Neutrino Fits, Tensions and Puzzles

Patrick Huber

Center for Neutrino Physics – Virginia Tech

SLAC Summer Institute 2019 SLAC August 12 – 23, 2019

P. Huber – VT CNP – p. 1

All theorists are liars

Neutrino physics has a rich history of anomalies:

It took 40 years for Ray Davis and John Bahcall to be taken seriously with the solar neutrino anomaly.

The atmospheric neutrino anomaly did not last quite that long, but still was labeled an anomaly till Super-K came around in 1998.

Much of the anomalous nature stemmed from theoretical prejudice: neutrinos are massless, neutrino mixing angles are small, astrophysics isn't an exact science, chemistry is really scary asf.

Of course, I happen to be a theorist ...

Why sterile?

We have measured in neutrino oscillation:

- $\Delta m_{21}^2 \sim 8 \cdot 10^{-5} \,\mathrm{eV}^2$ and $\theta_{12} \sim 1/2$
- $\Delta m^2_{31} \sim 2 \cdot 10^{-3} \,\mathrm{eV}^2$ and $\theta_{23} \sim \pi/4$
- $\theta_{13} \sim 0.16$

This implies a lower bound on the mass of the heaviest neutrino

 $\sqrt{2 \cdot 10^{-3} \,\mathrm{eV}^2} \sim 0.04 \,\mathrm{eV}$ This IS BSM physics!

Any $\Delta m^2 \gg \Delta m_{21}^2$, Δm_{31}^2 requires a 4th neutrino, BUT only three neutrinos with $m_{\nu} \leq m_Z$ couple to the Z \Rightarrow "sterile" neutrino.

Evidence in favor

Or at least at odds with a simple 3-flavor framework

- LSND $\bar{\nu}_{\mu} \rightarrow \bar{\nu}_{e}$
- MiniBooNE $\bar{\nu}_{\mu} \rightarrow \bar{\nu}_{e}$ and $\nu_{\mu} \rightarrow \nu_{e}$
- Reactors $\nu_e \rightarrow \nu_e$
- Gallium $\nu_e \rightarrow \nu_e$

LSND and MiniBooNE





 $P(\bar{\nu}_{\mu} \rightarrow \bar{\nu}_{e}) \simeq 0.003$

Statistically significant: $4 - 6\sigma$

Fermilab SBN



Figure courtesy D. Schmitz and C. Adams Signal to noise not so different from LSND... will a near detector of completely different design help?





Pion decay at rest at JSNS, Gd-doped scintillator. JSNS2, 2017

Direct test of the LSND result \rightarrow should have been done 20 years ago!

The reactor anomaly



Daya Bay, 2014

Mueller *et al.*, 2011, 2012 – where are all the neutrinos gone?

Contributors to the anomaly

6% deficit of $\bar{\nu}_e$ from nuclear reactors at short distances

- 3% increase in reactor neutrino fluxes
- decrease in neutron lifetime
- inclusion of long-lived isotopes (non-equilibrium correction)

The effects is therefore only partially due to the fluxes, but the error budget is clearly dominated by the fluxes.

Neutron lifetime



Neutrinos from fission



β -branches



β -spectrum from fission



²³⁵U foil inside the High Flux Reactor at ILL

Electron spectroscopy with a magnetic spectrometer

Same method used for ²³⁹Pu and ²⁴¹Pu

For ²³⁸U recent measurement by Haag *et al.*, 2013

Schreckenbach, et al. 1985.

A priori calculations



Updated β -feeding functions from total absorption γ spectroscopy (safe from pandemonium) for key isotopes.

For ²³⁸, ²³⁹Pu and ²⁴¹Pu better than 5% agreement with beta decay data

²³⁵U, the odd-one-out?

