dimensions ( $L = \#$  of loops)
- Integration in Feynman parameters:  $P - 1$  dimensions ( $P = \#$  of propagators)
  - Multi-dim. integrals need large computing resources and converge slowly
- Variable transformations to avoid singularities and peaks



Analytical techniques and expansions:

Complexity increases with ...

... more loops;

... more external particles;

**... more different masses**

Numerical techniques:

Complexity increases with ...

... more loops;

... more external particles;

**... fewer masses**

“Kappa framework”

- Multiply Higgs Feynman rules with factor  $\kappa_X \neq 1$ :

$$H - \frac{f}{\bar{f}} = -\frac{igm_f}{2m_W} \kappa_f$$

$$H - \frac{W_\mu}{W_\nu} = igm_W g_{\mu\nu} \kappa_W$$

- breaks gauge-invariance  
→ in general not possible to compute EW corrections

Effective field theory framework

→ V. Sanz

$$\mathcal{L} = \mathcal{L}_{\text{SM}} + \sum_{d=5}^{\infty} \frac{1}{\Lambda^{d-4}} \sum_i c_i \mathcal{O}_i^{(d)}$$

- Operators required to satisfy SM gauge invariance
- Valid description for energies  $E \ll \Lambda$  ( $\Lambda \sim$  mass of heavy particles)
- Leading contribution to Higgs physics:  $d = 6$
- SMEFT:** Higgs doublet as in SM
- HEFT:** Higgs and Goldstone bosons treated independently

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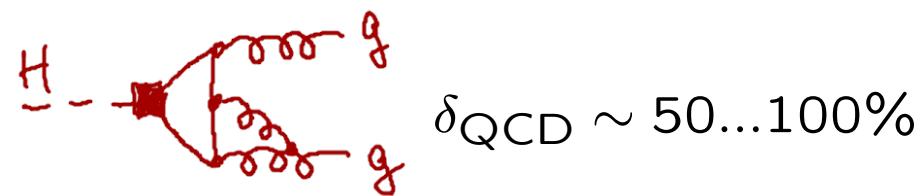
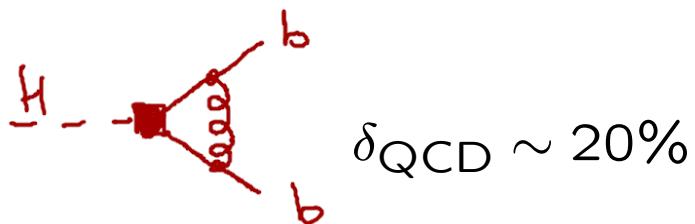
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No conclusive BSM evidence in Higgs physics

- BSM effects are small
- ignore rad. corr. to BSM contributions (?)

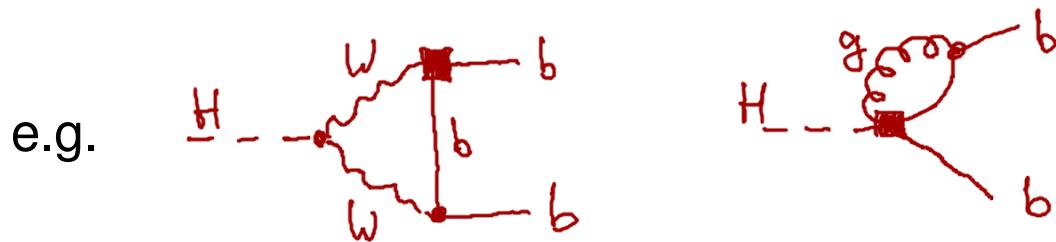
But some higher-order effects can be significant:

- Large SM QCD corrections to LO operator insertions:



- Practical calculations similar to SM

- New operators appearing in loops:

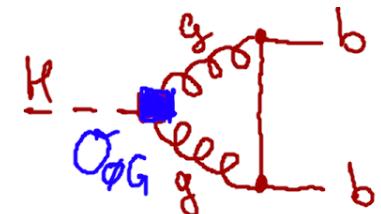


## Example 1: $hgg$ operator in $h \rightarrow b\bar{b}$

$$\frac{\delta\Gamma_{bb}}{\Gamma_{bb}} = C_F \frac{\alpha_s}{\pi} \frac{v^2}{\Lambda^2} c_{\phi G} \ln^2 \frac{m_b^2}{m_H^2} \approx 2.4 \frac{v^2}{\Lambda^2} c_{\phi G}$$