Estienne et al., 2019

Forbidden decays



 $e,\overline{\nu}$ final state can form a singlet or triplet spin state J=0 or J=1

Allowed: s-wave emission (l = 0)Forbidden: p-wave emission (l = 1)or l > 1

Significant nuclear structure dependence in forbidden $decays \rightarrow large unquantifiable uncertainties!$

P. Huber – VT CNP – p. 15

Look at past data

a	Experiment	f^{a}_{235}	f^{a}_{238}	f^{a}_{239}	f^{a}_{241}	$R_{a,\mathrm{SH}}^{\mathrm{exp}}$	σ^{\exp}_{a} [%]	$\sigma_a^{ m cor}$ [%]	L_a [m]
1	Bugey-4	0.538	0.078	0.328	0.056	0.932	1.4	1.4	15
2	Rovno91	0.606	0.074	0.277	0.043	0.930	2.8	1.8	18
3	Rovno88-11	0.607	0.074	0.277	0.042	0.907	6.4	3.8	18
4	Rovno88-2I	0.603	0.076	0.276	0.045	0.938	6.4	3.8	18
5	Rovno88-1S	0.606	0.074	0.277	0.043	0.962	7.3	3.8	18
6	Rovno88-2S	0.557	0.076	0.313	0.054	0.949	7.3	3.8	25
7	Rovno88-3S	0.606	0.074	0.274	0.046	0.928	6.8	3.8	18
8	Bugey-3-15	0.538	0.078	0.328	0.056	0.936	4.2	4.1	15
9	Bugey-3-40	0.538	0.078	0.328	0.056	0.942	4.3	4.1	40
10	Bugey-3-95	0.538	0.078	0.328	0.056	0.867	15.2	4.1	95
11	Gosgen-38	0.619	0.067	0.272	0.042	0.955	5.4	3.8	37.9
12	Gosgen-46	0.584	0.068	0.298	0.050	0.981	5.4	3.8	45.9
13	Gosgen-65	0.543	0.070	0.329	0.058	0.915	6.7	3.8	64.7
14	ILL	1	0	0	0	0.792	9.1	8.0	8.76
15	Krasnoyarsk87-33	1	0	0	0	0.925	5.0	4.8	32.8
16	Krasnoyarsk87-92	1	0	0	0	0.942	20.4	4.8	92.3
17	Krasnoyarsk94-57	1	0	0	0	0.936	4.2	2.5	57
18	Krasnoyarsk99-34	1	0	0	0	0.946	3.0	2.5	34
19	SRP-18	1	0	0	0	0.941	2.8	0.0	18.2
20	SRP-24	1	0	0	0	1.006	2.9	0.0	23.8
21	Nucifer	0.926	0.061	0.008	0.005	1.014	10.7	0.0	7.2
22	Chooz	0.496	0.087	0.351	0.066	0.996	3.2	0.0	pprox 1000
23	Palo Verde	0.600	0.070	0.270	0.060	0.997	5.4	0.0	pprox 800
24	Daya Bay	0.561	0.076	0.307	0.056	0.946	2.0	0.0	≈ 550
25	RENO	0.569	0.073	0.301	0.056	0.946	2.1	0.0	≈ 410
26	Double Chooz	0.511	0.087	0.340	0.062	0.935	1.4	0.0	≈ 415

Giunti, 2016

What does this tell us?



Giunti, 2016

Is ²³⁵U odd? Are the error bars for ²³⁵U just smaller?

Latest result of Daya Bay



Only an issue if the prediction of Pu239 in the Huber+Mueller model is correct. Hayes *et al.*, 2017

Daya Bay, 2017 and Diaz et al., 2019

The 5 MeV bump



Double Chooz 2019 Seen by all three reactor experiments Tracks reactor power





NEOS, 2016

24m from a large core (power reactor), confirms bump

NEOS vs Daya Bay



Huber, 2017

There is more U235 in NEOS, since core is fresh \Rightarrow 3 - 4 σ evidence against Pu as sole source of bump, but equal bump size is still allowed at better than 2 σ .

Latest data vs bump



PROSPECT 2018 Disfavors 235 U as sole culprit at 2.1 σ



Daya Bay 2019 Requires a bump in 235 U at 4σ

Explanations?

Direct summation of latest ENSDF database with allowed beta-spectrum shape Sonzogni *et al.*, 2016



This direct summation, as all other direct summations, does not agree with the Schreckenbach measurement.

What happened?

Fission yield data has been suspected previously Hayes *et al.* 2015 and this what Sonzogni *et al.*, 2016 found:



Who is the odd-one-out?

Fission yields for germanium-86 wrong in ENDF/B but not in JEFF.

BSM explanation for the bump



Berryman, Brdar, PH, 2018

Requires a sterile neutrino consistent with the reactor anomaly and a new vector state X coupling to quarks.

Does it work?



Berryman, Brdar, PH, 2018

Excellent fit

Existence of highenergy neutrino flux is predicted

High energy flux is in agreement with Daya Bay bounds

Position and width of bump entirely determined by SM physics

Is it allowed?



Shown is axial coupling case, vector case even worse.

COHERENT data and old D_2O reactor data can not be reconciled with this model.

The bump remains unexplained with either BSM or nuclear physics!

NEOS and sterile neutrinos



NEOS reports a limit, but their best fit occurs at $\sin^2 2\theta = 0.05$ and $\Delta m^2 = 1.73 \,\mathrm{eV}^2$ with a χ^2 value 6.5 below the no-oscillation hypothesis.

adapted from NEOS, 2016 DANSS has a similar result.

DANSS and NEOS



Dentler et al. 2018

This is a spectral effect independent of rate and shape predictions!

Reactor fit



More than 3σ evidence for oscillation even without using a flux prediction!

Dentler et al. 2018

Gallium anomaly

	GAI	LLEX	SAGE		
k	G1	G2	S 1	S2	
source	⁵¹ Cr	⁵¹ Cr	⁵¹ Cr	³⁷ Ar	
R^k_{B}	0.953 ± 0.11	$0.812^{+0.10}_{-0.11}$	0.95 ± 0.12	$0.791 \pm {}^{+0.084}_{-0.078}$	
$R_{ m H}^k$	$0.84^{+0.13}_{-0.12}$	$0.71_{-0.11}^{+0.12}$	$0.84_{-0.13}^{+0.14}$	$0.70 \pm {+0.10 \atop -0.09}$	
radius [m]	1	.9	0.7		
height [m]	5	5.0	1.47		
source height [m]	2.7	2.38	0.72		

25% deficit of ν_e from radioactive sources at short distances

- Effect depends on nuclear matrix element
- R is a calibration constant

Nuclear matrix elements



Nuclear matrix elements – II



Kostensalo *et al.* 2019 Significance decreases from 3.0σ to 2.3σ , but all results in the $\nu_e/\bar{\nu}_e$ sector are fully consistent!

Disappearance and appearance

 $\nu_{\mu} \rightarrow \nu_{e}$ requires that the sterile neutrino mixes with both ν_{e} and ν_{μ}

 \Rightarrow there must be effects in *both* $\nu_e \rightarrow \nu_e$ and $\nu_\mu \rightarrow \nu_\mu$

Up to factors of 2, the energy averaged probabilities obey

$$P_{\mu e} \lesssim (1 - P_{\mu \mu})(1 - P_{ee})$$

Disappearance data



 $\sin^2 2\theta_{e\mu} = 4|U_{e4}U_{\mu4}|^2$ with $1 - P_{ee} \propto |U_{e4}|^2$ and $1 - P_{\mu\mu} \propto |U_{\mu4}|^2$

Dentler, *et al.*, 2018

There is (and has been for decades) a strong tension between global appearance data and disappearance data.

Global fit



Gariazzo *et al.*, 2017 Caveat: not entirely up-to-date.

Finding a sterile neutrino

All pieces of evidence have in common that they are less than 5σ effects and they may be all due to the extraordinary difficulty of performing neutrino experiments, if not:

- N sterile neutrinos are the simplest explanation
- Tension with null results in disappearance remains

Due to their special nature as SM gauge singlets sterile neutrinos are strong candidates for being a portal to a hidden sector – significant experimental activity.

$\nu_{\mathbf{e}}$ outlook



Current $\nu_e/\bar{\nu}_e$ data points to a region at the edge of the sensitivity current reactor experiments.

Shown are the allowed regions from a fit to reactor data after the year 2010 separated into rate and spectrum. The reactor flux model uses the latest *ab initio* results from Estienne *et al.* for the central value and the Huber-Mueller results for the error bars. The gallium region is from Kostensalo *et al.*.

P. Huber – VT CNP – p. 38

Summary

Tension in global fits

- Maybe more complicated than sterile neutrino
- And/or not all data is right
- Lots of nuclear physics uncertainties

With NEOS and DANSS we have a positive indication from reactors independent of flux predictions.

In combination, light sterile neutrinos are one of the best cases for New Physics, anywhere!





P. Huber – VT CNP – p. 40