[Note:  $c_{\phi G}$  expected to be loop-induced]

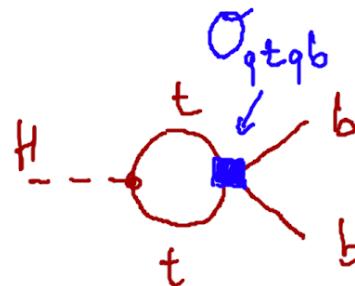
Gauld, Pecjak, Scott '16  
Cullen, Pecjak, Scott '19



## Example 2: $ttbb$ operator in $h \rightarrow b\bar{b}$

$$\frac{\delta\Gamma_{bb}}{\Gamma_{bb}} \approx 1.75 \frac{v^2}{\Lambda^2} c_{qtqb}^{(1)}$$

[Note:  $c_{qtqb}^{(1)}$  also makes a contribution to  $m_b$ , and fine-tuning arguments suggest it shouldn't be large]



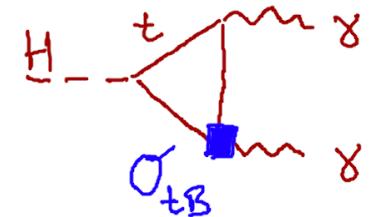
Gauld, Pecjak, Scott '15  
Cullen, Pecjak, Scott '19

Example 3: dipole-type  $\gamma$ -top operator in  $h \rightarrow \gamma\gamma$

Vryonidou, Zhang '18  
Dawson, Giardino '19

$$\frac{\delta\Gamma_{bb}}{\Gamma_{bb}} \approx \mathcal{O}(10-20) \times \frac{v^2}{\Lambda^2} c_{tB}^{(1)}$$

[Note:  $c_{tB}$  expected to be loop-induced]



## Backup slides

# SMEFT operators in Warsaw basis

Grzadkowski, Iskrzynski, Misiak, Rosiek '10

$X^3$		$\varphi^6$ and $\varphi^4 D^2$		$\psi^2 \varphi^3$	
$Q_G$	$f^{ABC} G_\mu^{A\nu} G_\nu^{B\rho} G_\rho^{C\mu}$	$Q_\varphi$	$(\varphi^\dagger \varphi)^3$	$Q_{e\varphi}$	$(\varphi^\dagger \varphi)(\bar{l}_p e_r \varphi)$
$Q_{\tilde{G}}$	$f^{ABC} \tilde{G}_\mu^{A\nu} G_\nu^{B\rho} G_\rho^{C\mu}$	$Q_{\varphi \square}$	$(\varphi^\dagger \varphi) \square (\varphi^\dagger \varphi)$	$Q_{u\varphi}$	$(\varphi^\dagger \varphi)(\bar{q}_p u_r \tilde{\varphi})$
$Q_W$	$\varepsilon^{IJK} W_\mu^{I\nu} W_\nu^{J\rho} W_\rho^{K\mu}$	$Q_{\varphi D}$	$(\varphi^\dagger D^\mu \varphi)^* (\varphi^\dagger D_\mu \varphi)$	$Q_{d\varphi}$	$(\varphi^\dagger \varphi)(\bar{q}_p d_r \varphi)$
$Q_{\tilde{W}}$	$\varepsilon^{IJK} \tilde{W}_\mu^{I\nu} W_\nu^{J\rho} W_\rho^{K\mu}$				
$X^2 \varphi^2$		$\psi^2 X \varphi$		$\psi^2 \varphi^2 D$	
$Q_{\varphi G}$	$\varphi^\dagger \varphi G_{\mu\nu}^A G^{A\mu\nu}$	$Q_{eW}$	$(\bar{l}_p \sigma^{\mu\nu} e_r) \tau^I \varphi W_{\mu\nu}^I$	$Q_{\varphi l}^{(1)}$	$(\varphi^\dagger i \overleftrightarrow{D}_\mu \varphi)(\bar{l}_p \gamma^\mu l_r)$
$Q_{\varphi \tilde{G}}$	$\varphi^\dagger \varphi \tilde{G}_{\mu\nu}^A G^{A\mu\nu}$	$Q_{eB}$	$(\bar{l}_p \sigma^{\mu\nu} e_r) \varphi B_{\mu\nu}$	$Q_{\varphi l}^{(3)}$	$(\varphi^\dagger i \overleftrightarrow{D}_\mu^I \varphi)(\bar{l}_p \tau^I \gamma^\mu l_r)$
$Q_{\varphi W}$	$\varphi^\dagger \varphi W_{\mu\nu}^I W^{I\mu\nu}$	$Q_{uG}$	$(\bar{q}_p \sigma^{\mu\nu} T^A u_r) \tilde{\varphi} G_{\mu\nu}^A$	$Q_{\varphi e}$	$(\varphi^\dagger i \overleftrightarrow{D}_\mu \varphi)(\bar{e}_p \gamma^\mu e_r)$
$Q_{\varphi \tilde{W}}$	$\varphi^\dagger \varphi \tilde{W}_{\mu\nu}^I W^{I\mu\nu}$	$Q_{uW}$	$(\bar{q}_p \sigma^{\mu\nu} u_r) \tau^I \tilde{\varphi} W_{\mu\nu}^I$	$Q_{\varphi q}^{(1)}$	$(\varphi^\dagger i \overleftrightarrow{D}_\mu \varphi)(\bar{q}_p \gamma^\mu q_r)$
$Q_{\varphi B}$	$\varphi^\dagger \varphi B_{\mu\nu} B^{\mu\nu}$	$Q_{uB}$	$(\bar{q}_p \sigma^{\mu\nu} u_r) \tilde{\varphi} B_{\mu\nu}$	$Q_{\varphi q}^{(3)}$	$(\varphi^\dagger i \overleftrightarrow{D}_\mu^I \varphi)(\bar{q}_p \tau^I \gamma^\mu q_r)$
$Q_{\varphi \tilde{B}}$	$\varphi^\dagger \varphi \tilde{B}_{\mu\nu} B^{\mu\nu}$	$Q_{dG}$	$(\bar{q}_p \sigma^{\mu\nu} T^A d_r) \varphi G_{\mu\nu}^A$	$Q_{\varphi u}$	$(\varphi^\dagger i \overleftrightarrow{D}_\mu \varphi)(\bar{u}_p \gamma^\mu u_r)$
$Q_{\varphi WB}$	$\varphi^\dagger \tau^I \varphi W_{\mu\nu}^I B^{\mu\nu}$	$Q_{dW}$	$(\bar{q}_p \sigma^{\mu\nu} d_r) \tau^I \varphi W_{\mu\nu}^I$	$Q_{\varphi d}$	$(\varphi^\dagger i \overleftrightarrow{D}_\mu \varphi)(\bar{d}_p \gamma^\mu d_r)$
$Q_{\varphi \tilde{W}B}$	$\varphi^\dagger \tau^I \varphi \tilde{W}_{\mu\nu}^I B^{\mu\nu}$	$Q_{dB}$	$(\bar{q}_p \sigma^{\mu\nu} d_r) \varphi B_{\mu\nu}$	$Q_{\varphi ud}$	$i(\tilde{\varphi}^\dagger D_\mu \varphi)(\bar{u}_p \gamma^\mu d_r)$

# SMEFT operators in Warsaw basis

Grzadkowski, Iskrzynski, Misiak, Rosiek '10

$(\bar{L}L)(\bar{L}L)$		$(\bar{R}R)(\bar{R}R)$		$(\bar{L}L)(\bar{R}R)$	
$Q_{ll}$	$(\bar{l}_p \gamma_\mu l_r)(\bar{l}_s \gamma^\mu l_t)$	$Q_{ee}$	$(\bar{e}_p \gamma_\mu e_r)(\bar{e}_s \gamma^\mu e_t)$	$Q_{le}$	$(\bar{l}_p \gamma_\mu l_r)(\bar{e}_s \gamma^\mu e_t)$
$Q_{qq}^{(1)}$	$(\bar{q}_p \gamma_\mu q_r)(\bar{q}_s \gamma^\mu q_t)$	$Q_{uu}$	$(\bar{u}_p \gamma_\mu u_r)(\bar{u}_s \gamma^\mu u_t)$	$Q_{lu}$	$(\bar{l}_p \gamma_\mu l_r)(\bar{u}_s \gamma^\mu u_t)$
$Q_{qq}^{(3)}$	$(\bar{q}_p \gamma_\mu \tau^I q_r)(\bar{q}_s \gamma^\mu \tau^I q_t)$	$Q_{dd}$	$(\bar{d}_p \gamma_\mu d_r)(\bar{d}_s \gamma^\mu d_t)$	$Q_{ld}$	$(\bar{l}_p \gamma_\mu l_r)(\bar{d}_s \gamma^\mu d_t)$
$Q_{lq}^{(1)}$	$(\bar{l}_p \gamma_\mu l_r)(\bar{q}_s \gamma^\mu q_t)$	$Q_{eu}$	$(\bar{e}_p \gamma_\mu e_r)(\bar{u}_s \gamma^\mu u_t)$	$Q_{qe}$	$(\bar{q}_p \gamma_\mu q_r)(\bar{e}_s \gamma^\mu e_t)$
$Q_{lq}^{(3)}$	$(\bar{l}_p \gamma_\mu \tau^I l_r)(\bar{q}_s \gamma^\mu \tau^I q_t)$	$Q_{ed}$	$(\bar{e}_p \gamma_\mu e_r)(\bar{d}_s \gamma^\mu d_t)$	$Q_{qu}^{(1)}$	$(\bar{q}_p \gamma_\mu q_r)(\bar{u}_s \gamma^\mu u_t)$
		$Q_{ud}^{(1)}$	$(\bar{u}_p \gamma_\mu u_r)(\bar{d}_s \gamma^\mu d_t)$	$Q_{qu}^{(8)}$	$(\bar{q}_p \gamma_\mu T^A q_r)(\bar{u}_s \gamma^\mu T^A u_t)$
		$Q_{ud}^{(8)}$	$(\bar{u}_p \gamma_\mu T^A u_r)(\bar{d}_s \gamma^\mu T^A d_t)$	$Q_{qd}^{(1)}$	$(\bar{q}_p \gamma_\mu q_r)(\bar{d}_s \gamma^\mu d_t)$
				$Q_{qd}^{(8)}$	$(\bar{q}_p \gamma_\mu T^A q_r)(\bar{d}_s \gamma^\mu T^A d_t)$
$(\bar{L}R)(\bar{R}L)$ and $(\bar{L}R)(\bar{L}R)$		B-violating			
$Q_{ledq}$	$(\bar{l}_p^j e_r)(\bar{d}_s q_t^j)$	$Q_{duq}$	$\varepsilon^{\alpha\beta\gamma} \varepsilon_{jk} [(d_p^\alpha)^T C u_r^\beta] [(q_s^\gamma)^T C l_t^k]$		
$Q_{quqd}^{(1)}$	$(\bar{q}_p^j u_r) \varepsilon_{jk} (\bar{q}_s^k d_t)$	$Q_{qqu}$	$\varepsilon^{\alpha\beta\gamma} \varepsilon_{jk} [(q_p^{\alpha j})^T C q_r^{\beta k}] [(u_s^\gamma)^T C e_t]$		
$Q_{quqd}^{(8)}$	$(\bar{q}_p^j T^A u_r) \varepsilon_{jk} (\bar{q}_s^k T^A d_t)$	$Q_{qqq}$	$\varepsilon^{\alpha\beta\gamma} \varepsilon_{jn} \varepsilon_{km} [(q_p^{\alpha j})^T C q_r^{\beta k}] [(q_s^{\gamma m})^T C l_t^n]$		
$Q_{lequ}^{(1)}$	$(\bar{l}_p^j e_r) \varepsilon_{jk} (\bar{q}_s^k u_t)$	$Q_{duu}$	$\varepsilon^{\alpha\beta\gamma} [(d_p^\alpha)^T C u_r^\beta] [(u_s^\gamma)^T C e_t]$		
$Q_{lequ}^{(3)}$	$(\bar{l}_p^j \sigma_{\mu\nu} e_r) \varepsilon_{jk} (\bar{q}_s^k \sigma^{\mu\nu} u_t)$				